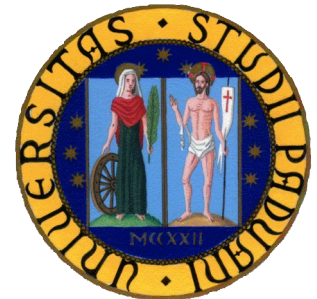




# Experimental Results on Flavor Physics



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On Behalf of LHCb Collaboration  
Presenting also results from BaBar and Belle

Workshop on High Energy Physics Phenomenology  
4-14 December 2015, Kanpur , India

# Outline

- Motivation for studying flavor physics
- CP violation and CKM physics
  - $\sin(2\beta)$  in  $B \rightarrow J/\psi K_s$  [Bfact+LHCb]
  - $\sin(2\beta)$  in  $\bar{B} \rightarrow D_{CP^{(*)}} h^0$  [Bfact]
  - $\Delta m_d$  [LHCb]
  - $|V_{ub}/V_{cb}|$  from  $\Lambda_b \rightarrow p \mu \nu$  [LHCb]
- Rare Decays
  - $B(s) \rightarrow \mu^+ \mu^-$  BF combination (CMS LHCb)
  - $B \rightarrow K^* \mu^+ \mu^-$  and  $\phi \mu^+ \mu^-$  angular analysis (LHCb)
- Universality
  - $R_K$  in  $B \rightarrow K l^+ l^-$  (LHCb)
  - $R_D$  e  $R_{D^*}$  in  $B \rightarrow D^* \tau \nu$  (BaBar, Belle, LHCb)

# Motivation for studying flavor physics

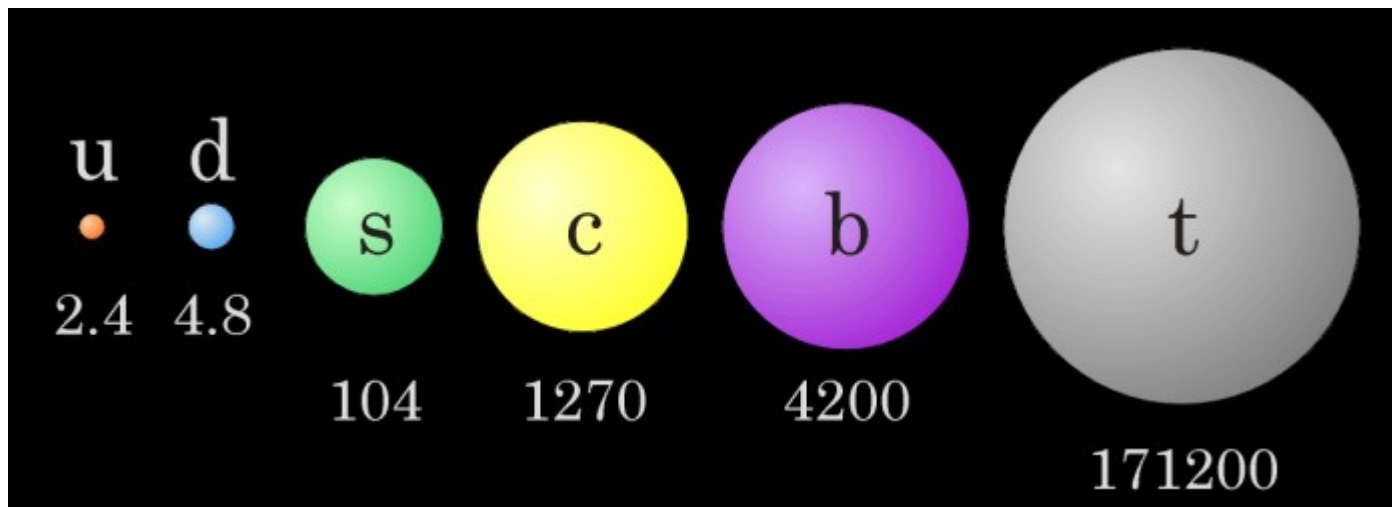
- In the Standard model the Yukawa couplings of the Higgs boson to the fermions determine the quark mixing matrix
- 10 out of 19 of the free parameters of the model are the masses and the mixing parameters of the quarks. Their study gives us fundamental informations. But there are also some fundamental open questions
- Why are there 3 generations?

**Elementary Particles**

Quarks	<b><i>u</i></b> up	<b><i>c</i></b> charm	<b><i>t</i></b> top	Force Carriers
	<b><i>d</i></b> down	<b><i>s</i></b> strange	<b><i>b</i></b> bottom	
Leptons	<b><math>\nu_e</math></b> <i>e</i> neutrino	<b><math>\nu_\mu</math></b> $\mu$ neutrino	<b><math>\nu_\tau</math></b> $\tau$ neutrino	<b><i>W</i></b> <i>W</i> boson
	<b><i>e</i></b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau	<b><i>Z</i></b> <i>Z</i> boson
3 →	I	II	III	← Generations

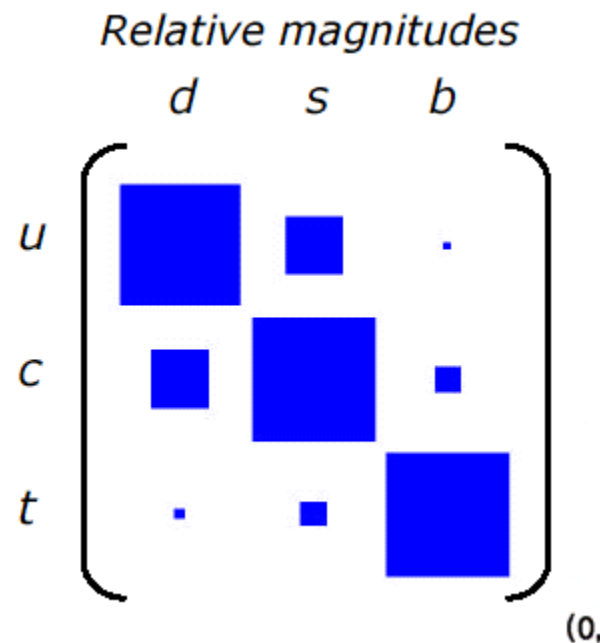
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- 10 out of 19 of the free parameters of the model are the masses and the mixing parameters of the quarks. Their study gives us fundamental informations. But there are also some open fundamental questions
- Why there is a striking hierarchy in the quark masses? Why the Higgs mass is at the EW scale?



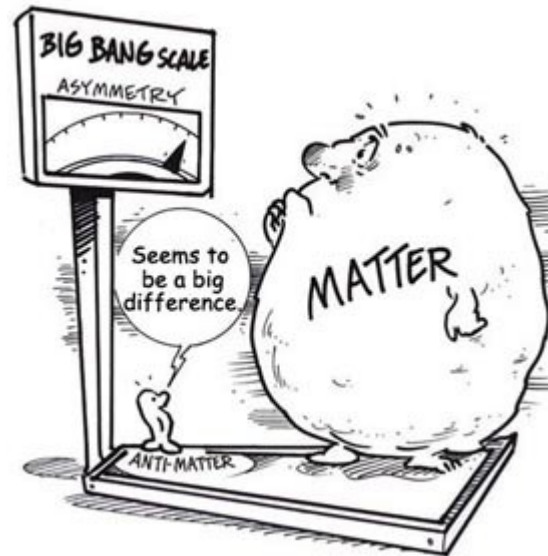
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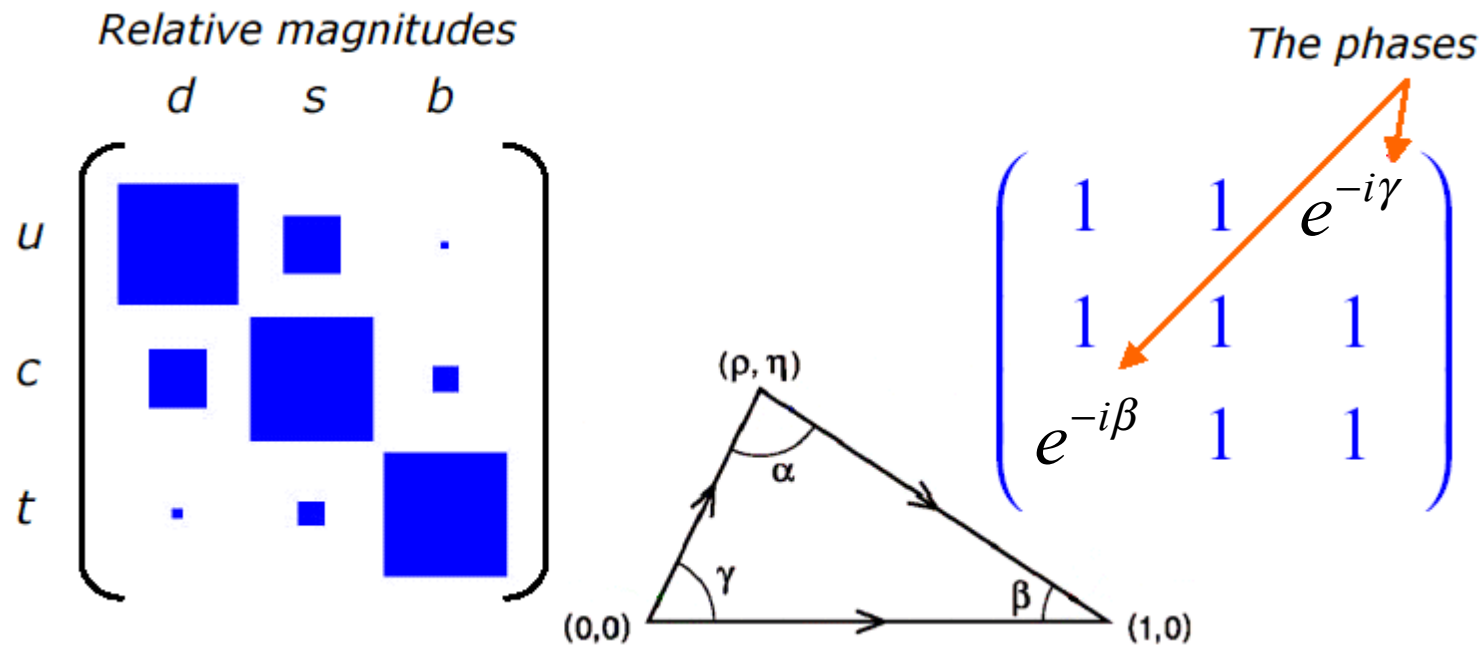
# Motivation for studying flavor physics

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- 10 out of 19 of the free parameters of the model are the masses and the mixing parameters of the quarks. Their study gives us fundamental informations. But there are also some open fundamental questions
- Why is there a large matter anti matter asymmetry in the universe?



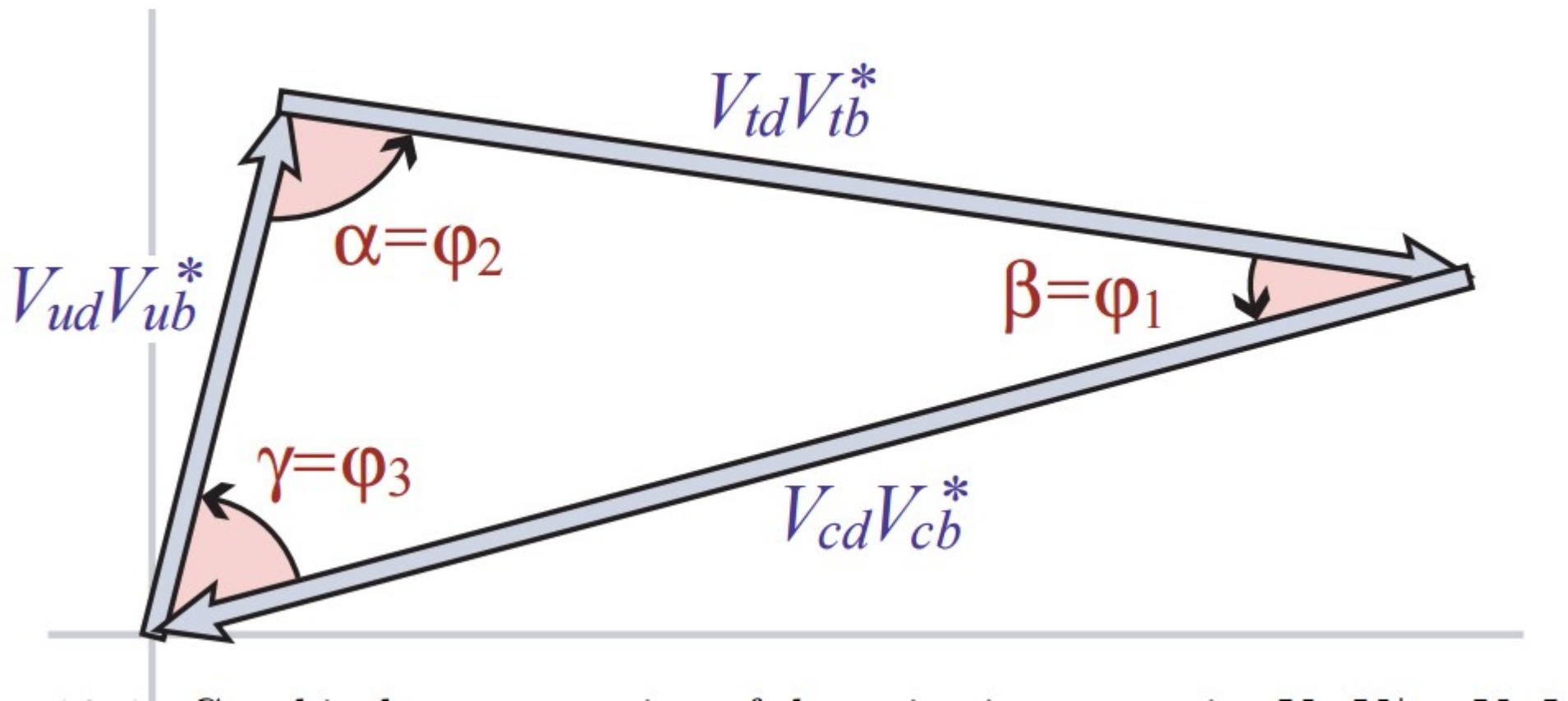
# CP Violation

- Baryon asymmetry requires large CP violation  $\frac{\Delta n_B^{expected}}{n_\gamma} = 10^{-17}$  while  $\Delta \frac{n_B}{n_\gamma} = 10^{-9}$
- In the standard model CP violation is described in by a single weak phase in the quark mixing (CKM) matrix



# The Unitarity triangle

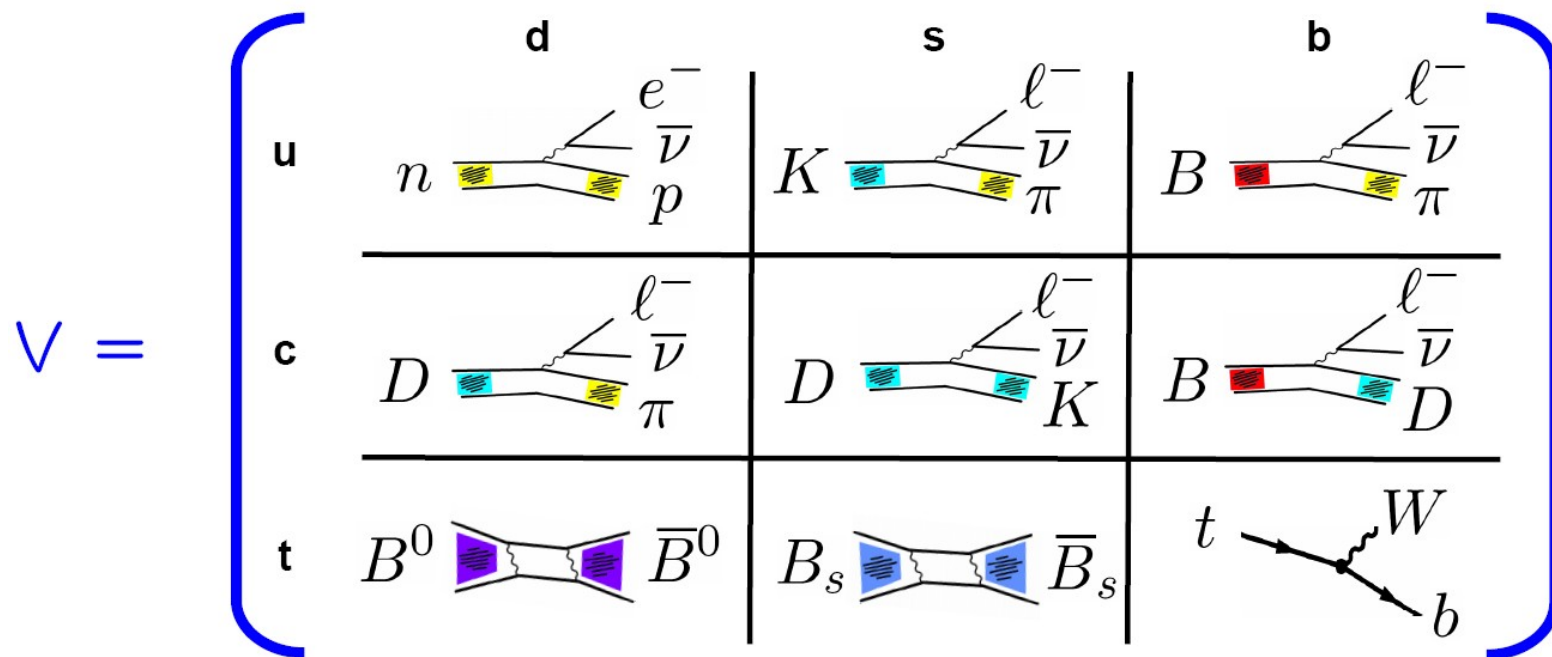
- The unitarity of the matrix can be represented as a triangle relation between its elements
- Measuring angles and sides of the triangle and over-constrains the model => consistency checks





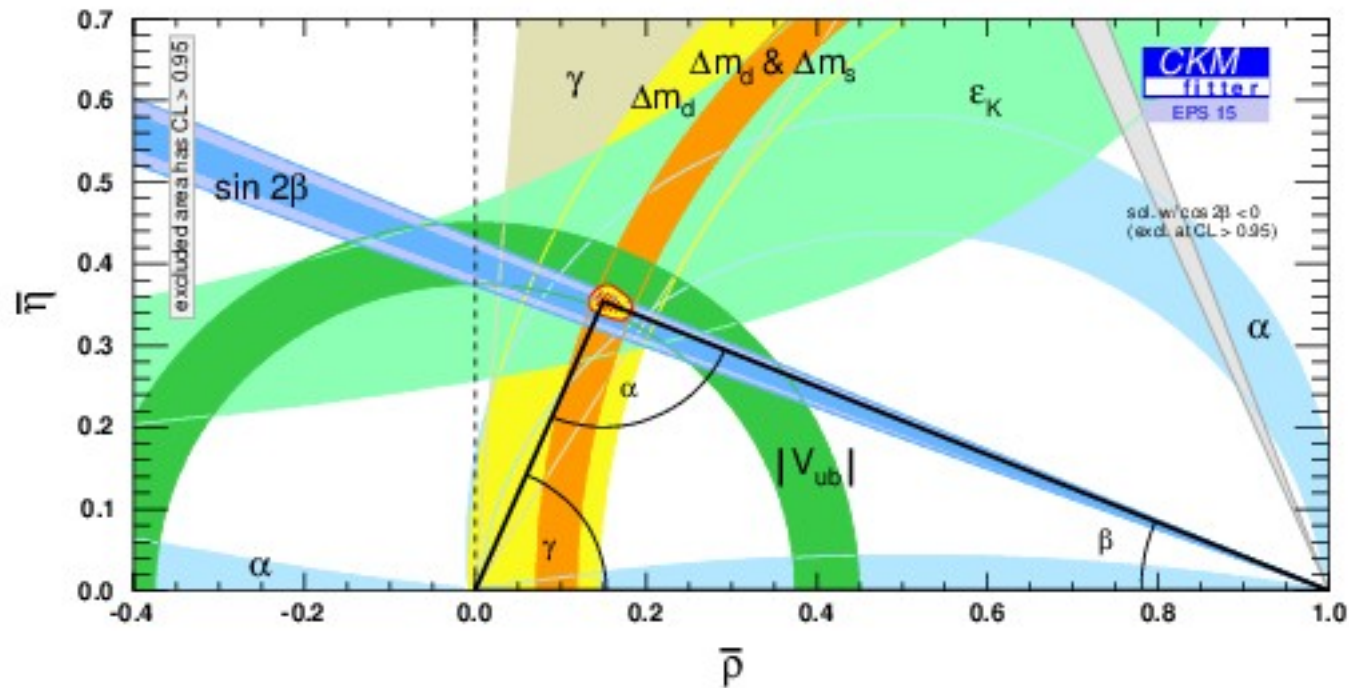
# CKM Matrix elements

- How can the sides and angles of the triangle be measured?
- Quarks can change flavor only through charged current interactions proportional to elements of the CKM matrix
- Interference effects between amplitudes give us information on the angles and sides



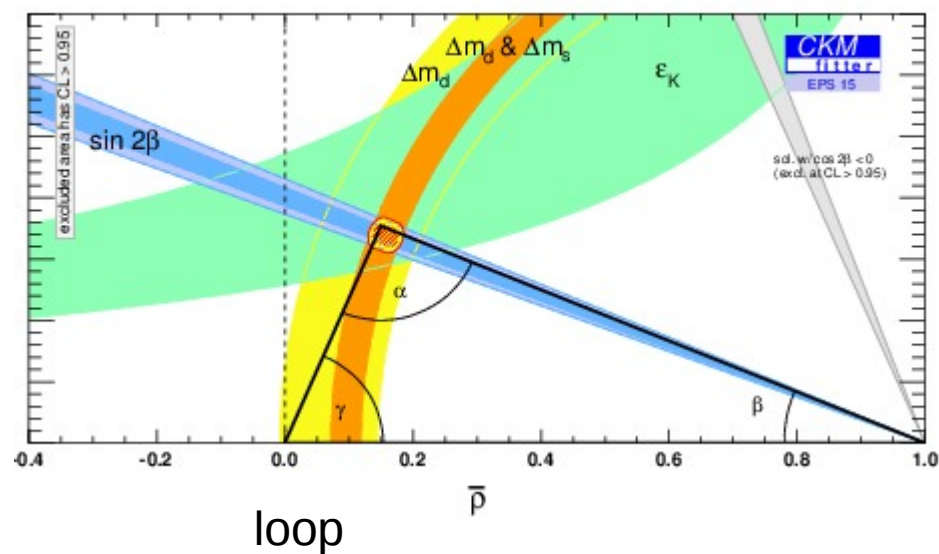
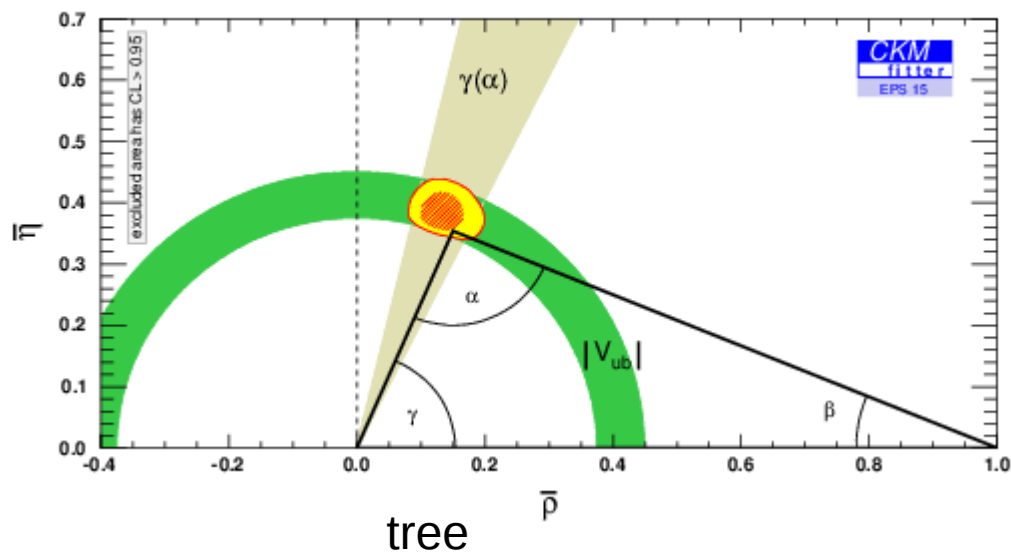
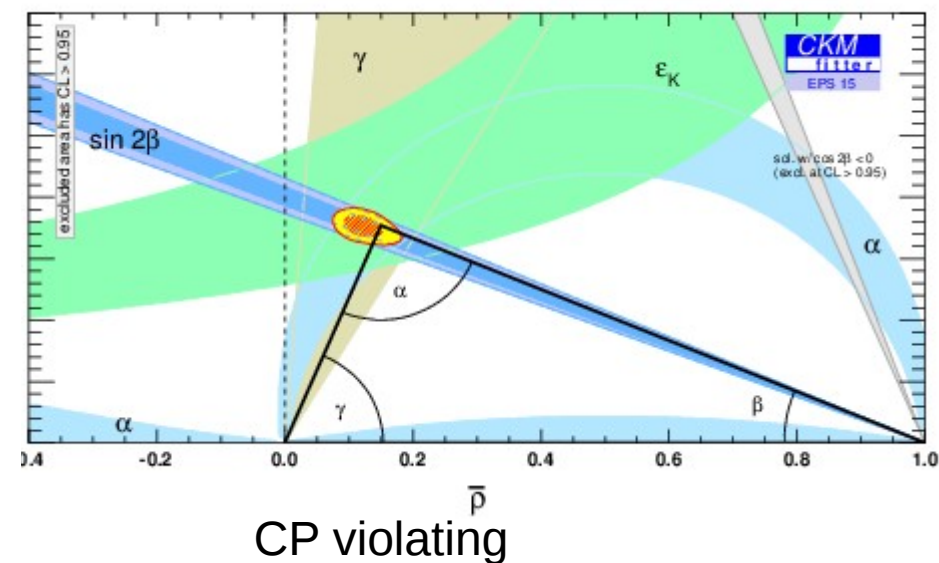
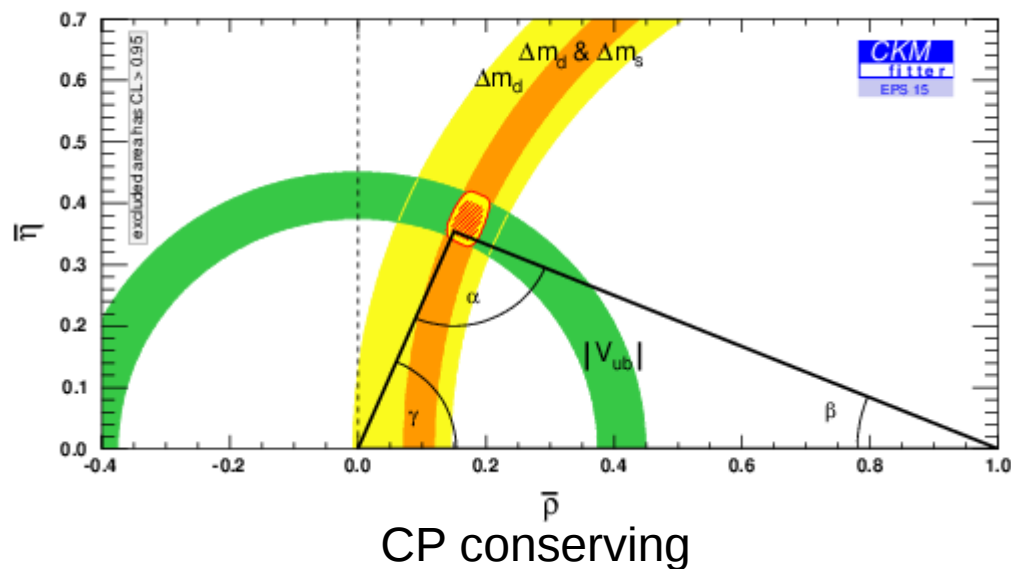
# The Unitarity triangle

- The CKM mechanism is extremely successful in describing the observed pattern of decays, mixing and CP violation



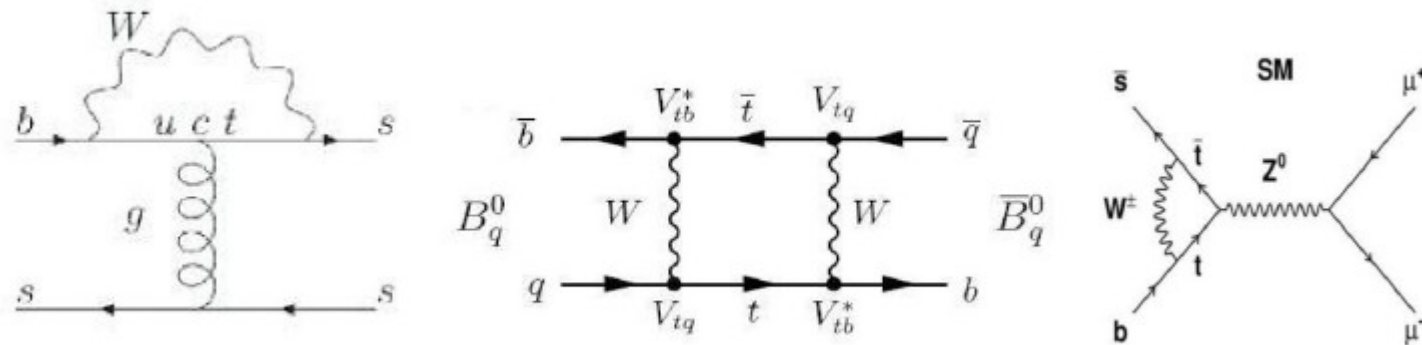
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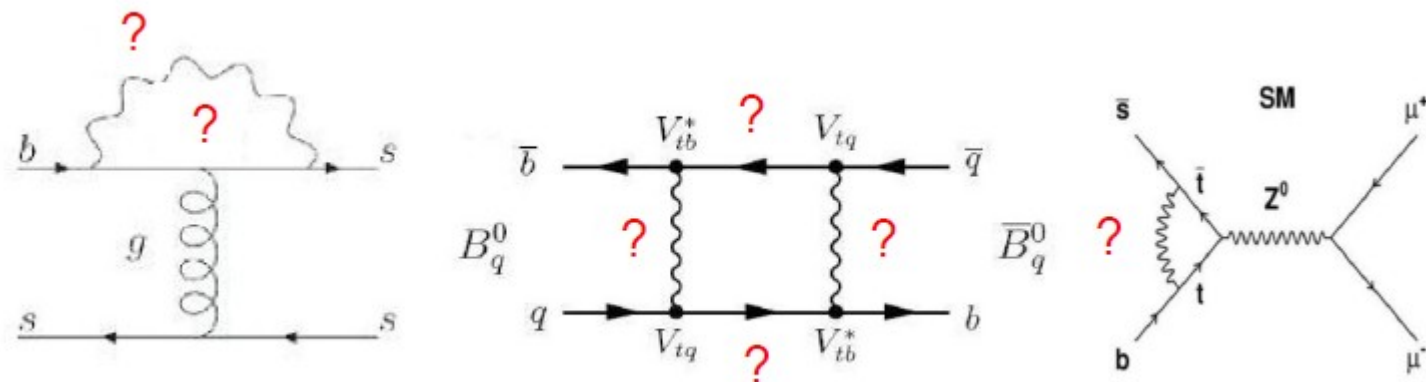


# Rare Decays

- In tree decays however new physics effects are loop suppressed
- Flavor changing neutral currents are already highly suppressed in the SM

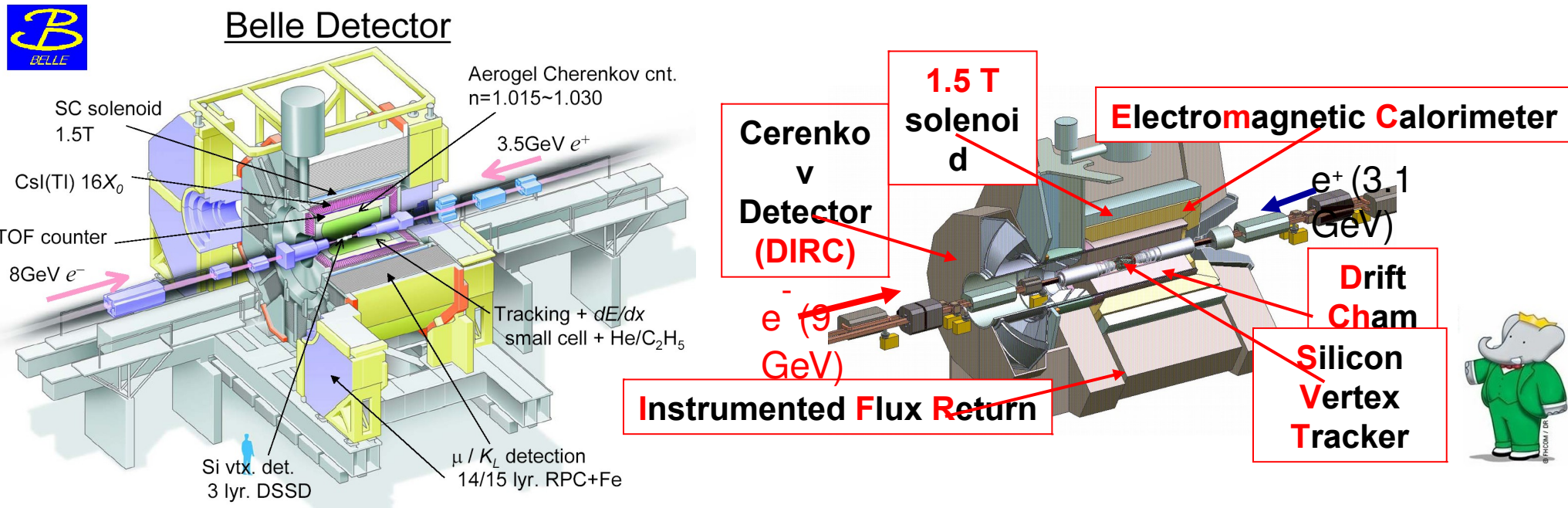


- Therefore new intermediate particles can appear in the loops giving unsuppressed contributions  $\Rightarrow$  clean signatures of New Physics
- In addition they can probe higher energy scales inaccessible with direct searches



# BFactories

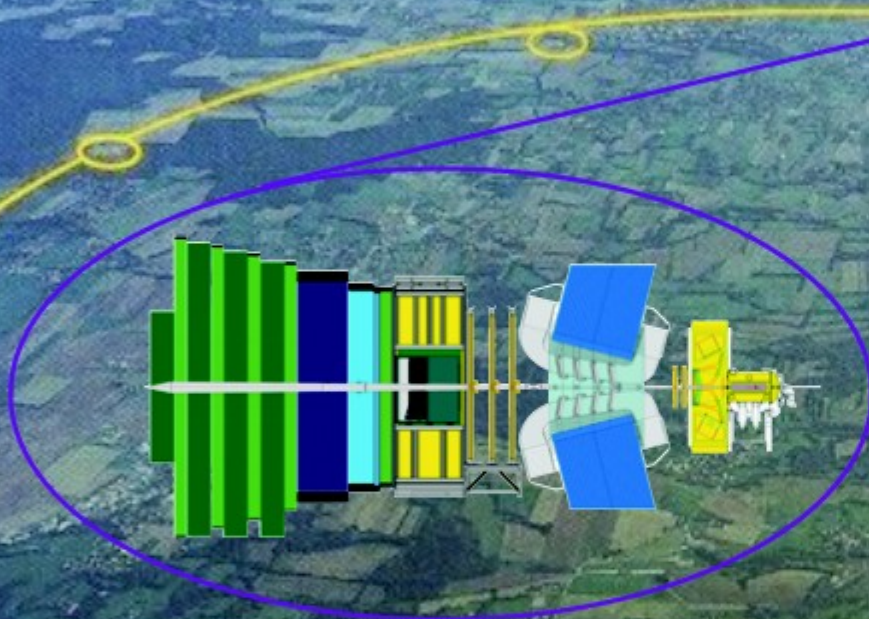
- Bfactories provide a clean environment to measure CKM parameters and rare decays
- $B\bar{B}$  pairs are produced at threshold from the  $\Upsilon(4S)$  resonance
- $\sigma(b\bar{b}) \sim 1.1 \text{ nb}$  and  $\sigma(q\bar{q}) \sim 3 \text{ nb}$ 
  - Almost hermetic detector, highly efficient trigger



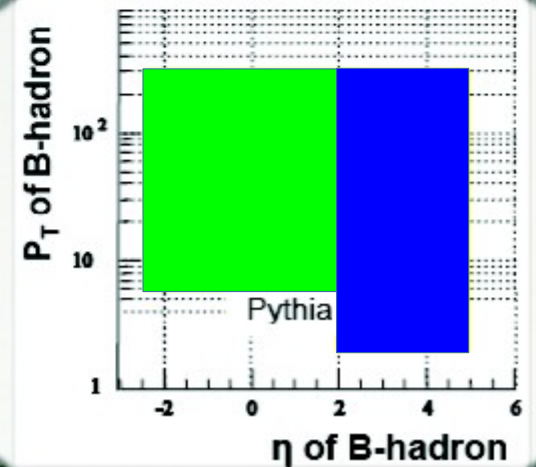
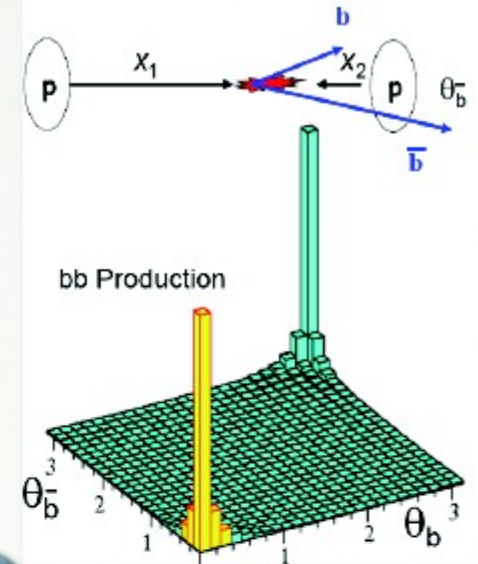


# The LHCb experiment

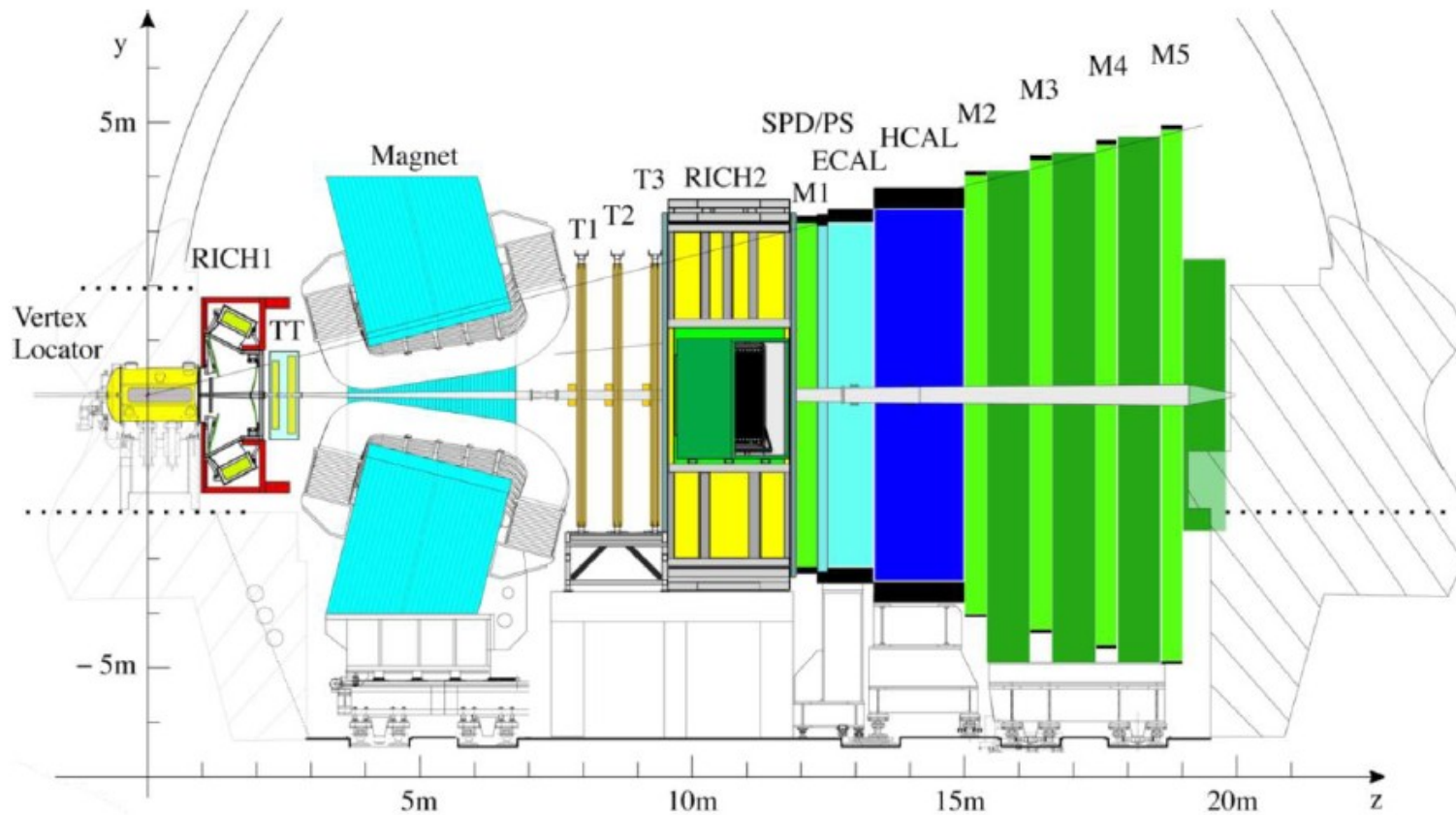
$$\begin{aligned}\sigma(pp \rightarrow \text{inel}) &= 60 \text{ mb} \\ \sigma(pp \rightarrow c\bar{c}) &= 6 \text{ mb} \\ \sigma(pp \rightarrow b\bar{b}) &= 0.3 \text{ mb}\end{aligned}$$



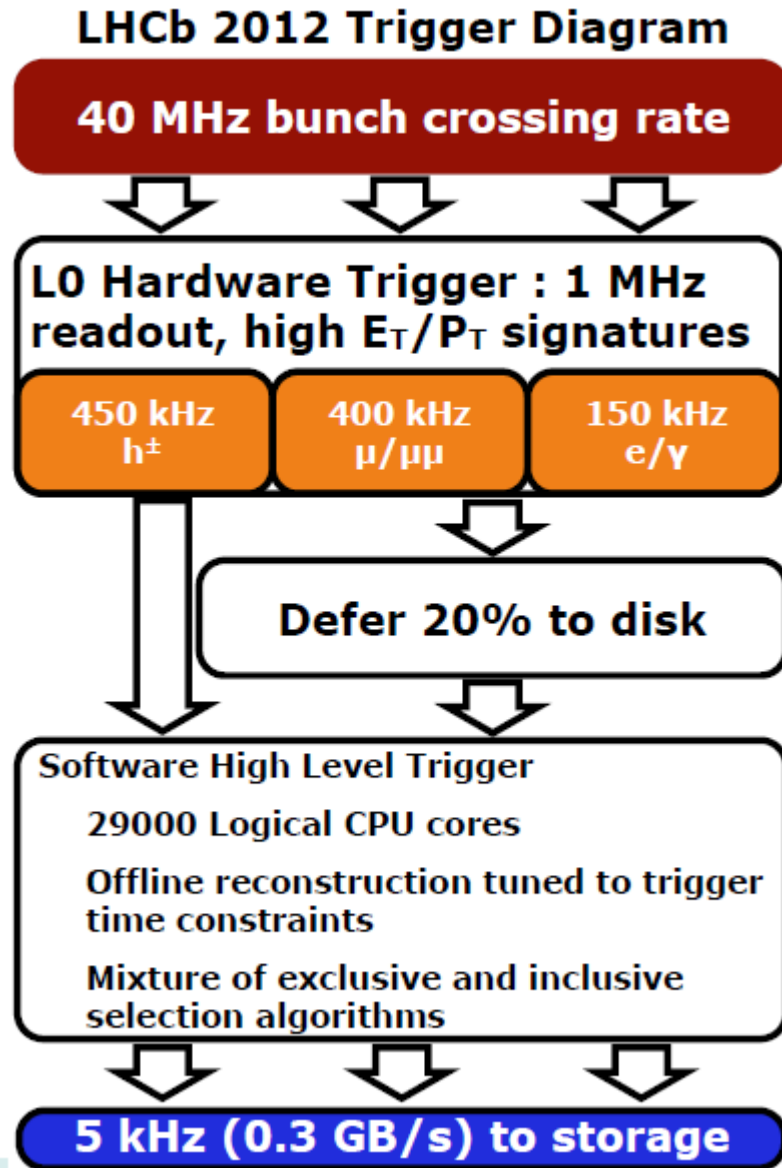
Gluon-Gluon-Fusion:







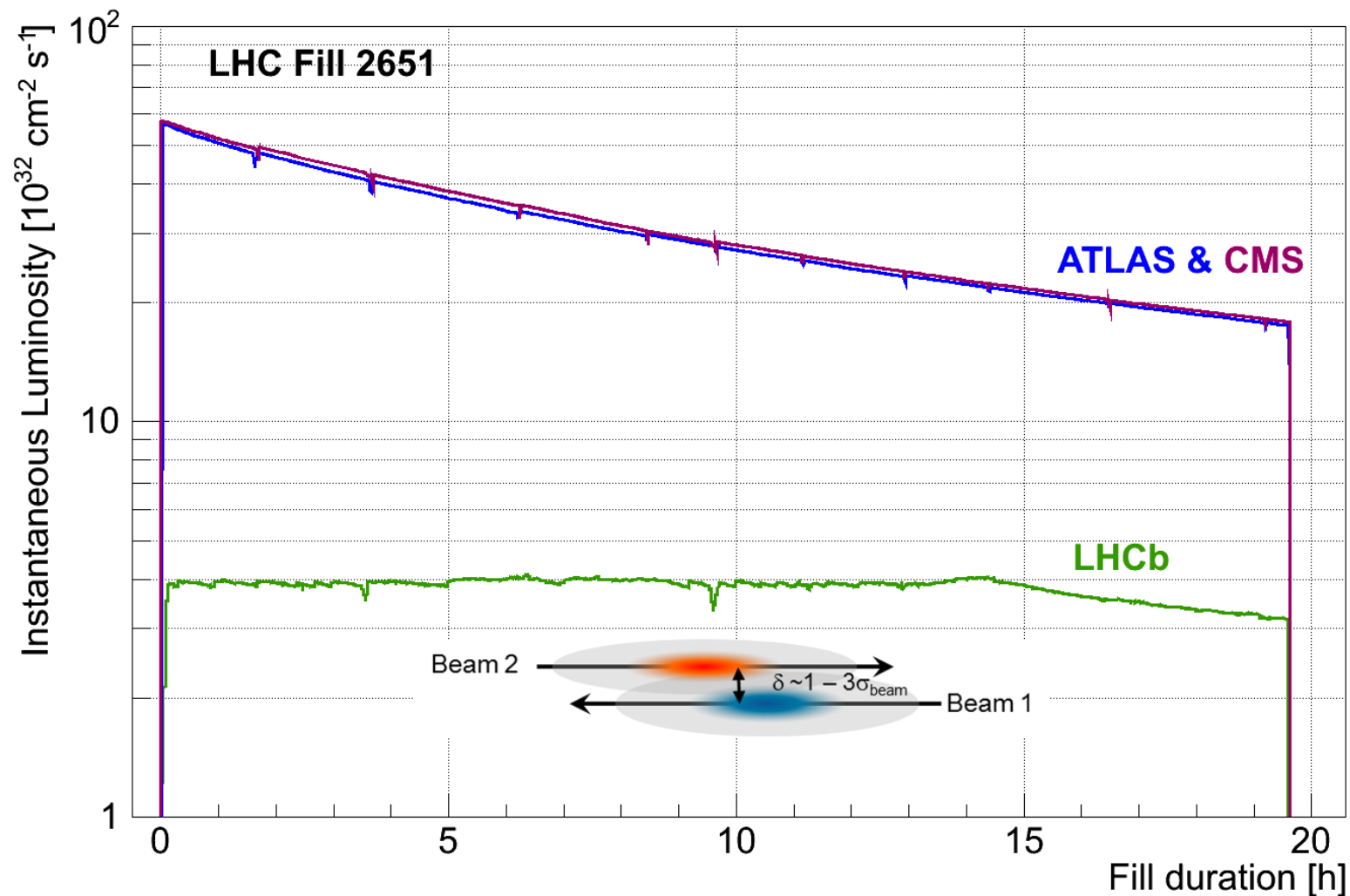
- Requirements:
  - Vertex resolution to resolve  $B_s$  oscillations
  - Excellent PID for flavor tagging and bkg rejection in rare decays
  - Momentum resolution to separate  $B/B_s, D/D_s$



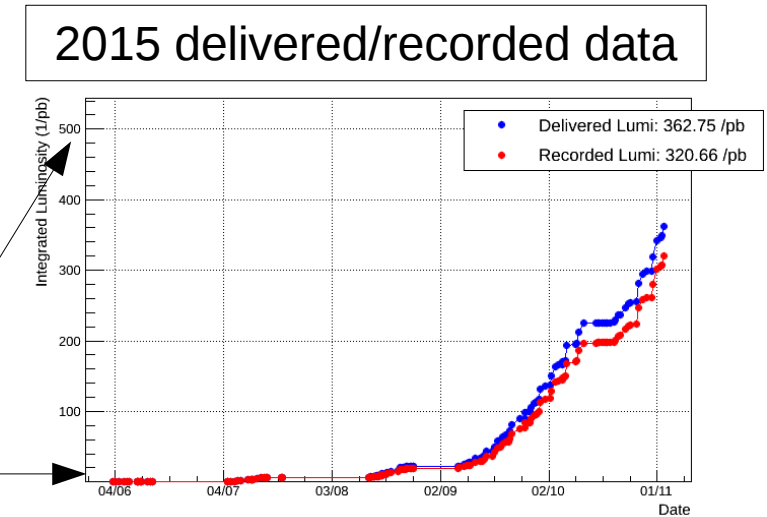
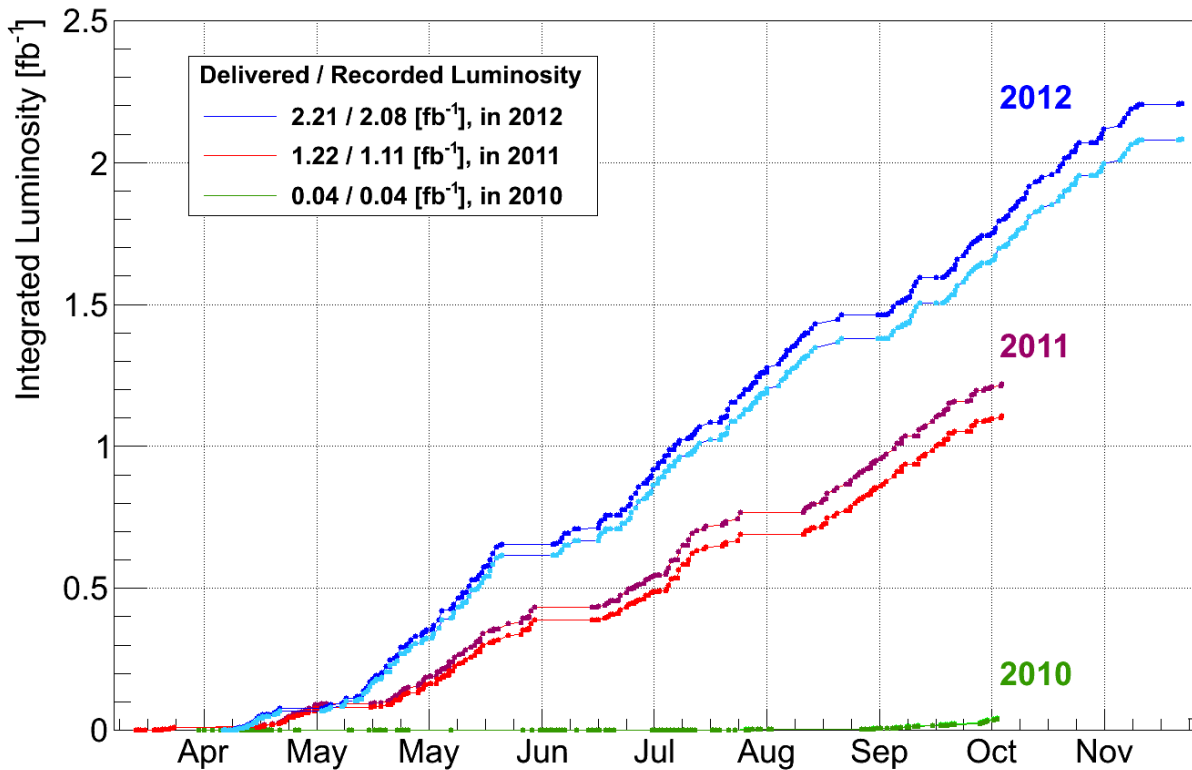
- High efficiency for manageable rates
- Background
  - Minimum bias
  - Other  $B$  and  $D$  decays
- Trigger variables
  - High  $p_t$  tracks
  - High impact parameter and displaced vertices
- Upgraded in Run II
  - Perform real-time detector calibrations and alignment
  - Full offline-like event selection
  - Turbo stream (5kHz)



- LHCb tracking/PID sensitive to pileup and occupancy
- Reduce luminosity by displacing the beams from head-on collisions



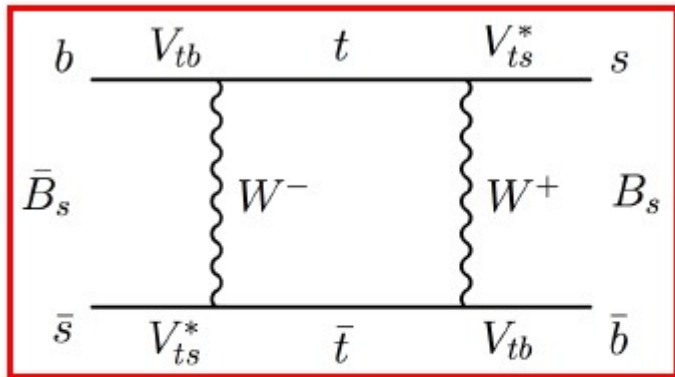
- LHC run 1 went from 2010 to 2012, during which LHCb collected 3 fb<sup>-1</sup> of data (this corresponds to  $\sim 3 \times 10^{11}$  b anti-b pairs being produced within LHCb).



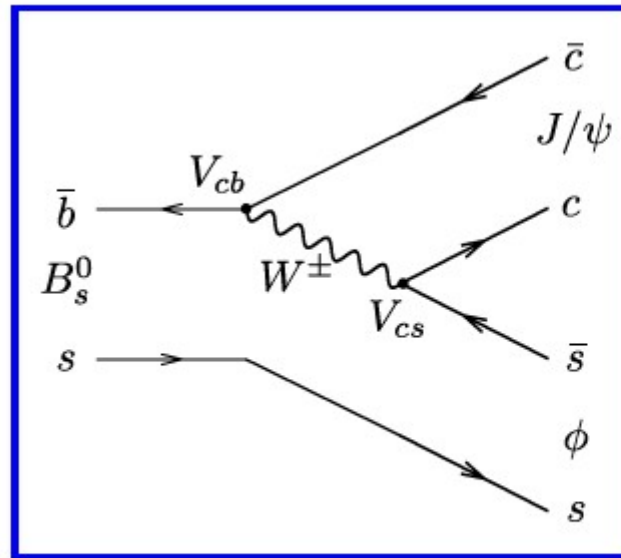
- Run 2 started after a 2 year shutdown, necessary to upgrade LHC energy.
- Will continue to end of 2018 ~we aim to increase our beauty sample by x3 or more

$\text{Sin}(2\beta)$  with  $B \rightarrow J/\psi K_s$

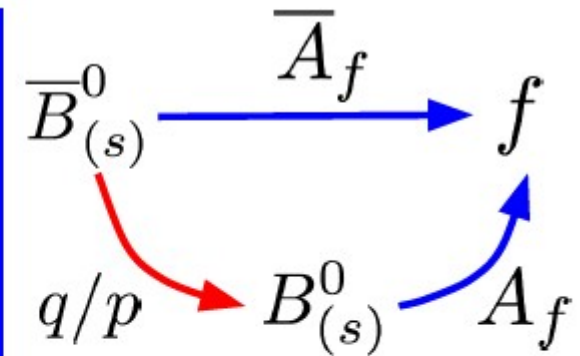
# How to observe CP violation



$$\phi_{mix} = 2 \arg(V_{tb} V_{ts}^*)$$



$$\phi_{dec} = \arg(V_{cb} V_{cs}^*)$$



$$|\lambda_f| \equiv \left| \frac{q}{p} \frac{A_f}{\bar{A}_f} \right| \approx 1$$

$$A_f = |a_1| e^{i(\delta_1 + \phi_1)} + |a_2| e^{i(\delta_2 + \phi_2)},$$

$$\bar{A}_f = |a_1| e^{i(\delta_1 - \phi_1)} + |a_2| e^{i(\delta_2 - \phi_2)}.$$

$$A_{CP}(t) \equiv \frac{\Gamma_{B^0 \rightarrow f} - \Gamma_{\bar{B}^0 \rightarrow f}}{\Gamma_{B^0 \rightarrow f} + \Gamma_{\bar{B}^0 \rightarrow f}} = \frac{S_f \sin(\Delta m t) - C_f \cos(\Delta m t)}{\cosh(\Delta \Gamma t/2) + A_{\Delta \Gamma} \sinh(\Delta \Gamma t/2)}$$

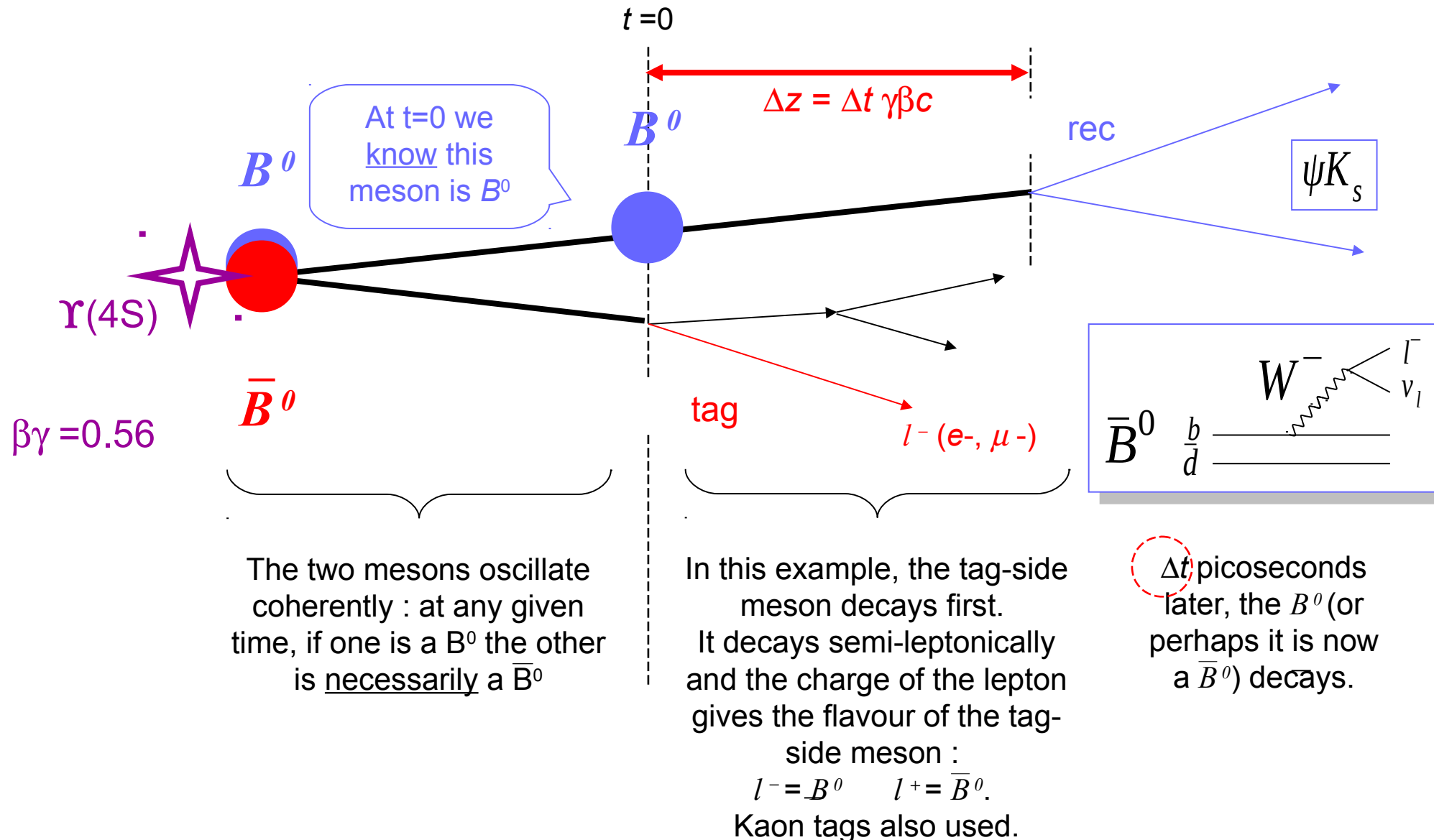
$$C_f \equiv \frac{1 - |\lambda_f|^2}{1 + |\lambda_f|^2}$$

$$S_f \equiv \frac{2 \sin \phi_{d,s}}{1 + |\lambda_f|^2}$$

$$A_{\Delta \Gamma} \equiv -\frac{2 \cos \phi_{d,s}}{1 + |\lambda_f|^2}$$

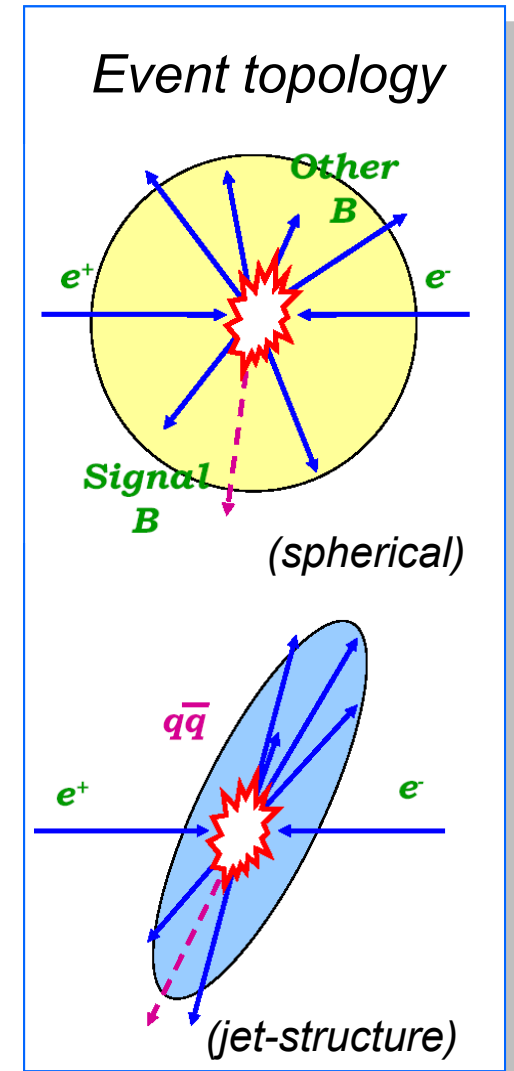
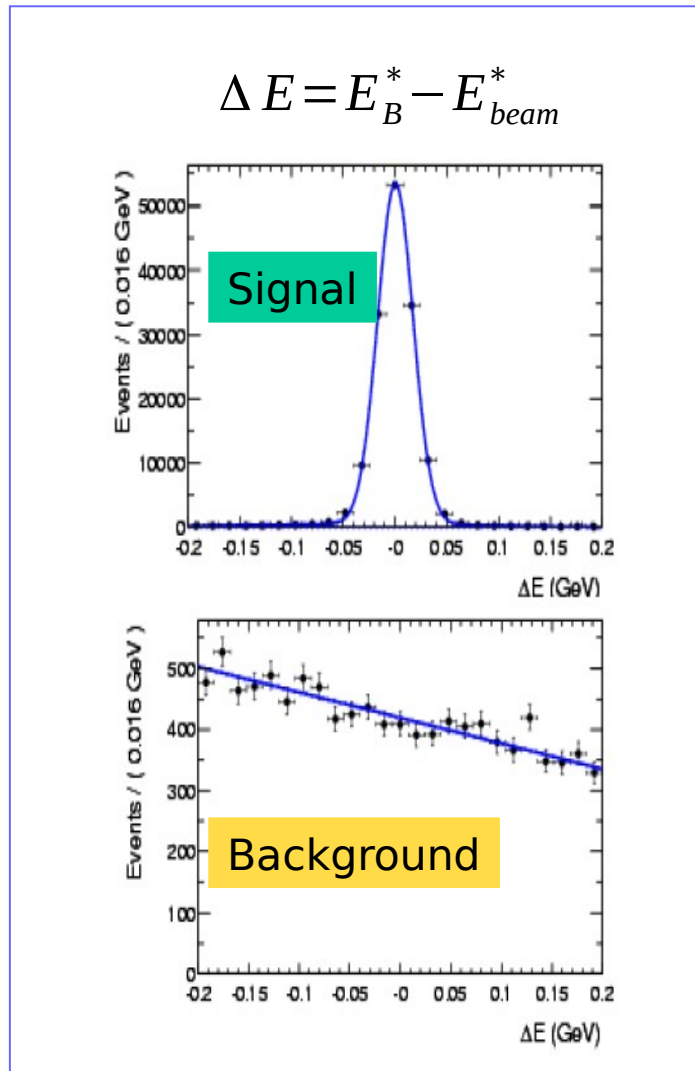
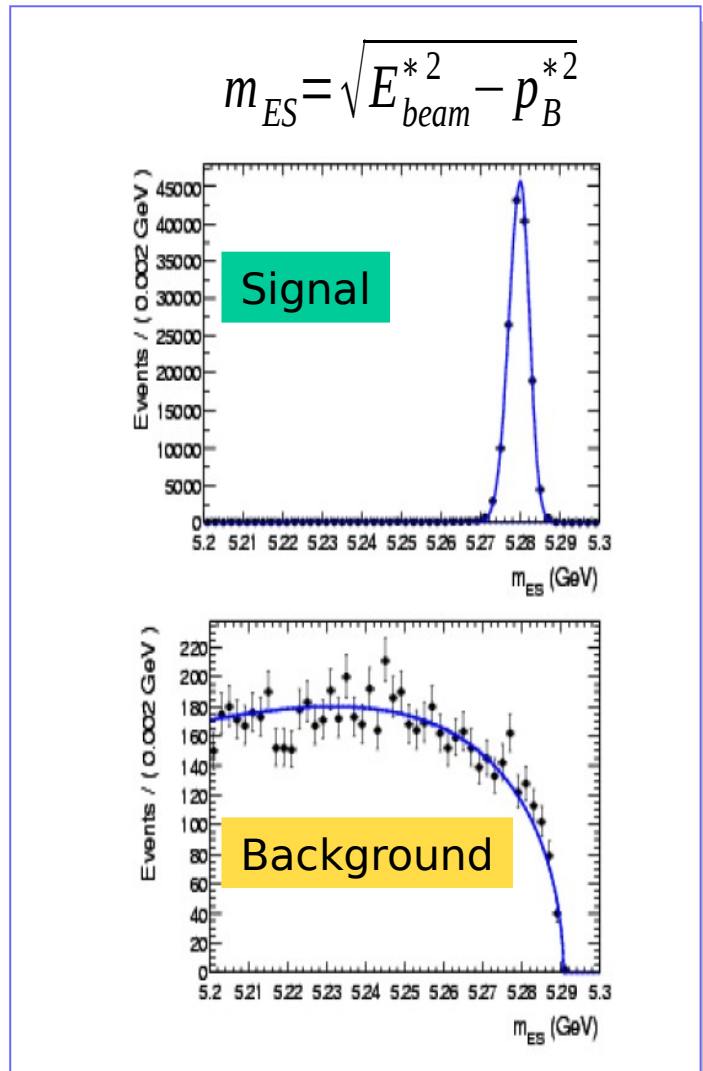
# Measurement of time dependent rates at Bfactories

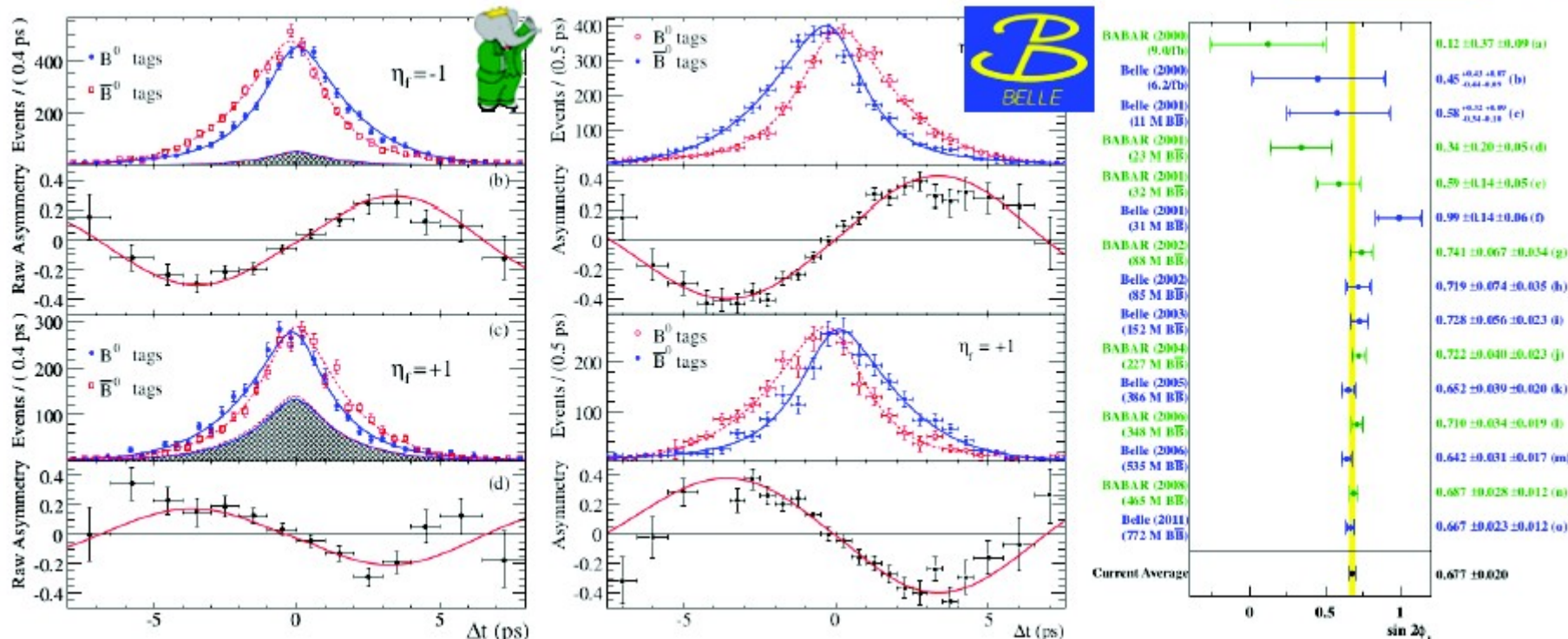
- Time dependent CP violation measured employing entanglement of initial states to know the flavor of the B at the reference time  $t=0$



# Key analysis techniques at B-Factories

Threshold kinematics: the initial energy of the  $Y(4S)$  system is known, and therefore the approximate energy and magnitude of momentum of each B

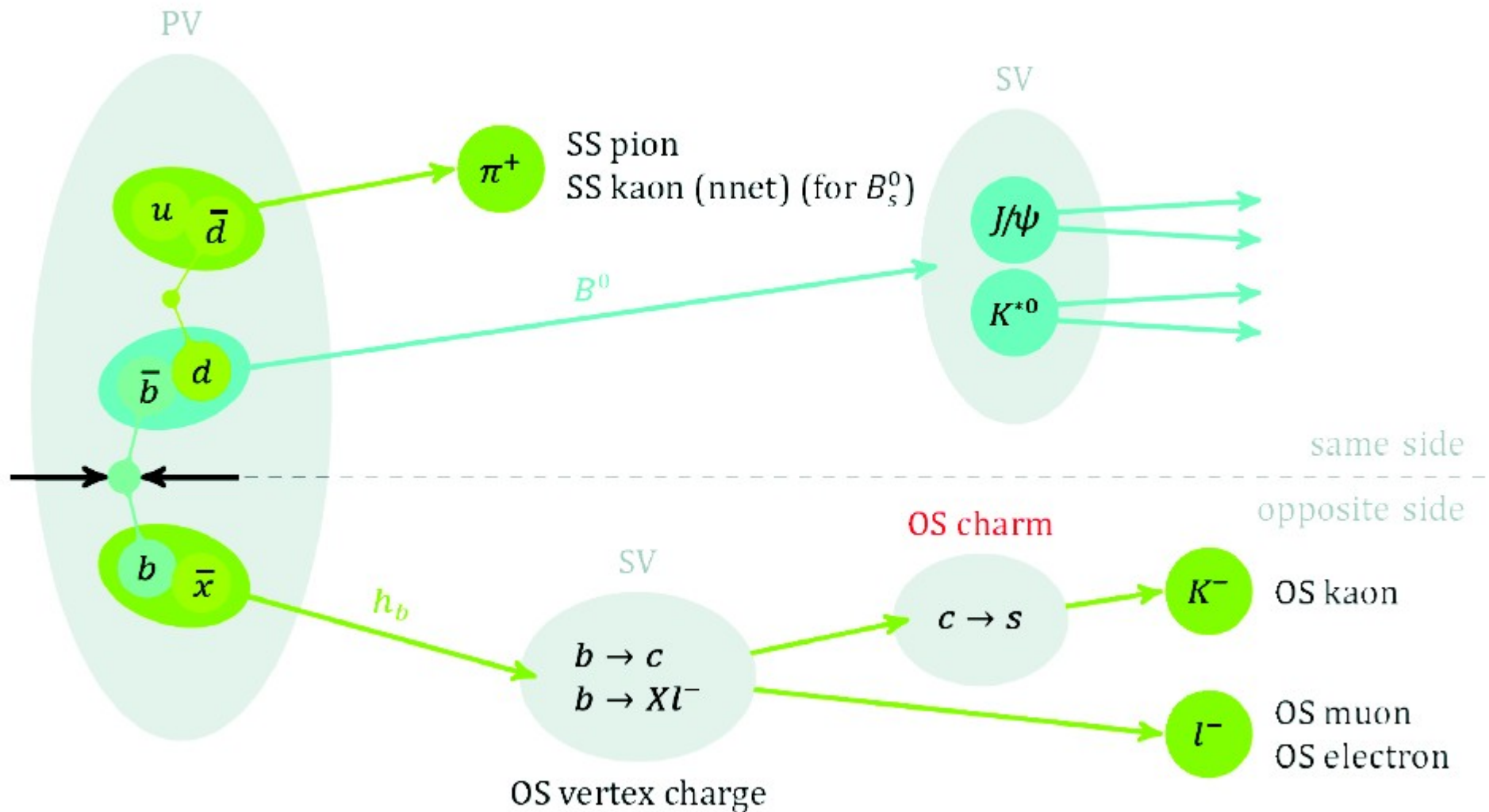




- BaBar:
  - $-\eta_f S_f = 0.687 \pm 0.028 \pm 0.012$  [PRD 79, 072009 (2009)]
  - $C_f = 0.024 \pm 0.020 \pm 0.016$
- Belle:
  - $-\eta_f S_f = 0.667 \pm 0.023 \pm 0.012$  [PRL 108, 171802 (2012)]
  - $C_f = 0.006 \pm 0.016 \pm 0.012$ .



# TD CPV at the LHC: flavor tagging

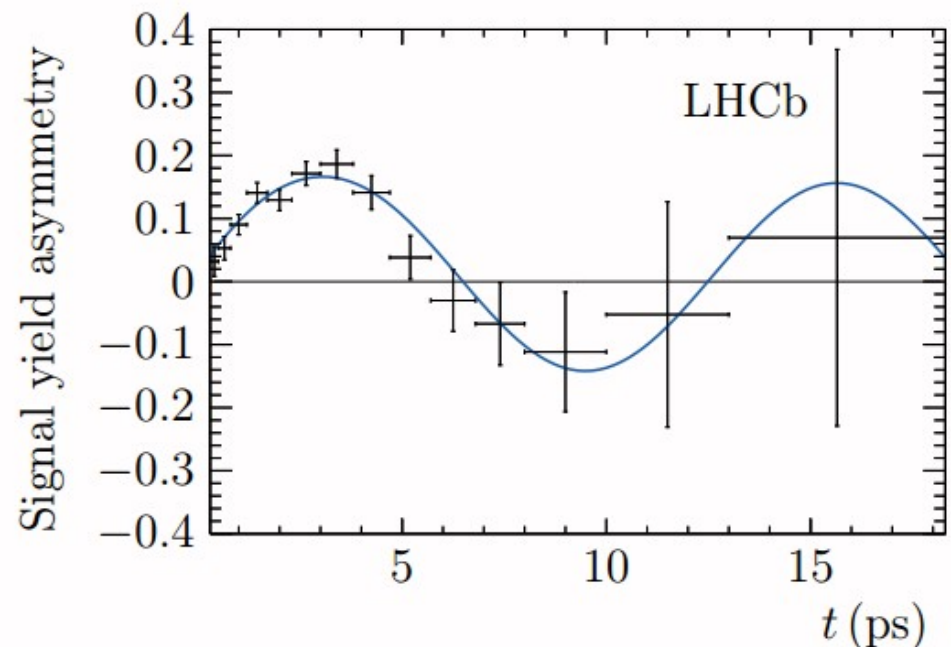
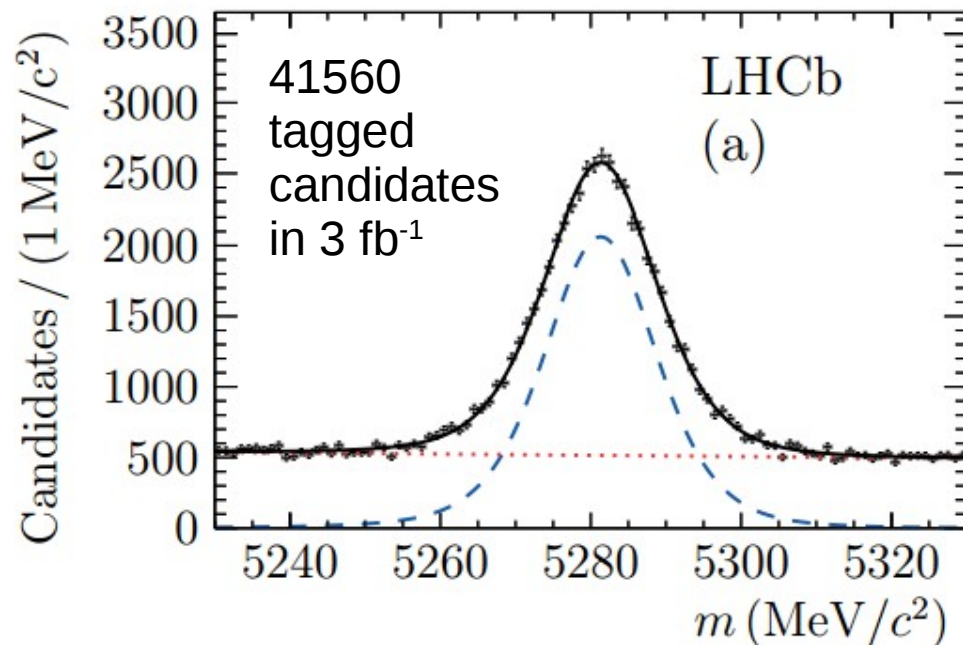


- OS tagging obtained similarly to BFactories, however
- no correlation between the evolution of the two B hadrons => intrinsic dilution => intrinsically small effective tagging efficiency



- Tagging using OS tagger
- Inclusion of SS pion tagger →  
eff=2.4% → 3%
  - Calibrated using B → J/ψ K\*<sub>0</sub>

- Time dependence of acceptance determined from data at small times (trigger effects) and MC for large times (reconstruction effects)
- $-\eta_f S_f = 0.731 \pm 0.035 \pm 0.020$ ,
- $C_f = -0.038 \pm 0.032 \pm 0.005$ ,



- Tagging using OS tagger
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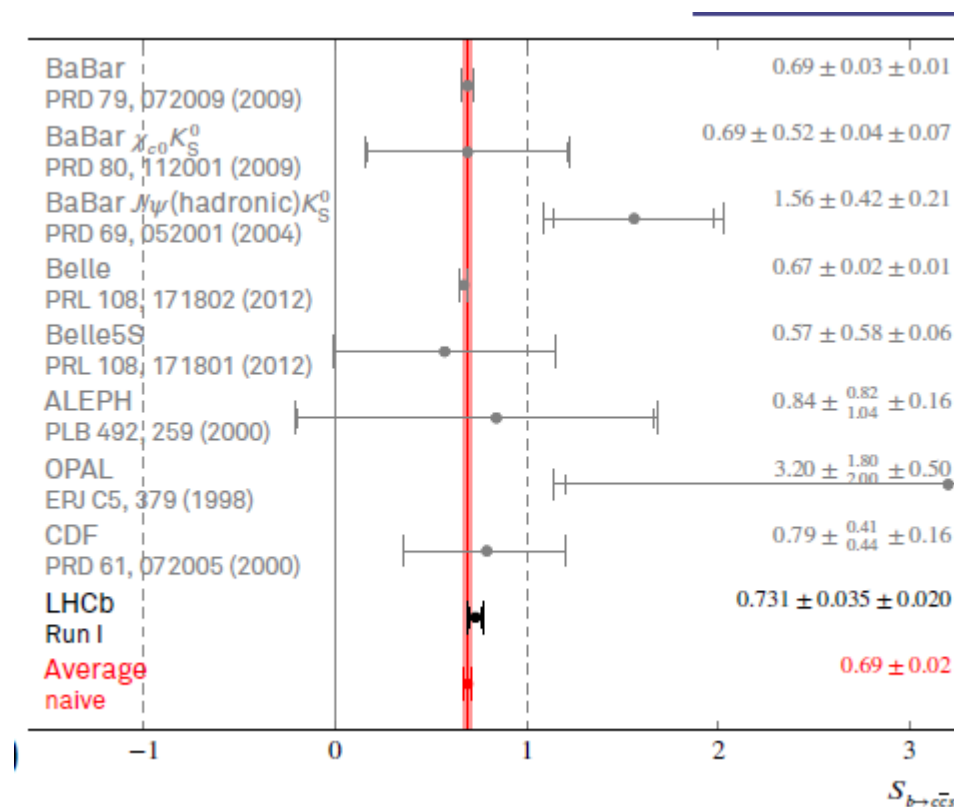
- Time dependence of acceptance determined from data at small times (trigger effects) and MC for large times (reconstruction effects),

- Now competitive with B-Factories
- B-Factories still better precision

$$S^{LHCb} = 0.731 \pm 0.035 \pm 0.020,$$

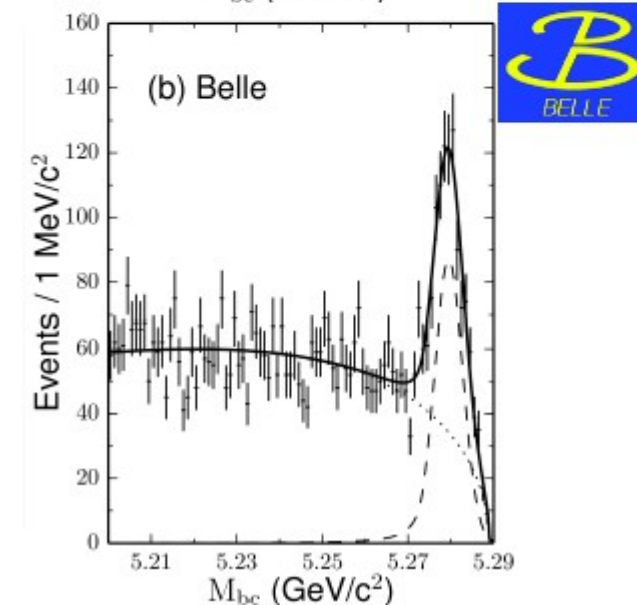
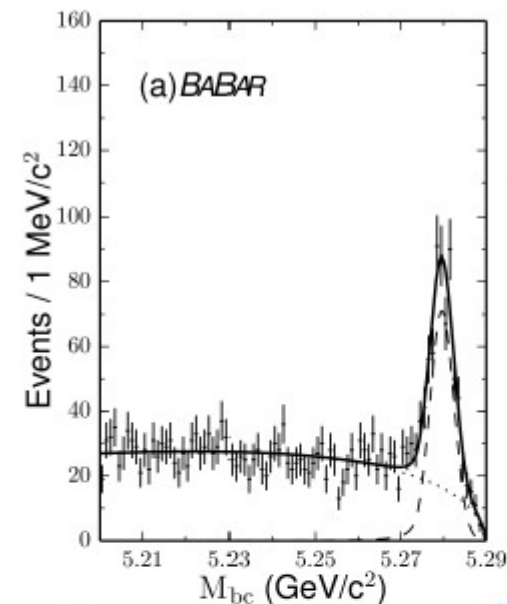
$$S^{BaBar} = 0.663 \pm 0.039 \pm 0.012,$$

$$S^{BELLE} = 0.670 \pm 0.029 \pm 0.013$$



- $\bar{B} \rightarrow D^{(*)} h^0$  decays proceed mainly through tree amplitudes  $q \rightarrow \bar{c} u d$
- If  $D$  decays to  $CP$  eigenstate  $K^+ K^-$ ,  $K_S \pi^0$ ,  $K_S \omega$  with  $h^0 = \pi^0, \eta, \omega$  then time dependent  $CP$  violation can occur via interference with oscillation
- Penguin free measurement of  $\sin(2\beta)$  that can be used as a reference for mixing induced  $CP$  violation in penguin mediated  $b \rightarrow s$   $B$  meson decays
- First combined fit to BaBar and Belle data ( $1.243 \cdot 10^9$   $B\bar{B}$  pairs)

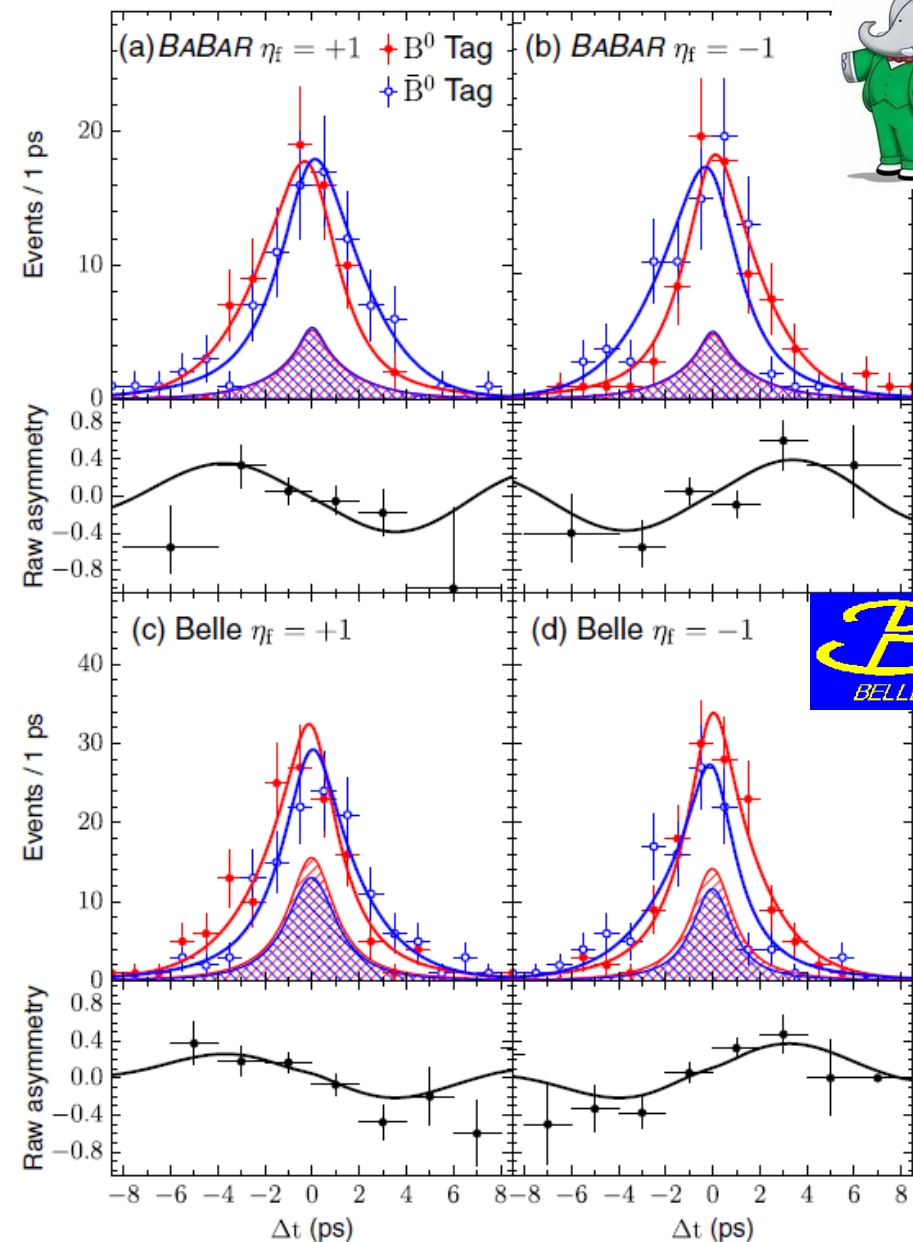
	BaBar	Belle
• $\bar{B}^0 \rightarrow D_{CP}^{(*)} h^0$ total	$508 \pm 31$	$757 \pm 44$



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$$-\eta_f \mathcal{S} = +0.66 \pm 0.10(\text{stat}) \pm 0.06(\text{syst}),$$

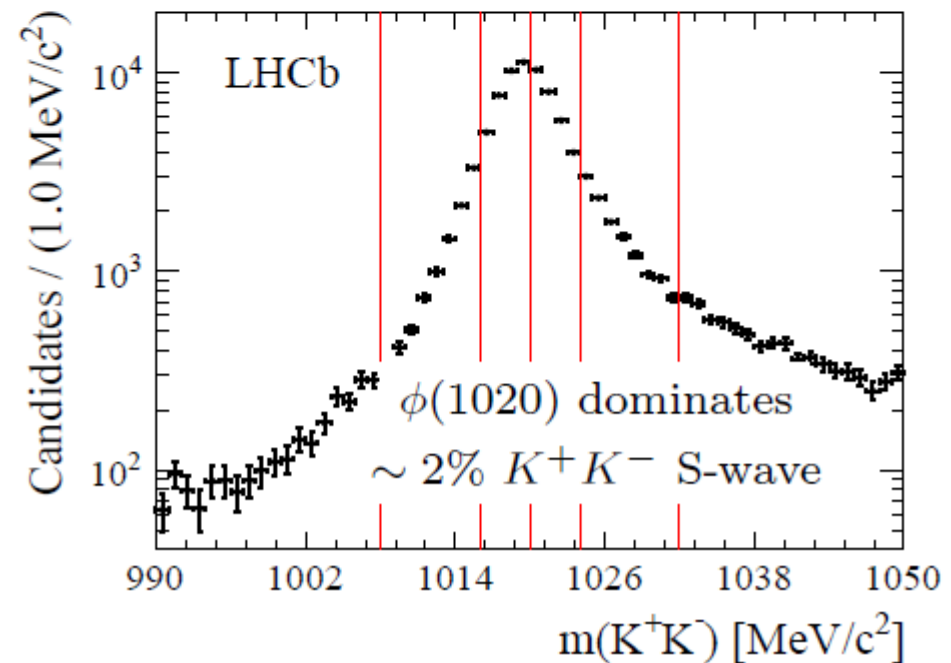
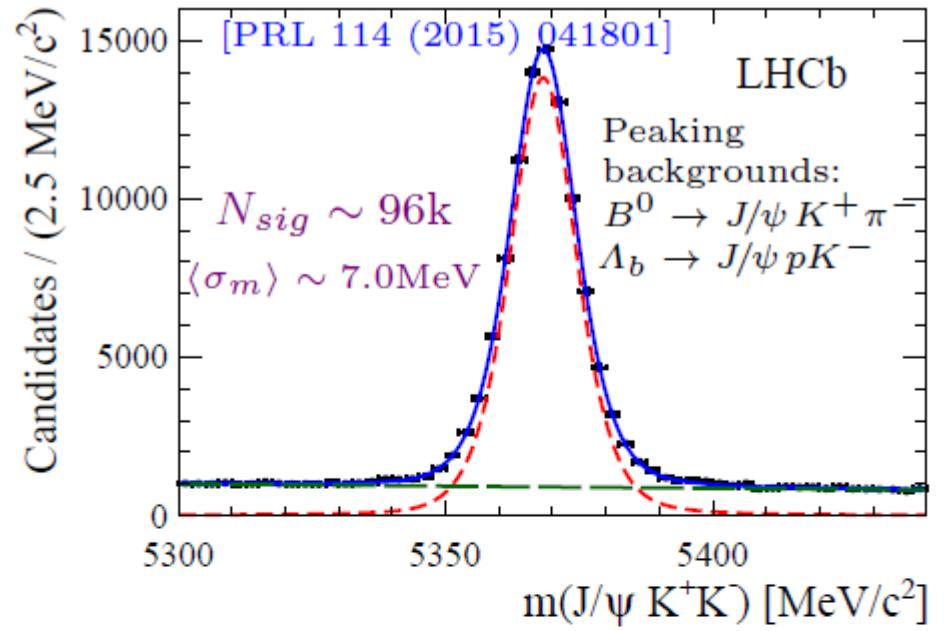
$$\mathcal{C} = -0.02 \pm 0.07(\text{stat}) \pm 0.03(\text{syst}).$$



# $\phi_s$ from $B_s^0 \rightarrow J/\psi \phi$

- $J/\psi \rightarrow \mu^+ \mu^-$
- $\Phi \rightarrow K^+ K^-$
- $B_s^0 \rightarrow J/\psi \phi$  is  $P \rightarrow VV$  decays so use angular information to disentangle  $CP$ -odd and  $CP$ -even components.
- Measure  $\phi_s, \Delta m_s, \Gamma_s, \Delta \Gamma_s, |\lambda_f|$  at the same time
- Precise determination of lifetime params

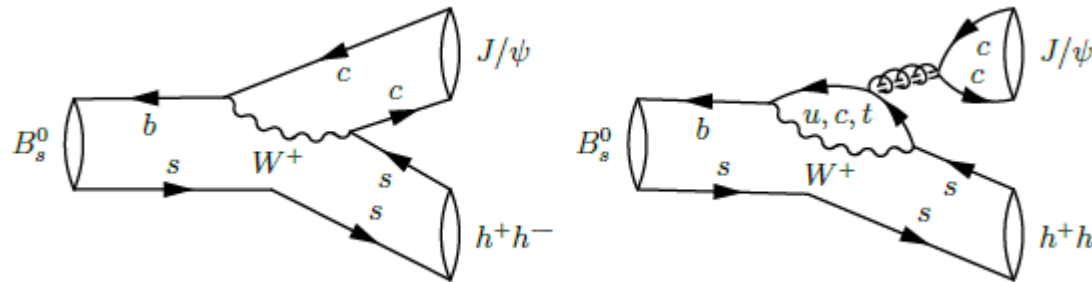
$\phi_s$	$-0.058 \pm 0.049 \pm 0.006$ rad
$ \lambda $	$0.964 \pm 0.019 \pm 0.007$
$\Gamma_s$	$0.6603 \pm 0.0027 \pm 0.0015$ ps <sup>-1</sup>
$\Delta \Gamma_s$	$0.0805 \pm 0.0091 \pm 0.0032$ ps <sup>-1</sup>
$\Delta m_s$	$17.711^{+0.055}_{-0.057} \pm 0.011$ ps <sup>-1</sup>



# Controlling penguin pollution

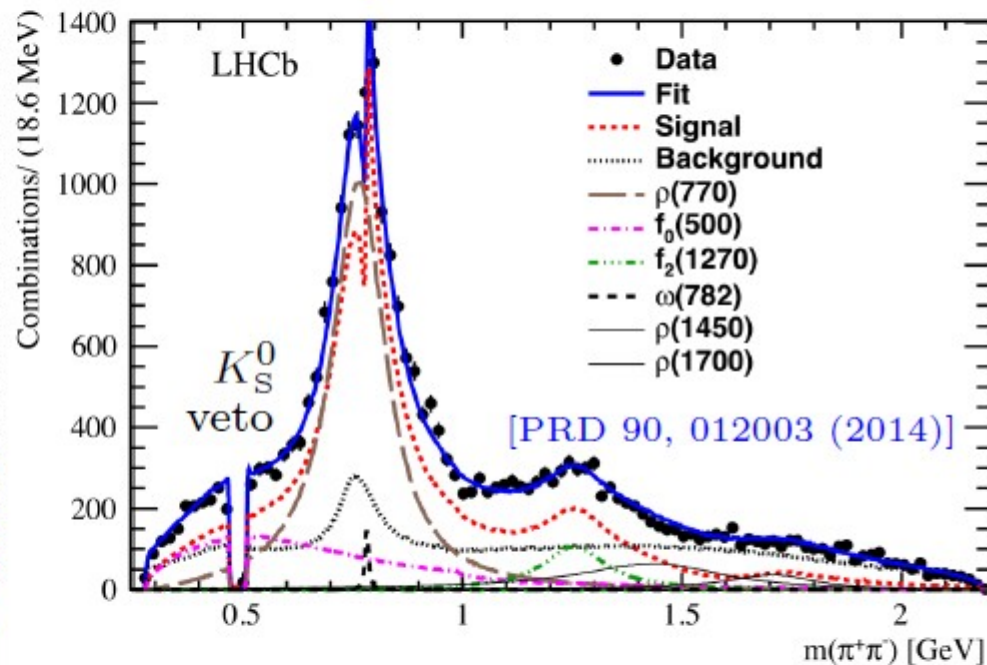
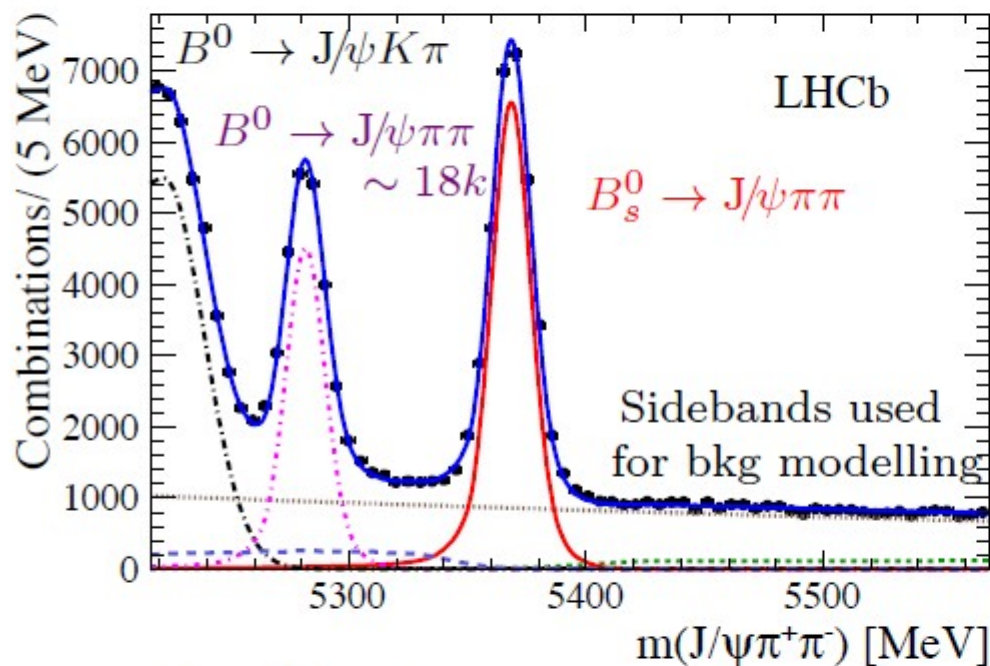
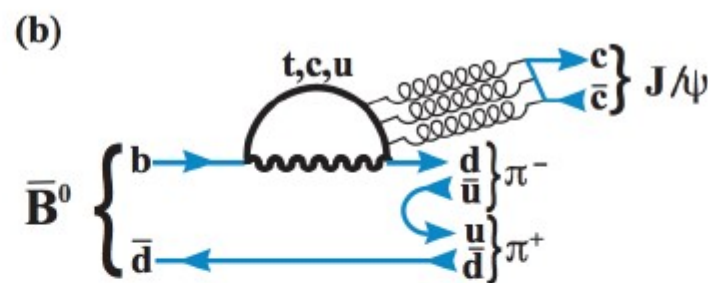
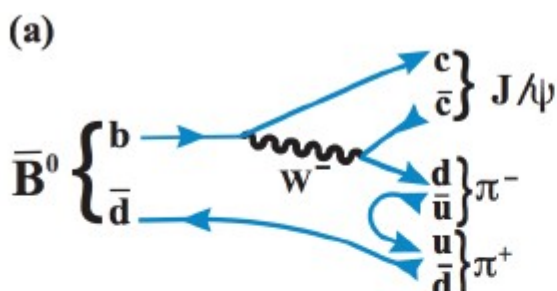
- Measured  $cp$  violating phase could be shifted due to a penguin amplitude with different weak phase

$$- \Phi_q^{\text{measured}} = \phi_q + \delta_{\text{penguin}} + \delta_{\text{NP}}$$



- $\Rightarrow$  measure using decays where penguin/tree ratio is not suppressed.
  - Use  $SU(3)$ -favour relations to link  $B_s^0$ s and  $B^0$  (broken at 20-30% level).





- Use  $\rho^0(770)$  component to measure:

$$\phi_d^{\text{eff}} = (41.7 \pm 9.6_{-6.3}^{+2.8})^\circ, \quad \alpha_{CP} \equiv \frac{1-|\lambda_f|}{1+|\lambda_f|} = (-32 \pm 28_{-7}^{+9}) \times 10^{-3}$$

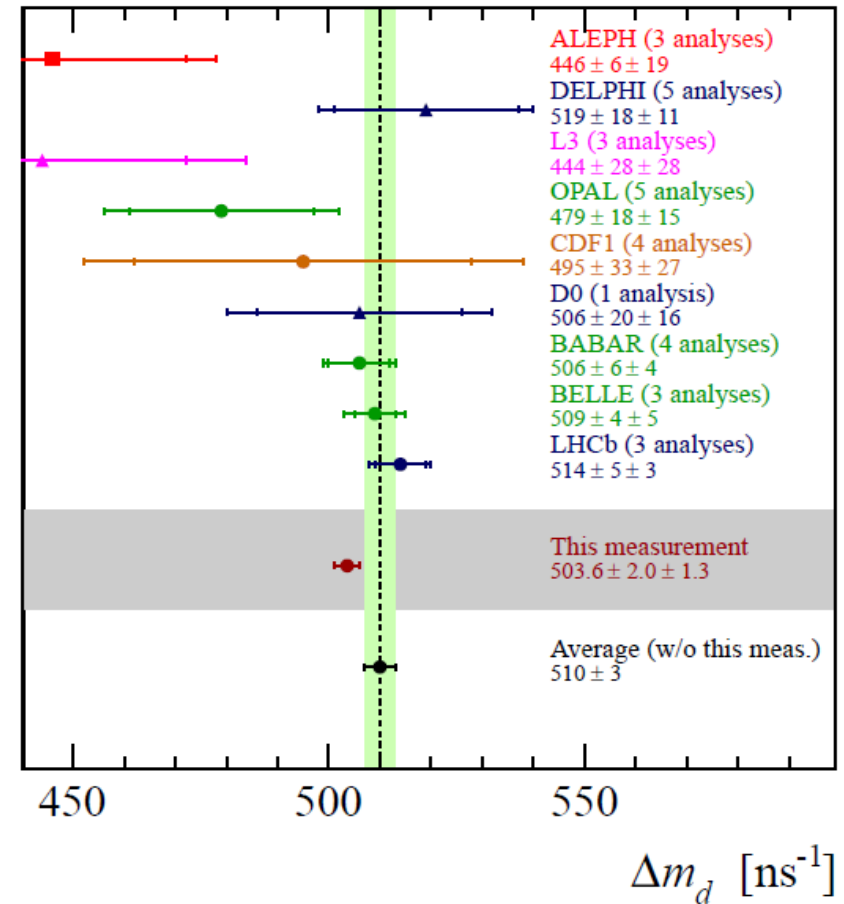
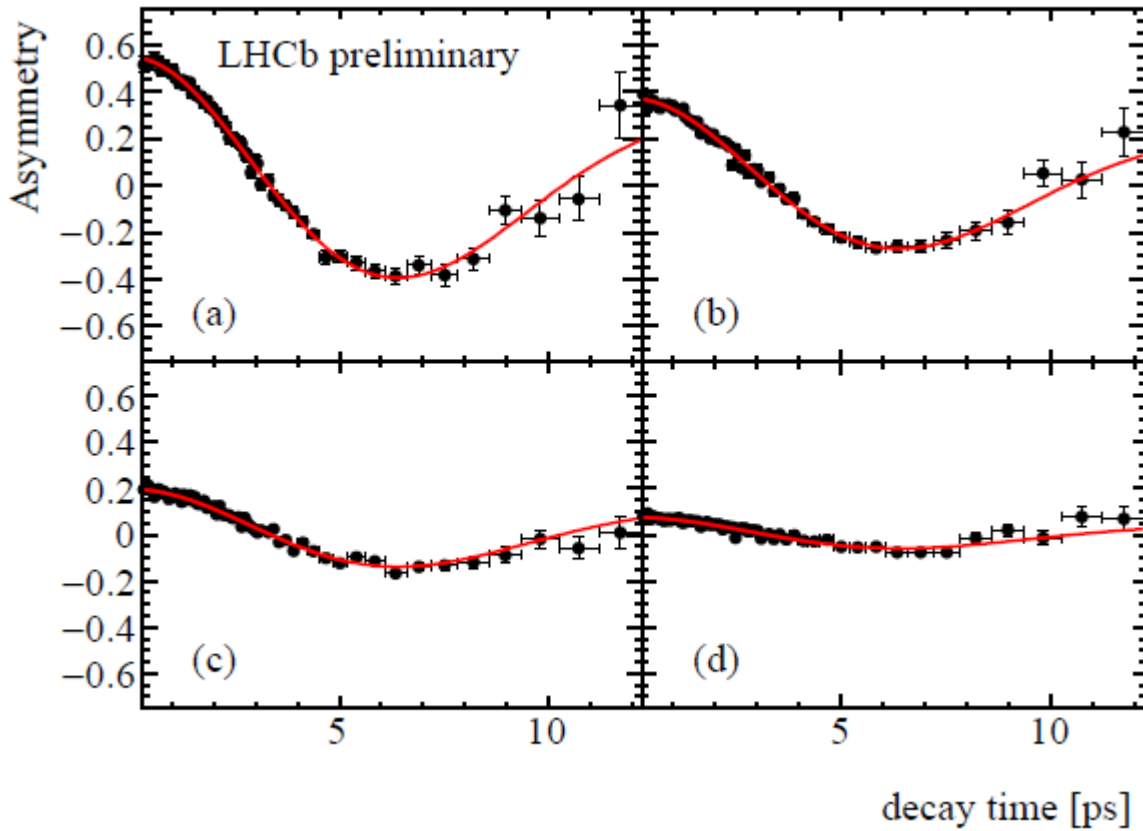
$$\Rightarrow \Delta\phi_d = (-0.9 \pm 9.7_{-6.3}^{+2.8})^\circ \quad (\text{equivalent to } 0.016 \pm 0.169_{-0.110}^{+0.049} \text{ rad})$$

Not covering  $\gamma$  and  $\alpha$



$$\Delta m_d$$





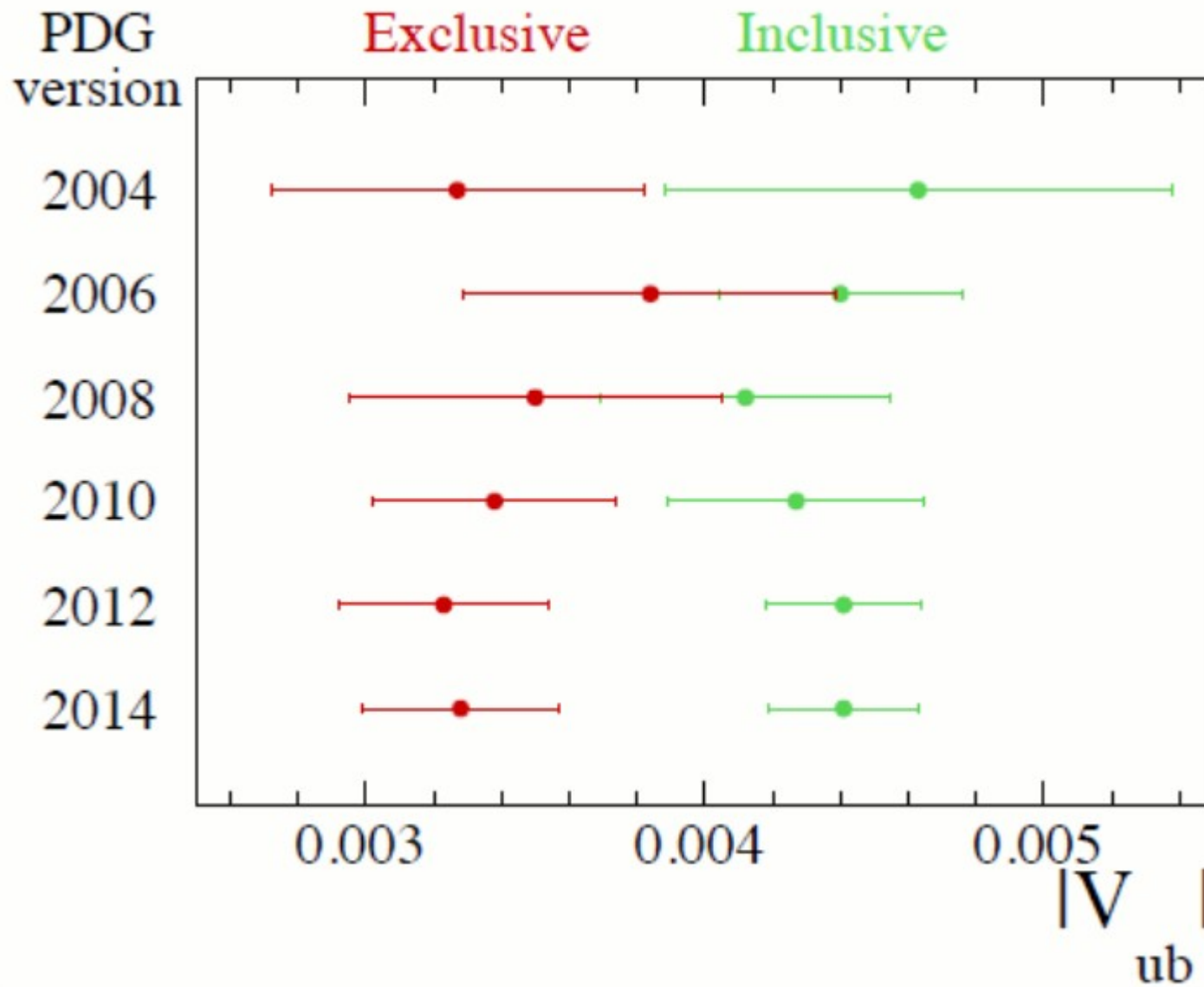
$$\Delta m_d = 503.6 \pm 2.0 (stat) \pm 1.3 (sys) \text{ GeV}/c^2$$

Most precise single measurement

$$\Lambda_b \rightarrow p \mu \nu$$

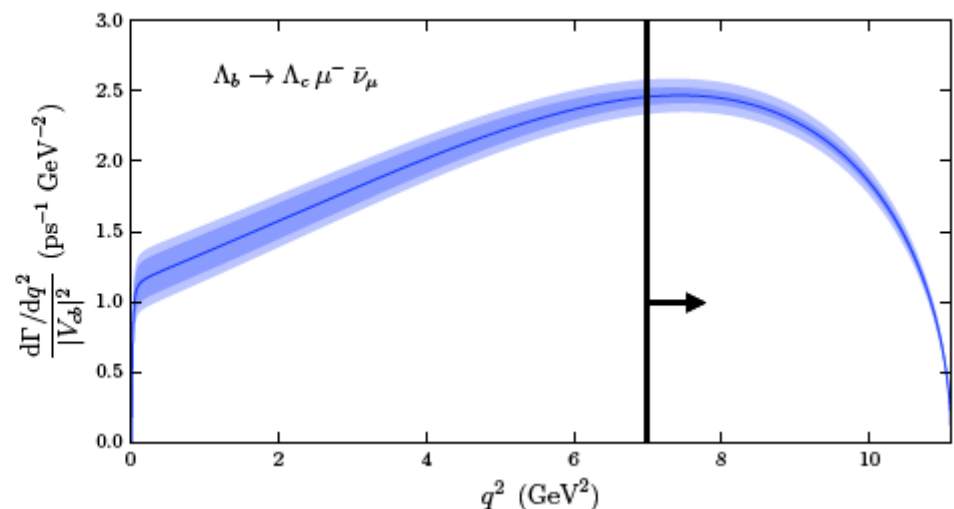
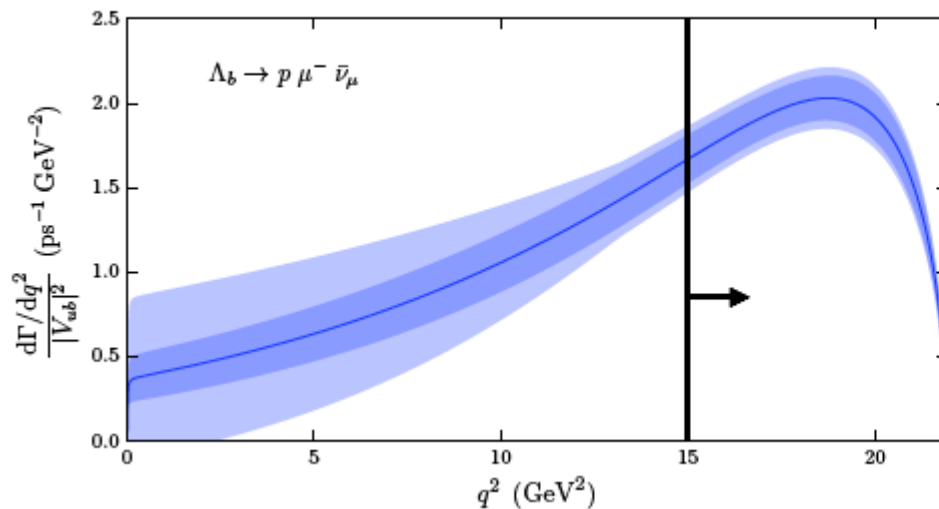
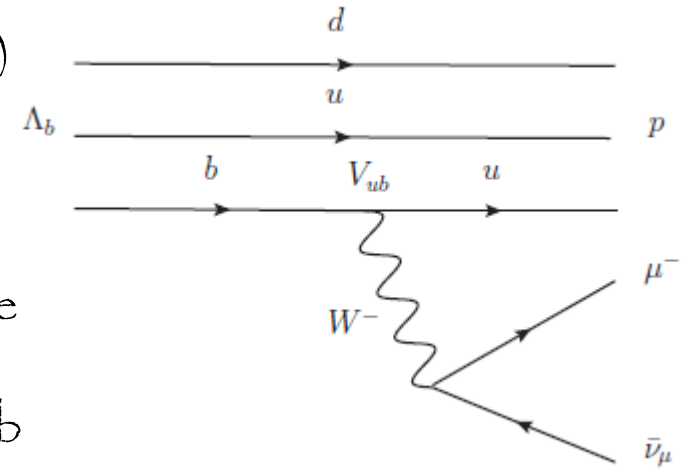
# $V_{ub}$

- Long standing issue with  $V_{ub}$

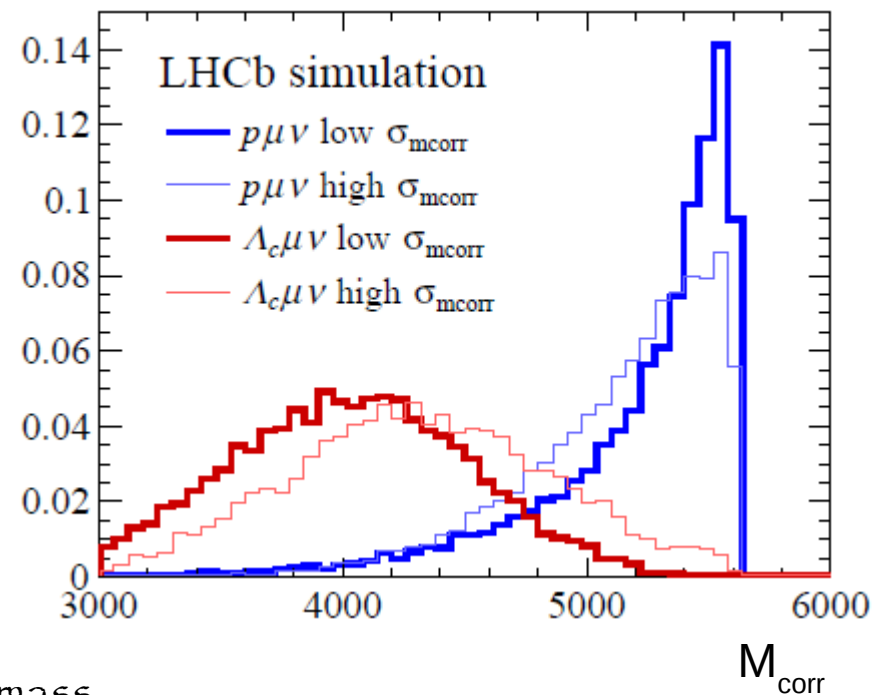
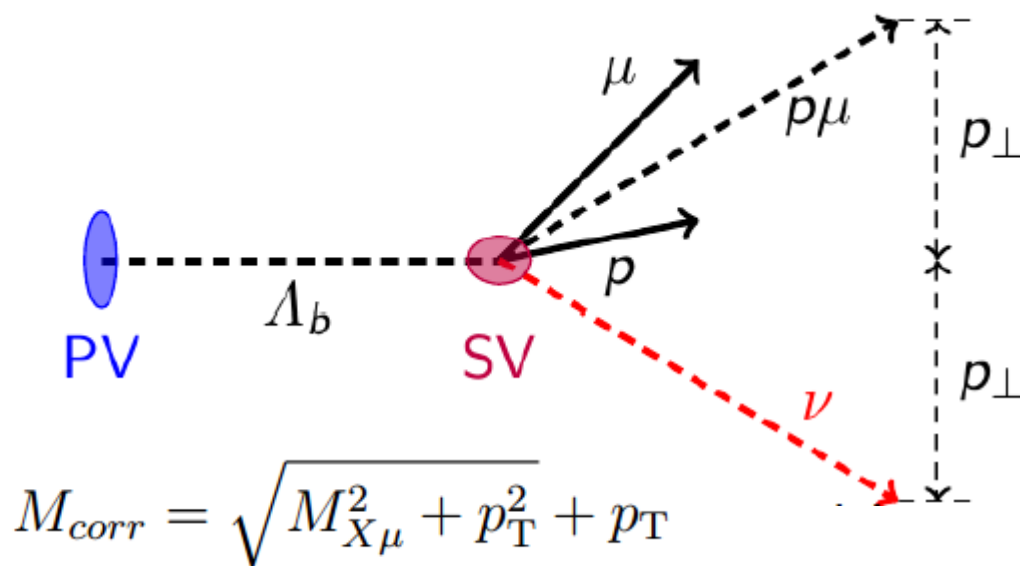


# $V_{ub}/V_{cb}$ from $\Lambda_b \rightarrow p \mu \nu$

- $V_{ub}$  measured at B-Factories using BF ( $B \rightarrow \pi \ell \nu$ )
- $\Lambda_b \rightarrow p \ell \nu$  is the baryonic version of  $B \rightarrow \pi \ell \nu$
- Measure the BF only in high  $q^2$  region where lattice calculation is known precisely  $\rightarrow$  uncertainty on  $V_{ub}$   $\sim 5\%$



W. Detmold, C. Lehner and S. Meinel: [arxiv:1503.01421\(hep-lat\)](https://arxiv.org/abs/1503.01421)

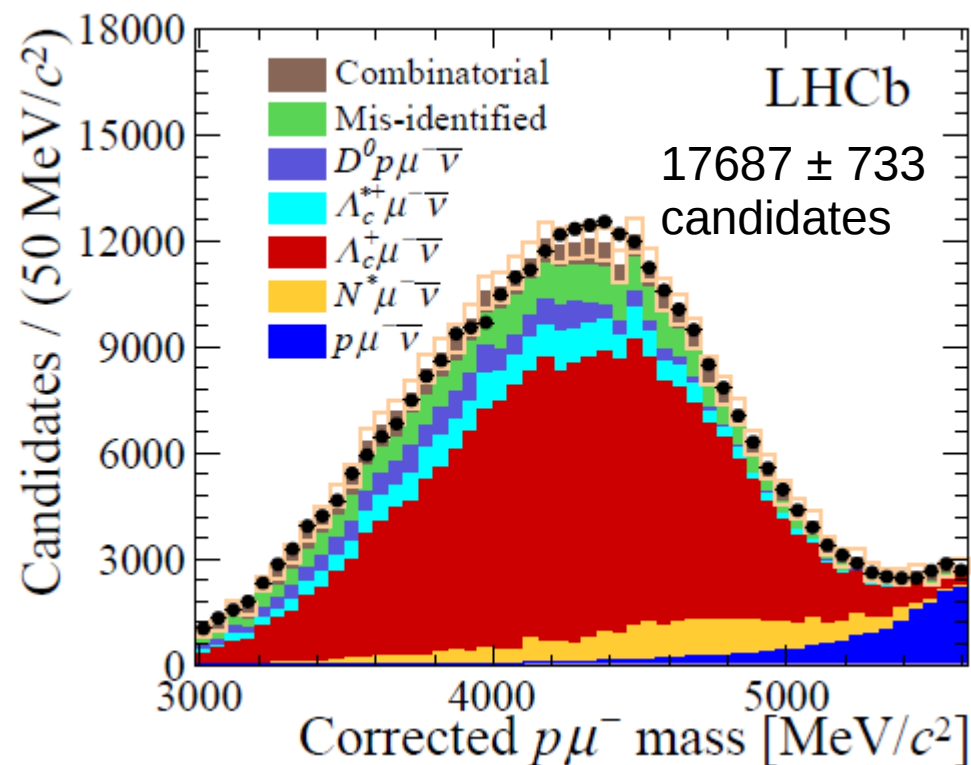
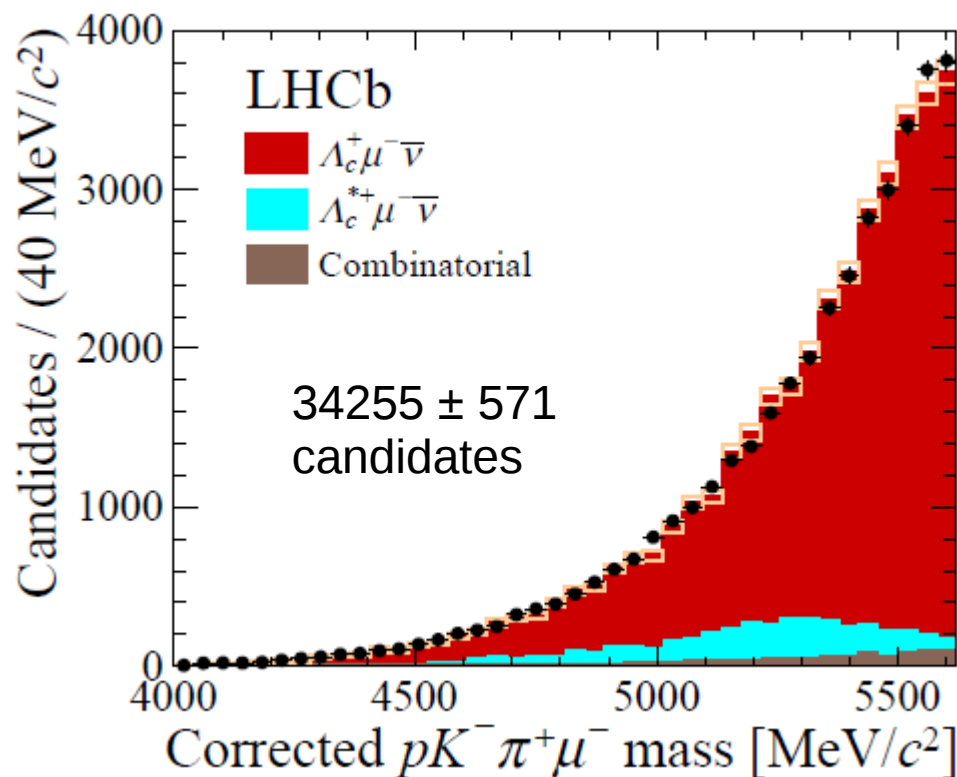


- Signal yield extracted from a fit to the corrected mass
  - Exploits the knowledge of the direction of the  $B$
  - Is the minimum mass of the system assuming a massless neutrino
  - The resolution can be improved cutting on the computer error on  $M_{corr}$
  - Selection mostly cut-based, a part from BDT used to isolate against background with additional charged tracks that were not reconstructed

Lattice calculations from  
Meinel et al.:  
[arxiv:1503.01421\(hep-lat\)](https://arxiv.org/abs/1503.01421)

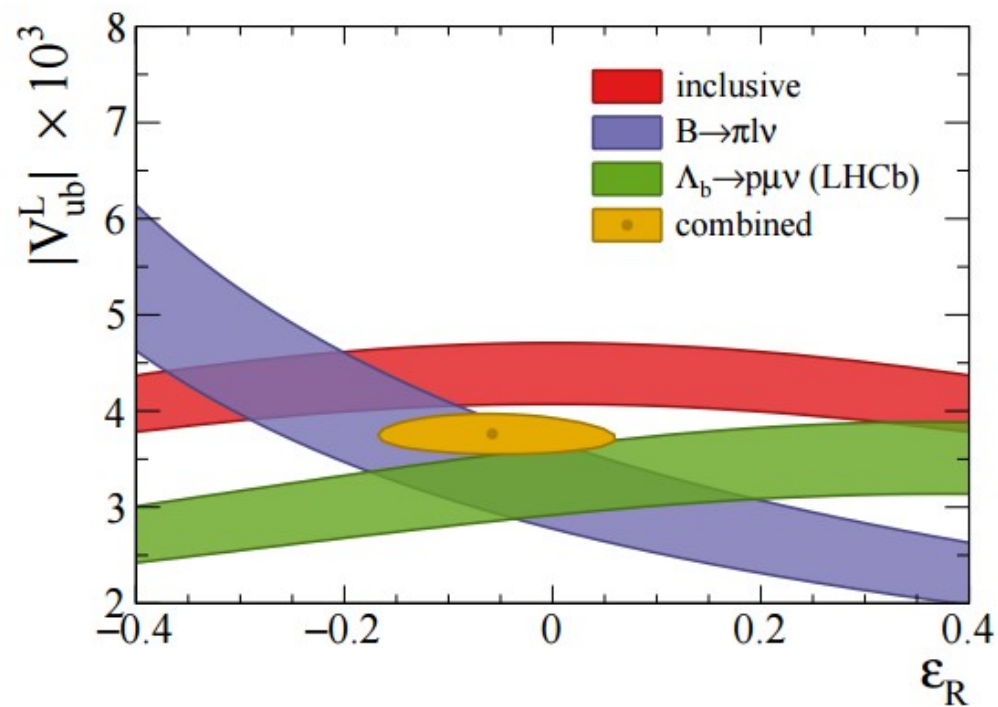
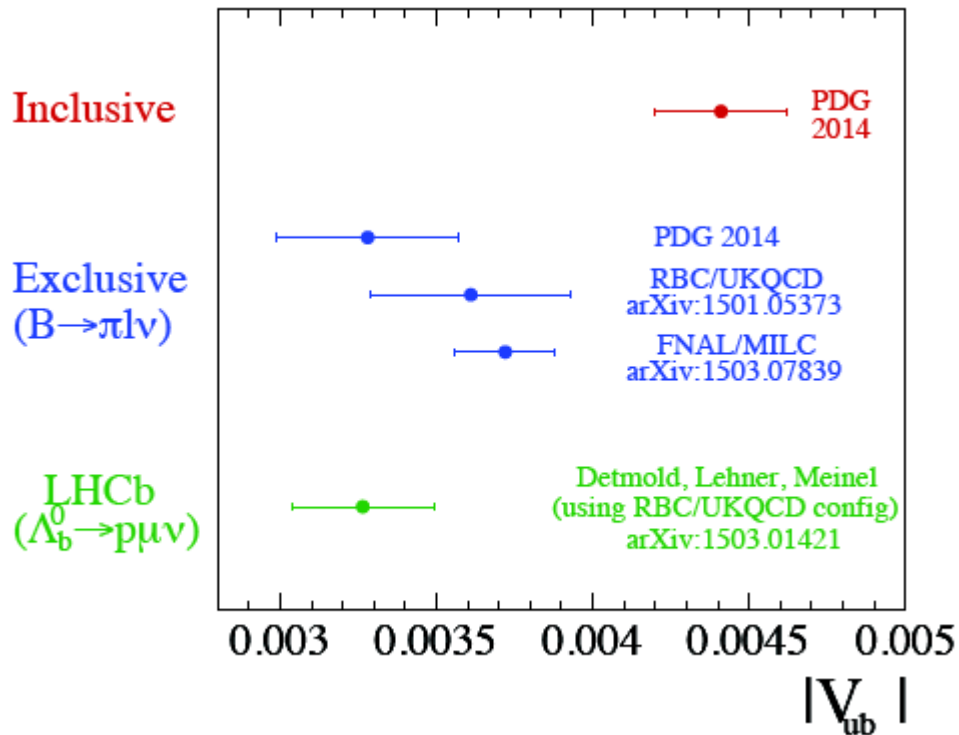
- Use  $\Lambda_b \rightarrow \Lambda_c \mu \nu$  at high  $q^2$  as normalization channel

$$\frac{\mathcal{B}(\Lambda_b^0 \rightarrow p \mu^- \bar{\nu}_\mu)_{q^2 > 15 \text{ GeV}^2/c^4}}{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda_c^+ \mu^- \bar{\nu}_\mu)_{q^2 > 7 \text{ GeV}^2/c^4}} = \frac{|V_{ub}|^2 G(\Lambda_b^0 \rightarrow p \mu^- \bar{\nu}_\mu)_{q^2 > 15 \text{ GeV}^2/c^4}}{|V_{cb}|^2 G(\Lambda_b^0 \rightarrow \Lambda_c^+ \mu^- \bar{\nu}_\mu)_{q^2 > 7 \text{ GeV}^2/c^4}}$$



- Largest background from  $\Lambda_b \rightarrow \Lambda_c \mu \nu$
- Most signal like background  $\Lambda_b \rightarrow N^* \mu \nu$





- Largest experimental systematic:  $\sim 5\%$  external measurement of  $BF(\Lambda_c \rightarrow p K \pi)$

Then trigger, tracking and selection efficiency:  $\sim 3\%$  each

- $V_{ub} = 3.27 \pm 0.15(\text{exp}) \pm 0.16(\text{theory}) \pm 0.06(|V_{cb}|)$
- Confirms discrepancy between inclusive and exclusive
- Inconsistent with a significant right handed current

# Prospects on Sides and angles

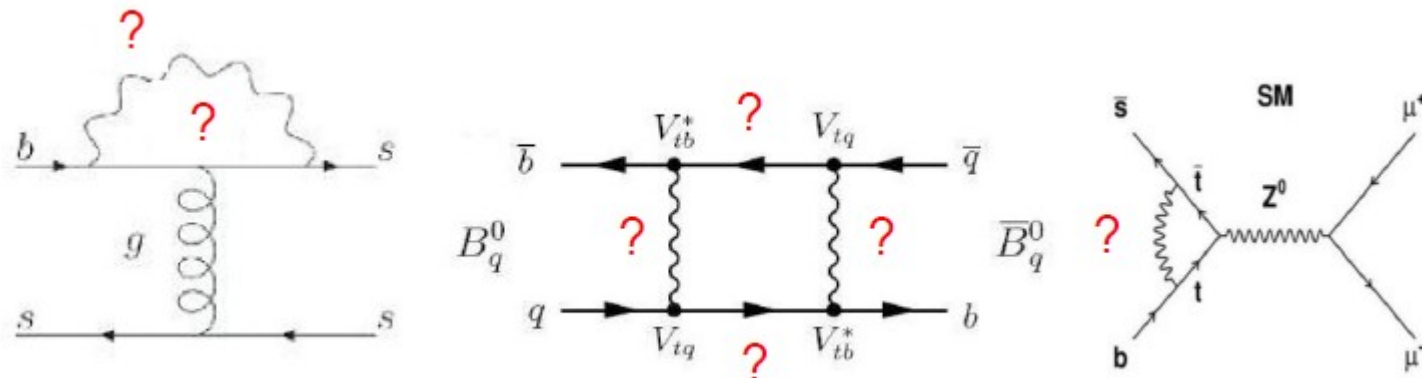
- LHCb reached comparable sensitivity to B-Factories on  $\beta$ 
  - Keep under control penguin contributions becomes important
- Both  $\beta$  and  $\gamma$  are statistically limited
  - There is a margin do better by improving the tagging
- $|V_{ub}/V_{cb}|$  can also be improved
  - Lattice calculations
  - External input  $\mathcal{B}(\Lambda_c \rightarrow p K \pi)$
  - New decays, eg  $B_s \rightarrow K \mu \nu$
- Also precision on  $\Delta m_d, \Delta m_s$  can be improved by improving lattice calculations

# Rare Decays

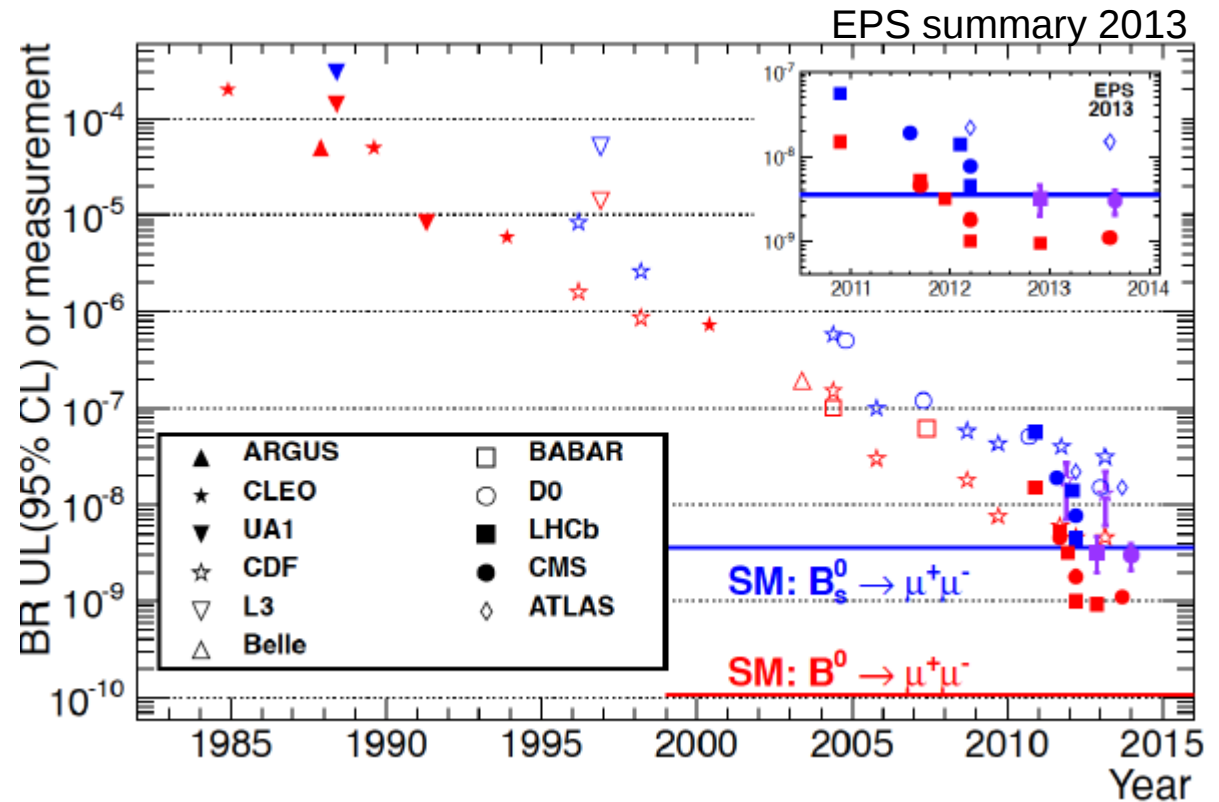
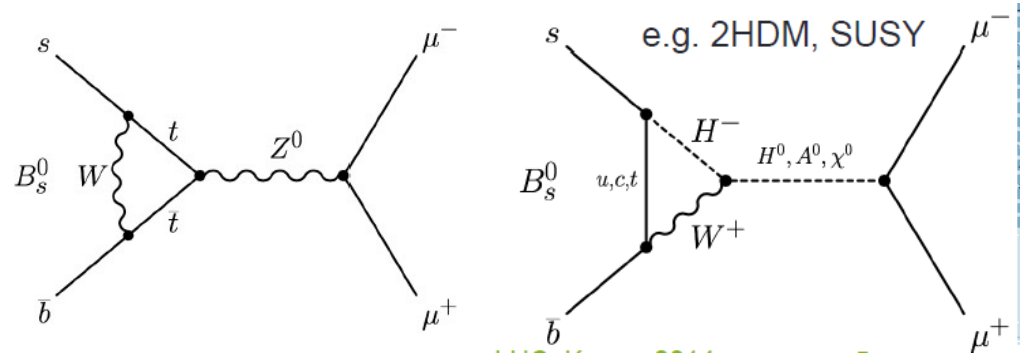
# Rare Decays

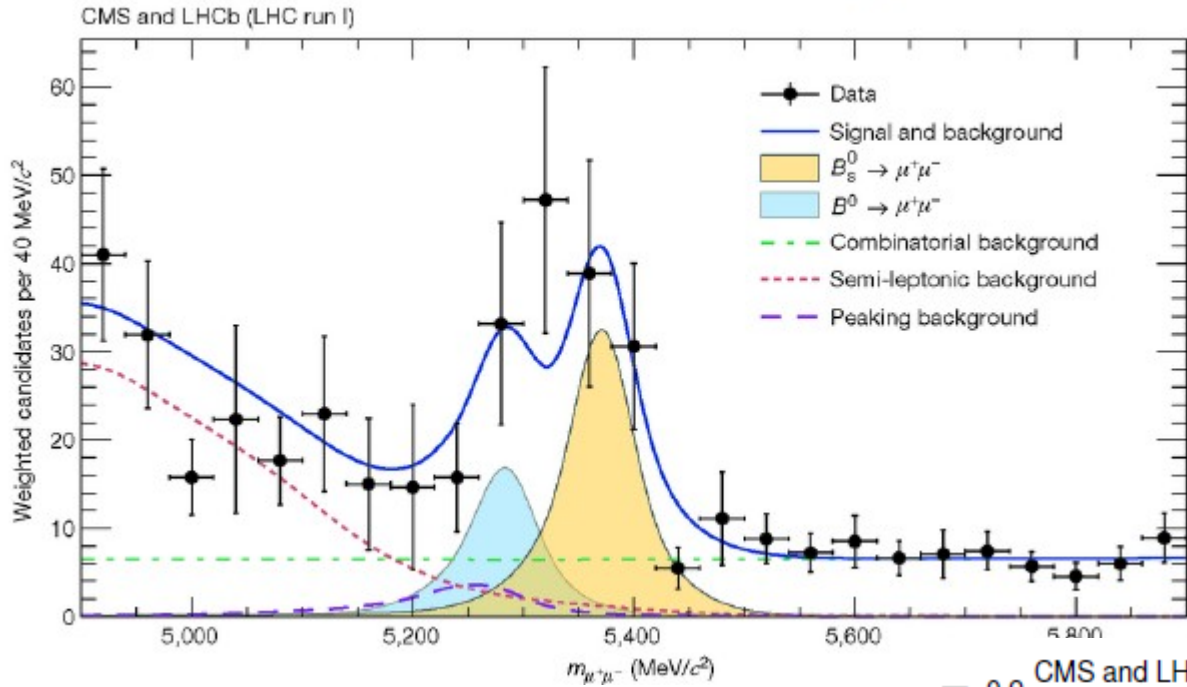
- In tree decays new physics effects are loop suppressed
- Flavor changing neutral currents are naturally highly suppressed in SM

- New intermediate particles can appear in the loops giving a contribution not suppressed with respect to the standard model => clean signatures of New Physics



- Very rare FCNC process with additional suppression due to CKM and helicity
- New physics contributions not necessarily suppressed
- Long history of search





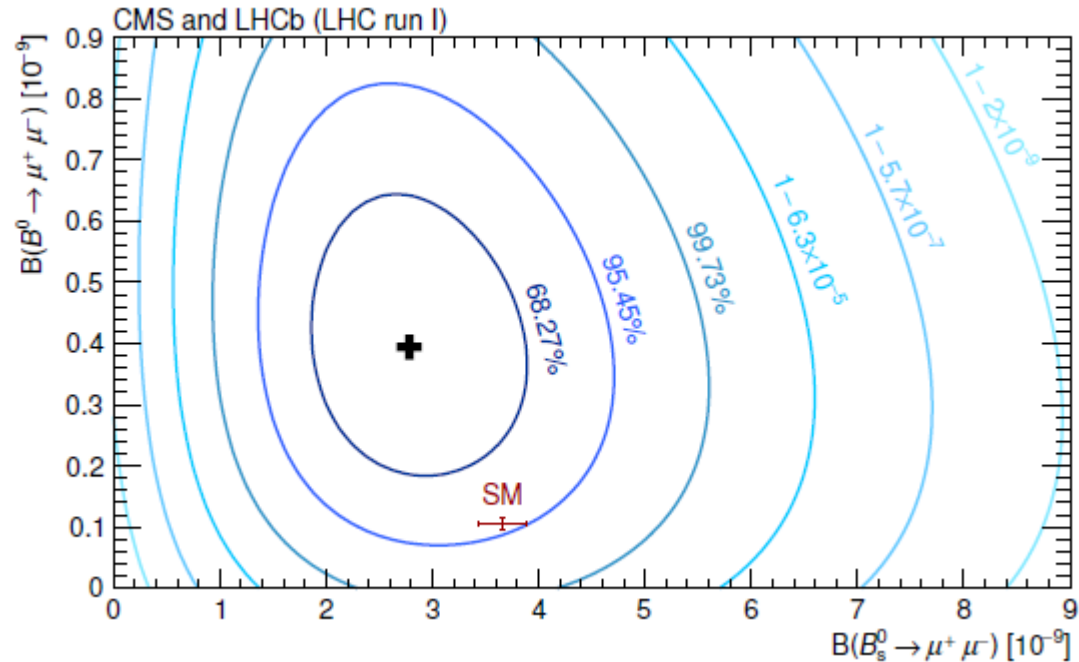
- First observation of  $B_s \rightarrow \mu\mu$  and first evidence of the  $B_d \rightarrow \mu\mu$  decay
- Compatible with the SM
- Strong constraints on new (pseudo) scalar couplings

$$\mathcal{B}(B_s^0 \rightarrow \mu^+\mu^-) = 2.8^{+0.7}_{-0.6} \times 10^{-9}$$

$$\mathcal{B}(B_s^0 \rightarrow \mu^+\mu^-)_{\text{SM}} = 3.66 \pm 0.23 \times 10^{-9}$$

$$\mathcal{B}(B_d^0 \rightarrow \mu^+\mu^-) = 3.9^{+1.6}_{-1.4} \times 10^{-10}$$

$$\mathcal{B}(B_d^0 \rightarrow \mu^+\mu^-)_{\text{SM}} = 1.06 \pm 0.09 \times 10^{-10}$$

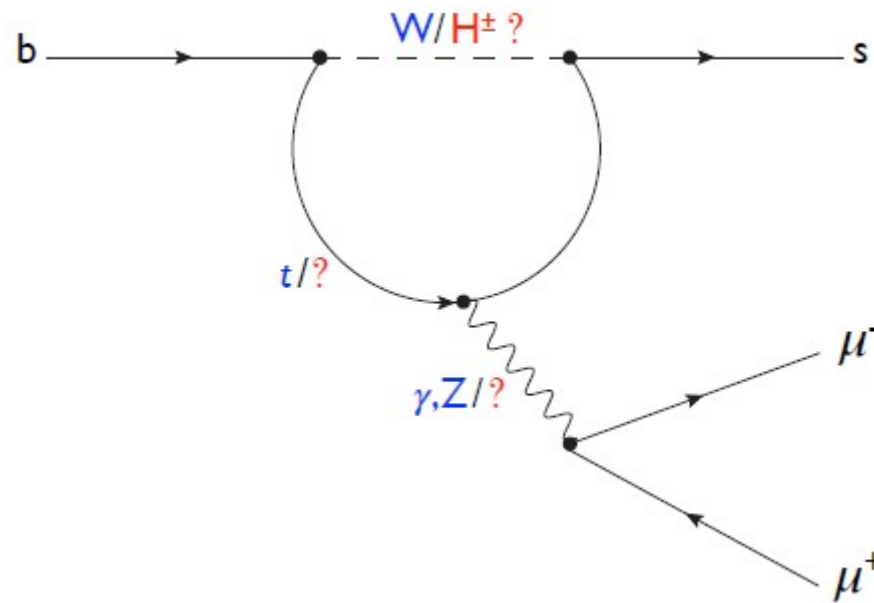




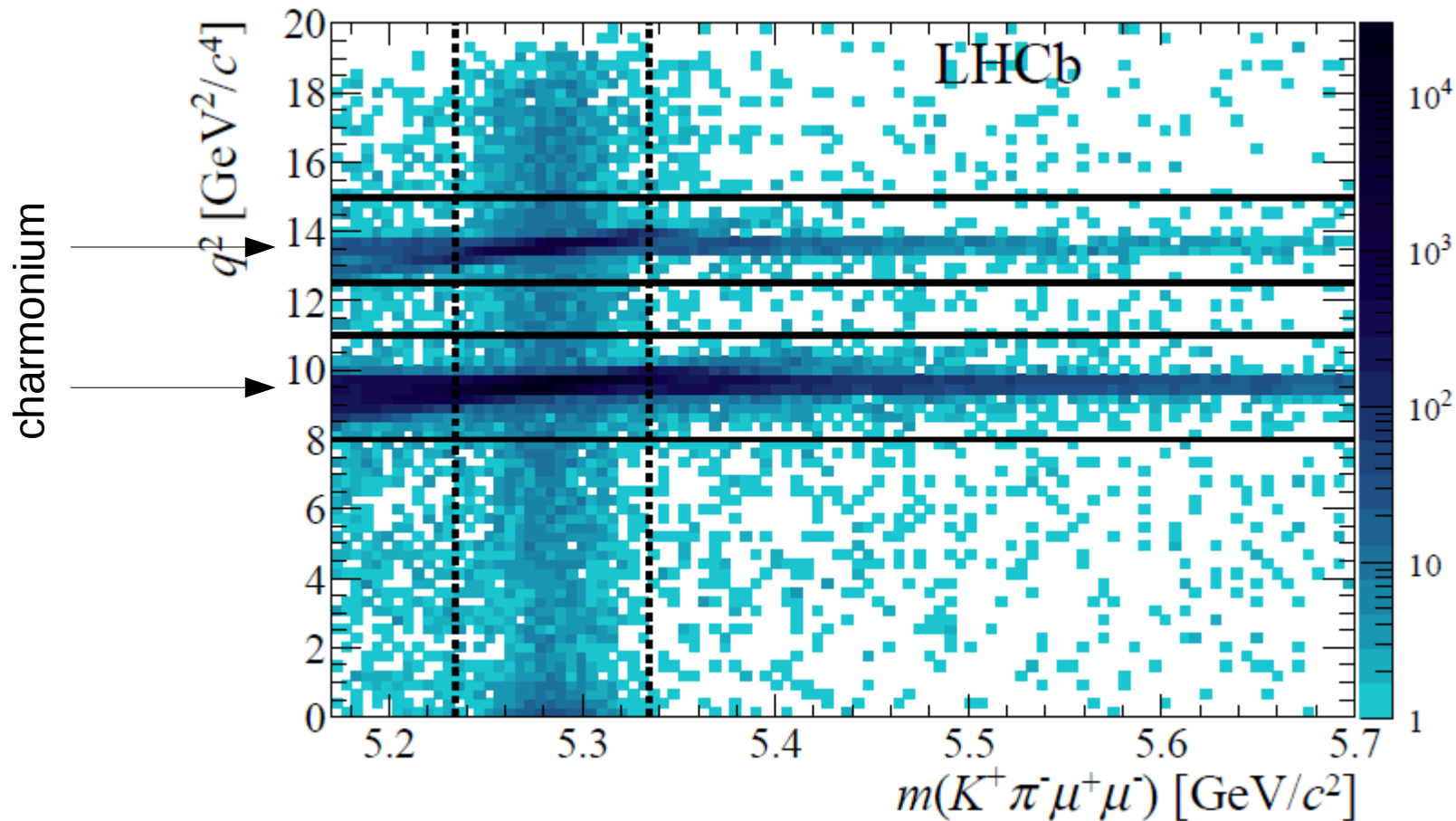
$B \rightarrow K^* \mu^+ \mu^-$ ,  $B \rightarrow \phi \mu^+ \mu^-$  full angular analysis

# $B \rightarrow K^* \mu^+ \mu^-$ EW penguin

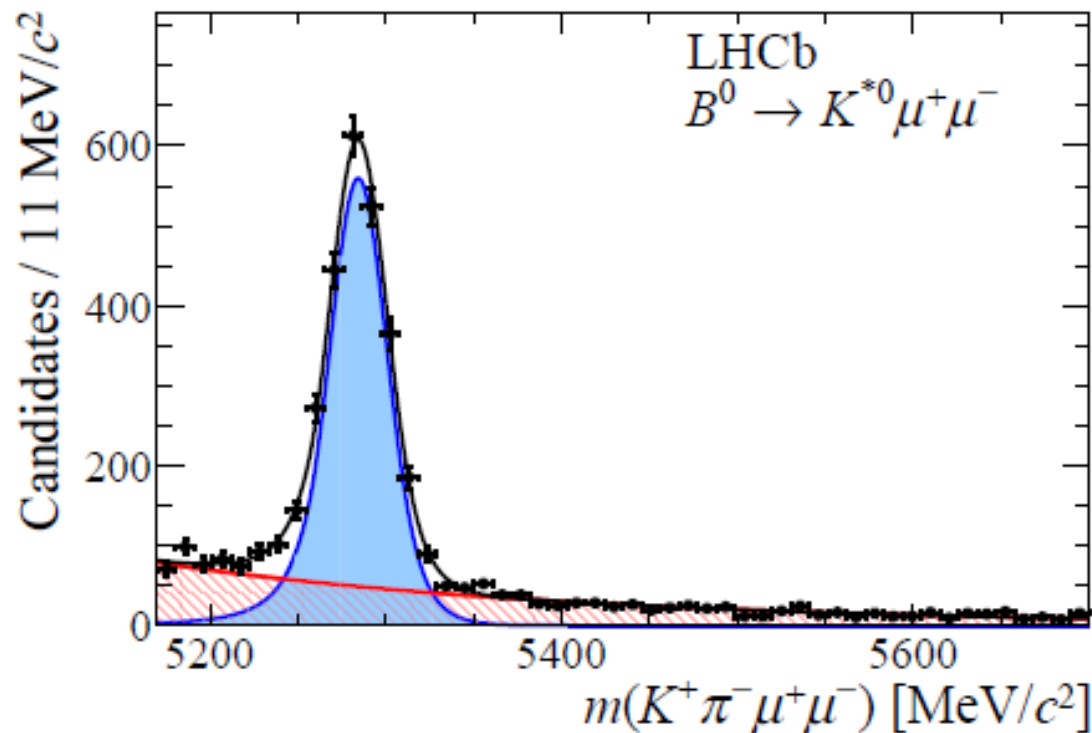
- New particles may appear in the loop and give contribution not suppressed w.r.t. SM one
- Many observables: angular distributions, decay rates, asymmetries



- Peaking background from charmonium  $\rightarrow \mu\mu$  resonances: veto using PID and kinematic variables
- Combinatorial background reduced using BDT
- $q^2 = m_{\mu\mu}^2$



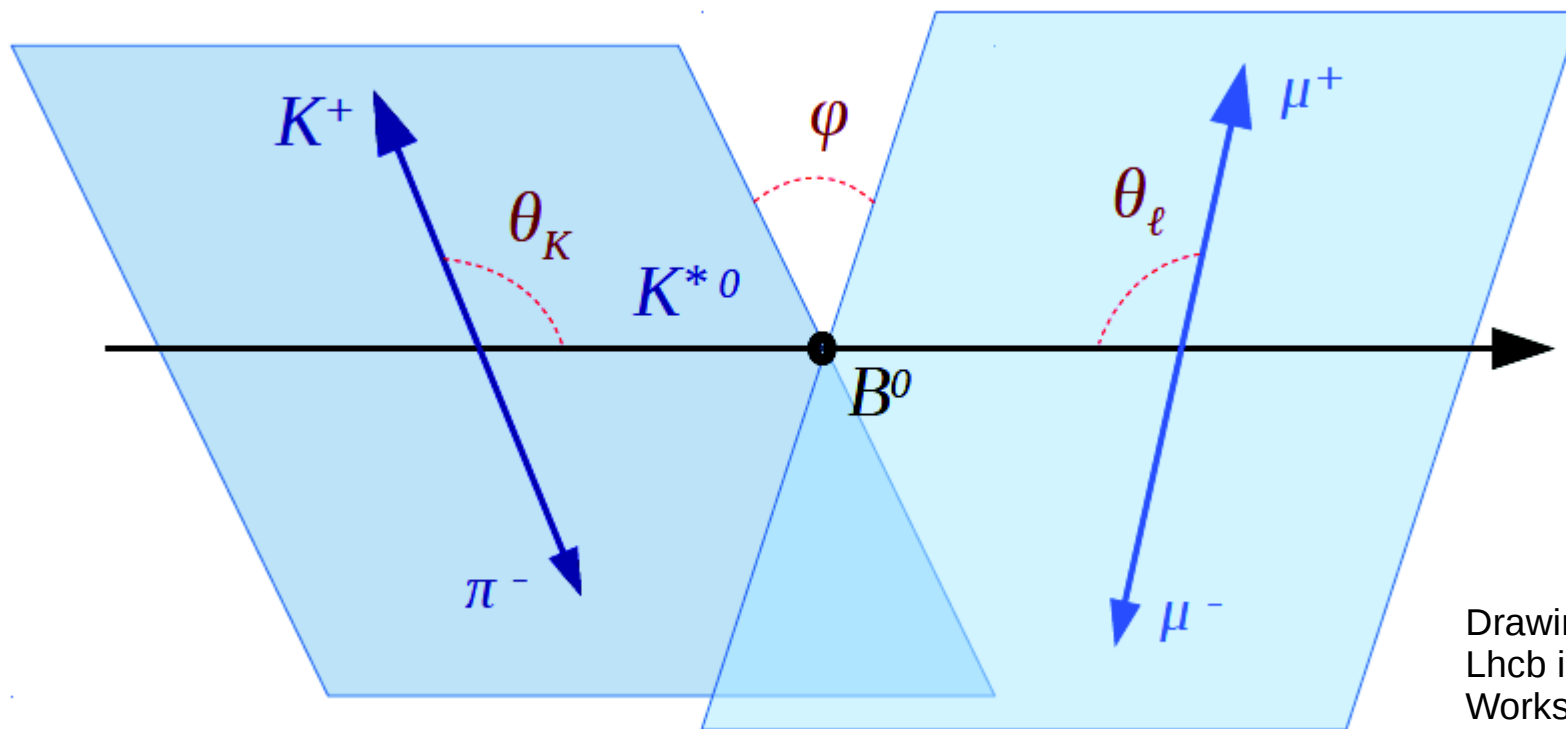
- Peaking background from charmonium  $\rightarrow \mu\mu$  resonances: veto using PID and kinematic variables
- Combinatorial background reduced using BDT
- $q^2 = m_{\mu\mu}^2$



- System fully described by  $q^2$  and three angles  $\vec{\Omega} = (\cos \theta_l, \cos \theta_k, \phi)$

$$\frac{d^4\Gamma[\bar{B}^0 \rightarrow \bar{K}^{*0} \mu^+ \mu^-]}{dq^2 d\vec{\Omega}} = \frac{9}{32\pi} \sum_j I_j(q^2) f_j(\vec{\Omega}) \quad S_j = (I_j + \bar{I}_j) / \left( \frac{d\Gamma}{dq^2} + \frac{d\bar{\Gamma}}{dq^2} \right)$$

$$\frac{d^4\bar{\Gamma}[B^0 \rightarrow K^{*0} \mu^+ \mu^-]}{dq^2 d\vec{\Omega}} = \frac{9}{32\pi} \sum_j \bar{I}_j(q^2) f_j(\vec{\Omega}) \quad A_j = (I_j - \bar{I}_j) / \left( \frac{d\Gamma}{dq^2} + \frac{d\bar{\Gamma}}{dq^2} \right)$$



Drawing by S.Cunliffe  
Lhcb implications  
Workshop 2015

- System fully described by  $q^2$  and three angles  $\vec{\Omega} = (\cos \theta_l, \cos \theta_k, \phi)$

$$\frac{d^4\Gamma[\bar{B}^0 \rightarrow \bar{K}^{*0} \mu^+ \mu^-]}{dq^2 d\vec{\Omega}} = \frac{9}{32\pi} \sum_j I_j(q^2) f_j(\vec{\Omega}) \quad S_j = (I_j + \bar{I}_j) / \left( \frac{d\Gamma}{dq^2} + \frac{d\bar{\Gamma}}{dq^2} \right)$$

$$\frac{d^4\bar{\Gamma}[B^0 \rightarrow K^{*0} \mu^+ \mu^-]}{dq^2 d\vec{\Omega}} = \frac{9}{32\pi} \sum_j \bar{I}_j(q^2) f_j(\vec{\Omega}) \quad A_j = (I_j - \bar{I}_j) / \left( \frac{d\Gamma}{dq^2} + \frac{d\bar{\Gamma}}{dq^2} \right)$$

- Full set of  $S_j$  and  $A_j$  never measured before

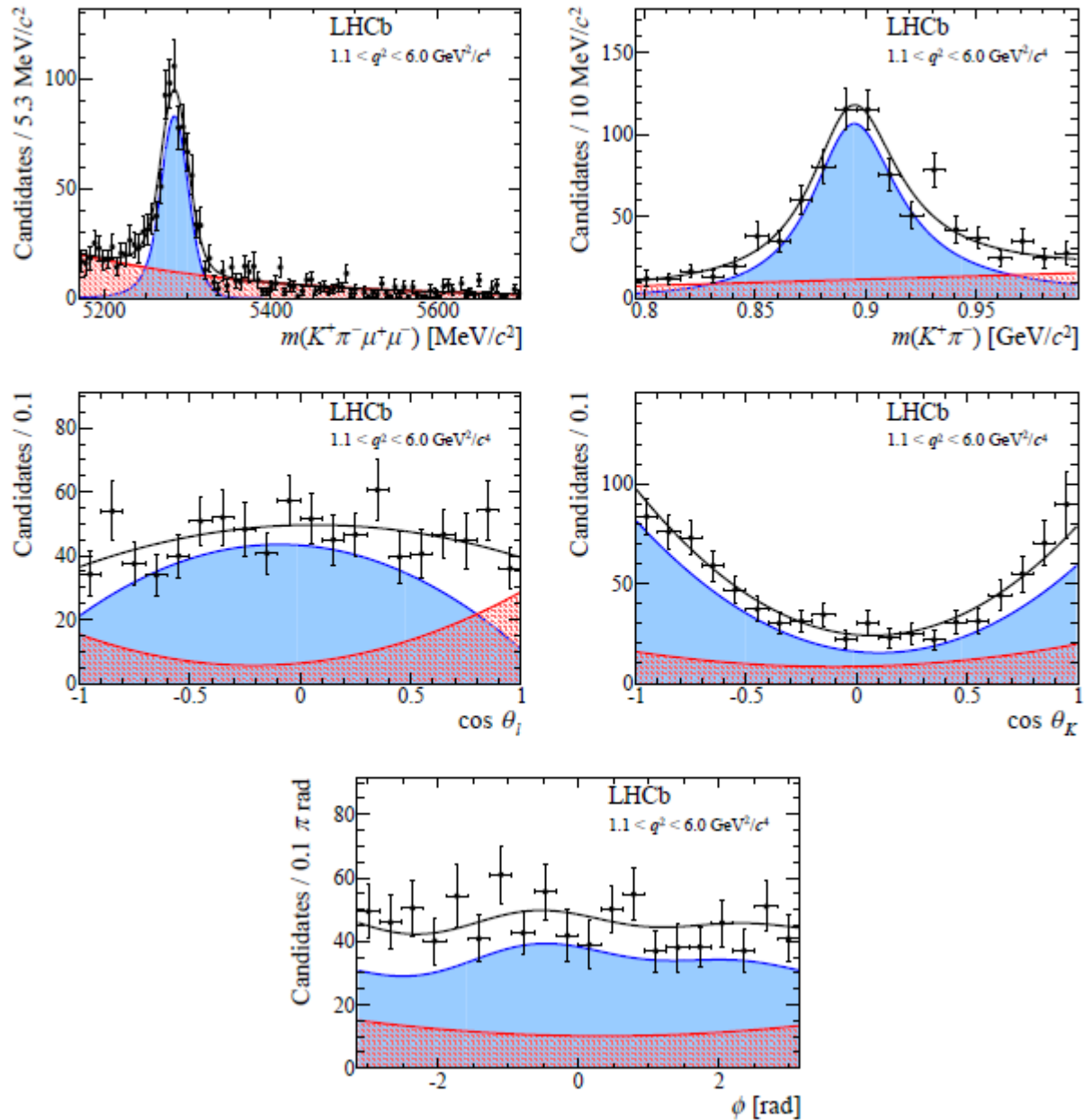


- Clever combinations have reduced dependence on form factors

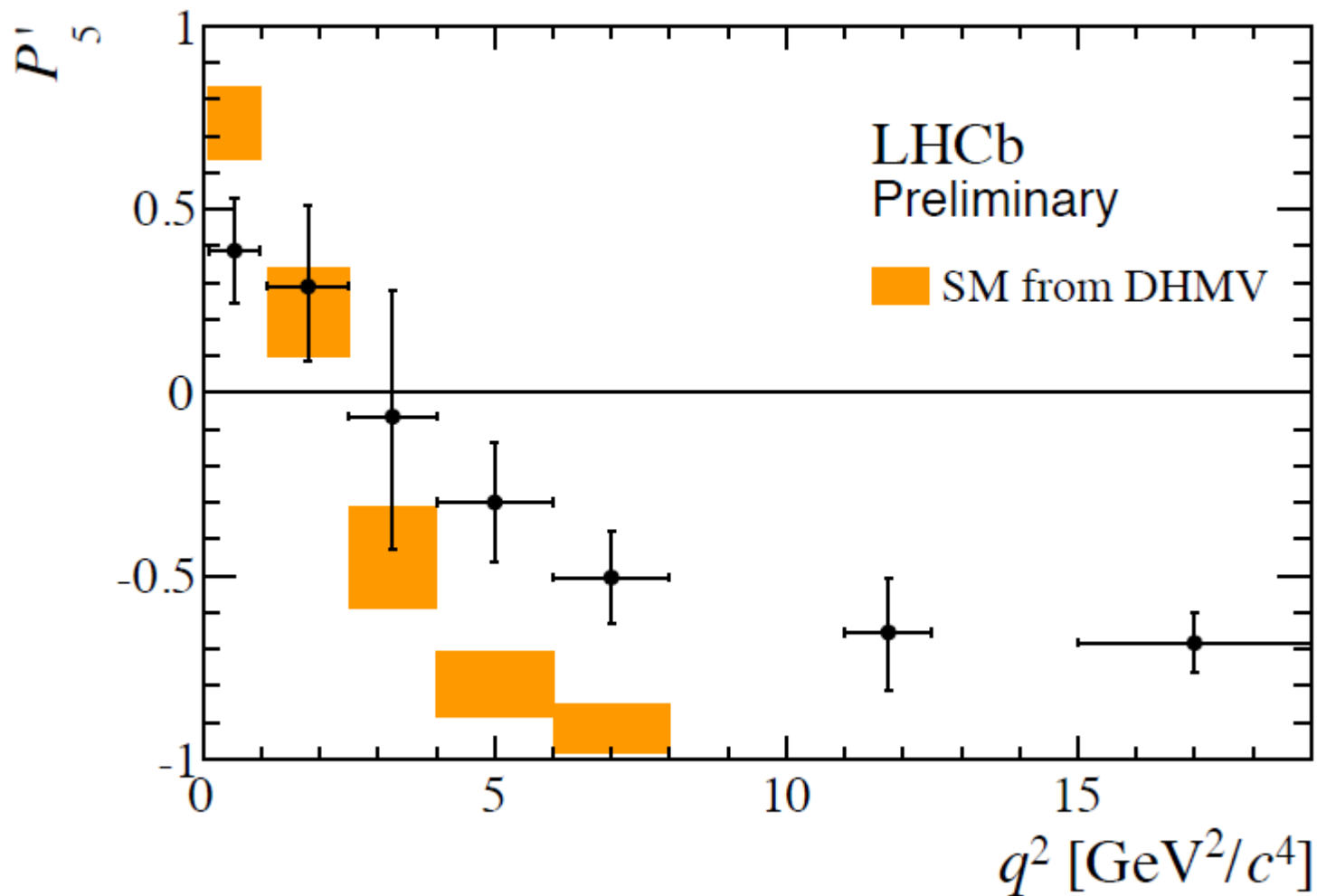
$$P'_{4,5} = S_{4,5} / \sqrt{F_L(1 - F_L)}$$

- Account for  $S$ -wave  $K\pi$  component which can help separate vector from scalar by fitting simultaneously the  $m_{K\pi}$
- Correct for dependence of efficiency as a function of angles

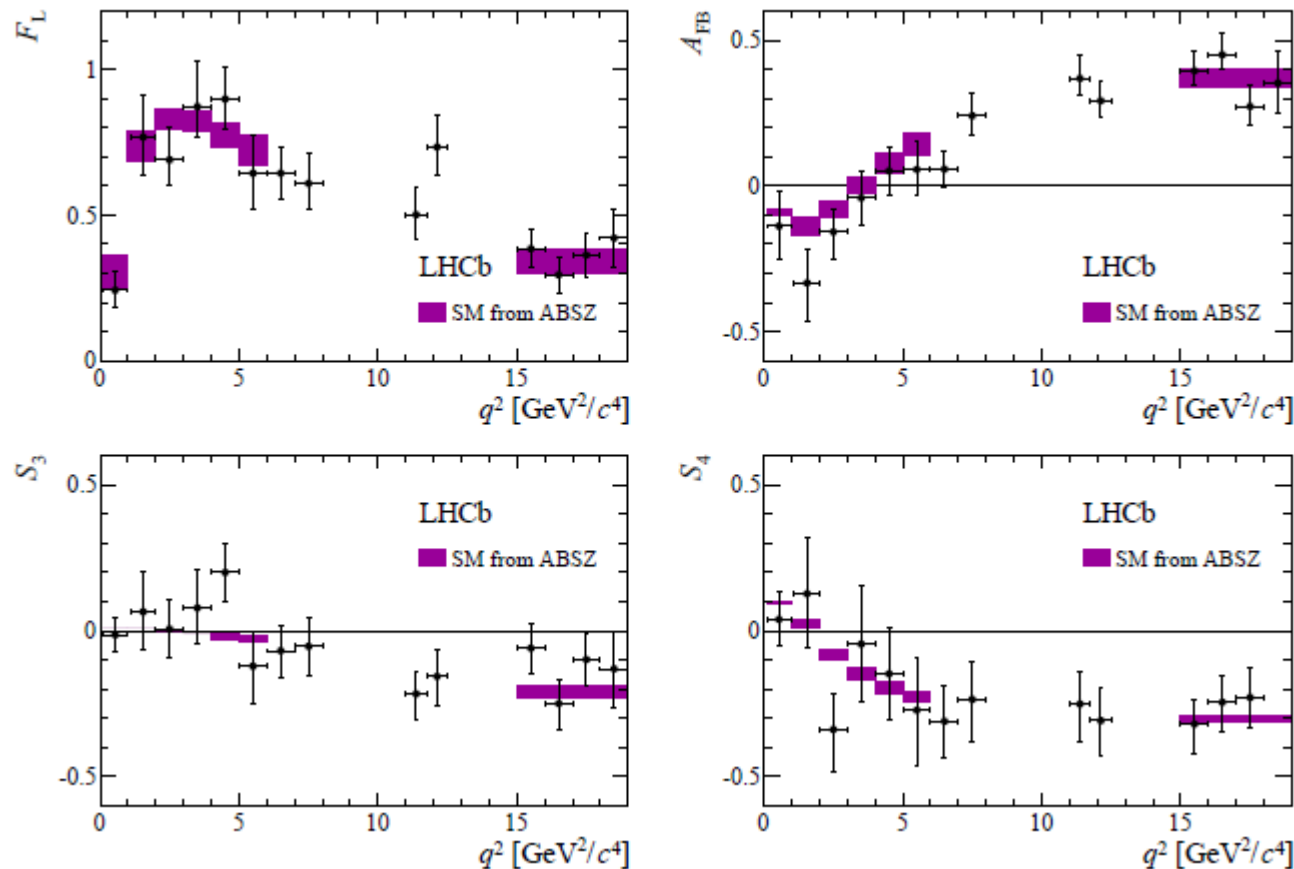




- Inconsistency with SM remains, the level of the inconsistency evaluated using the complete set of observables

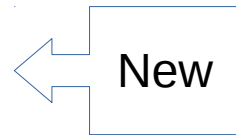


- Two other methods used to fit the parameters of the angular distributions
- Moments
  - Compute spherical moments of angular distribution and extract coefficients of spherical harmonics



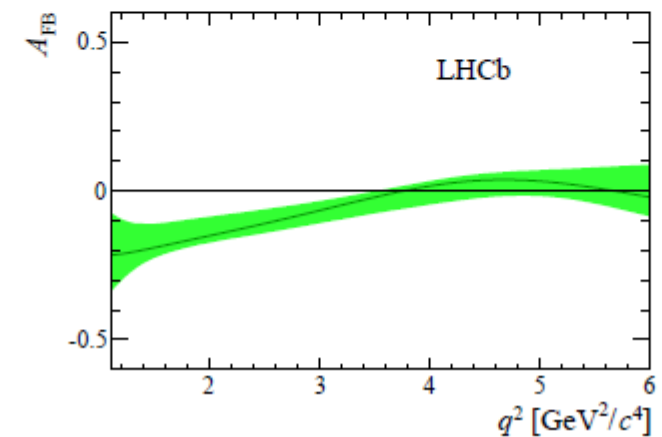
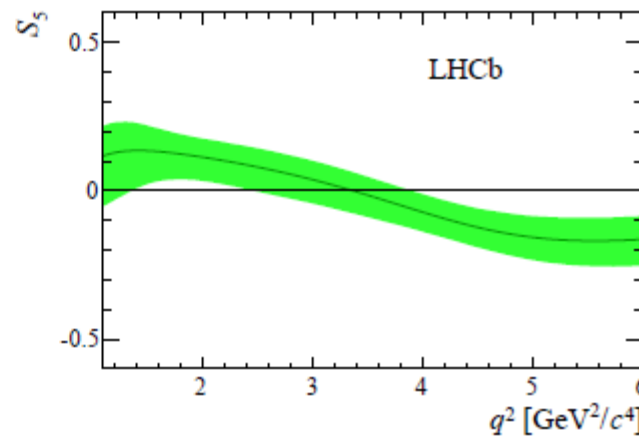
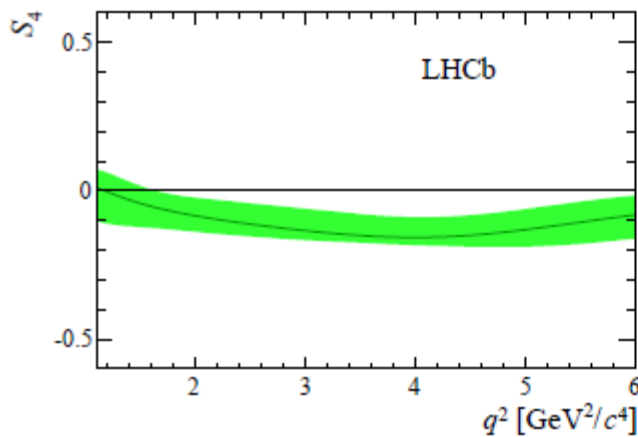
- Two other methods used to fit the parameters of the angular distributions

- Amplitudes



- Fit directly the amplitudes making the assumption

$$A(q^2) = \alpha + \beta q^2 + \frac{\gamma}{q^2}$$



- Extract directly the zero crossing points of the observables

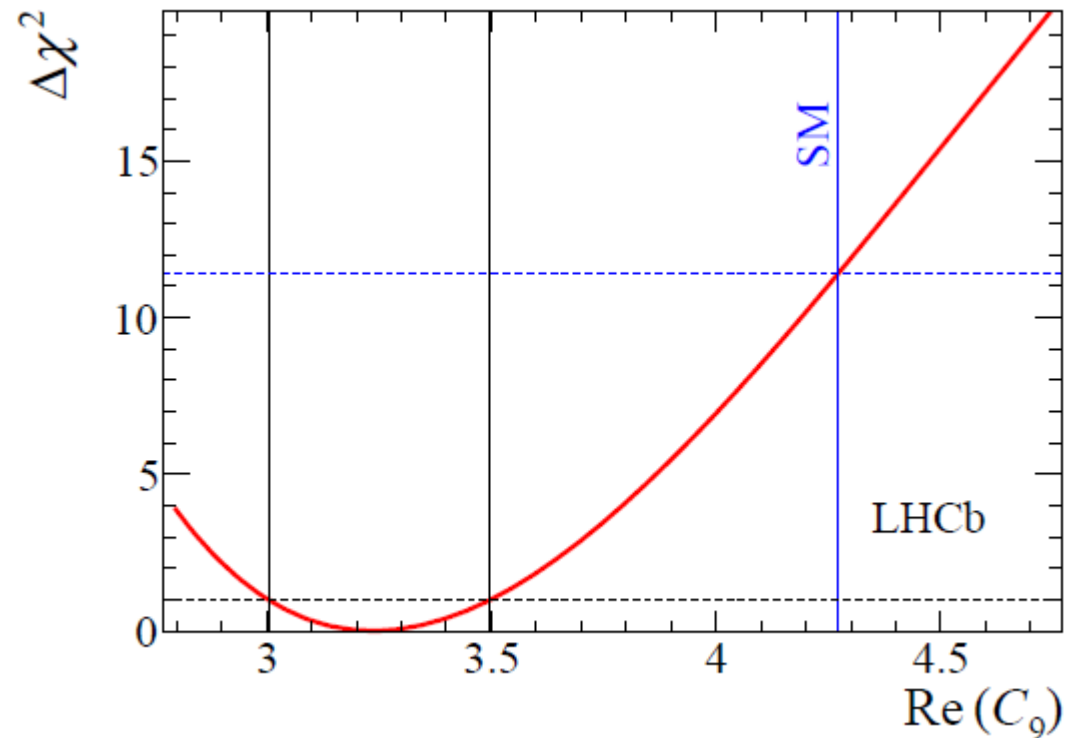
$$q_0^2(S_4) < 2.65 \text{ GeV}^2/c^4 \text{ at 95\% confidence level (C.L.) ,}$$

$$q_0^2(S_5) \in [2.49, 3.95] \text{ GeV}^2/c^4 \text{ at 68\% C.L. ,}$$

$$q_0^2(A_{\text{FB}}) \in [3.40, 4.87] \text{ GeV}^2/c^4 \text{ at 68\% C.L. .}$$

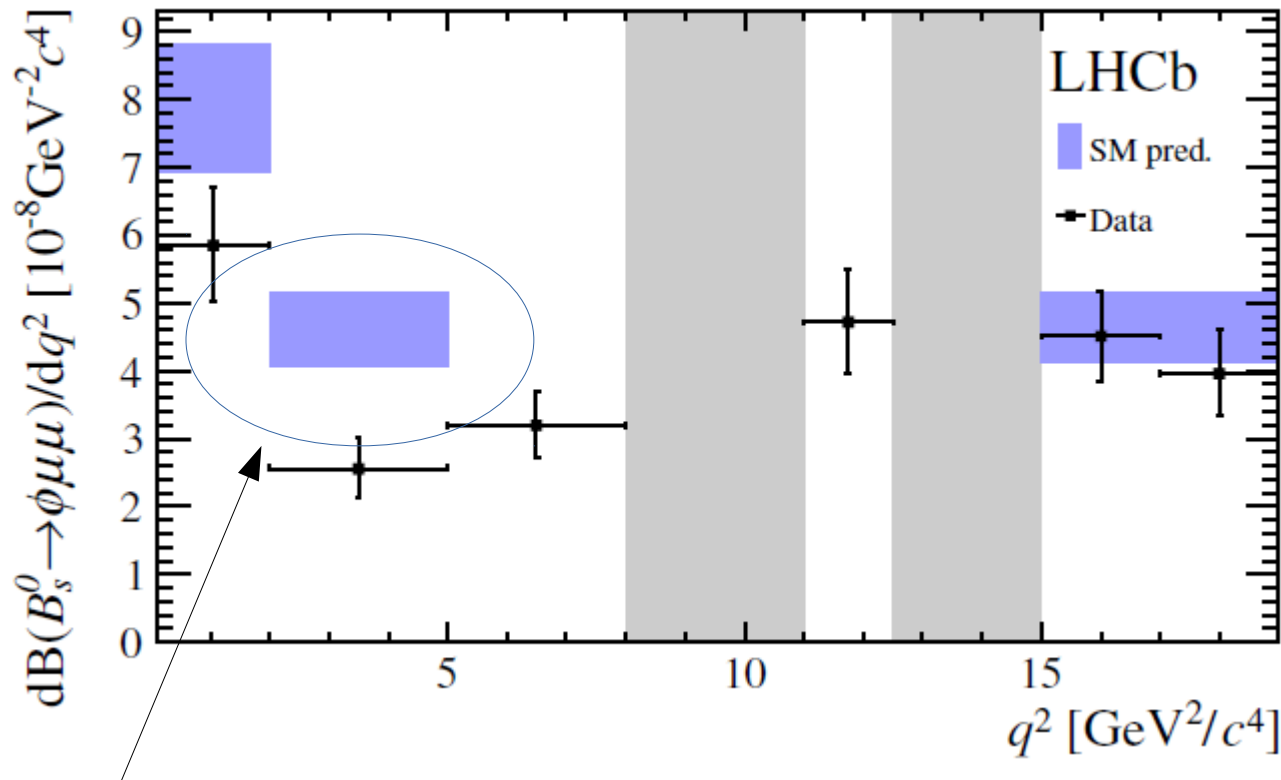


- A  $\chi^2$  fit for  $\text{Re}(C_9)$  is performed to the CP-averaged angular observables  $F_L$ ,  $A_{\text{FB}}$  and  $S_3$ - $S_9$  obtained from the likelihood fit to the data
- No additional effective couplings tested
- Compatibility with SM is  $3.4 \sigma$
- A fit to the moments gives consistent results but less precise



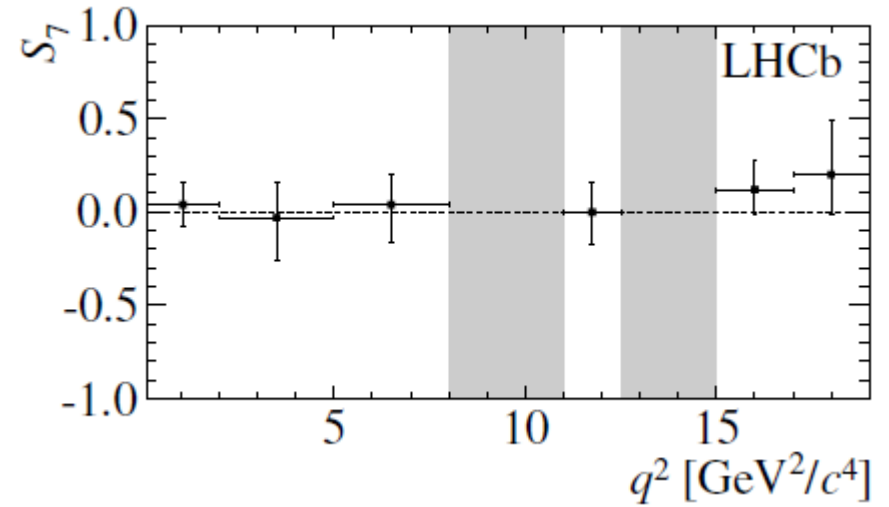
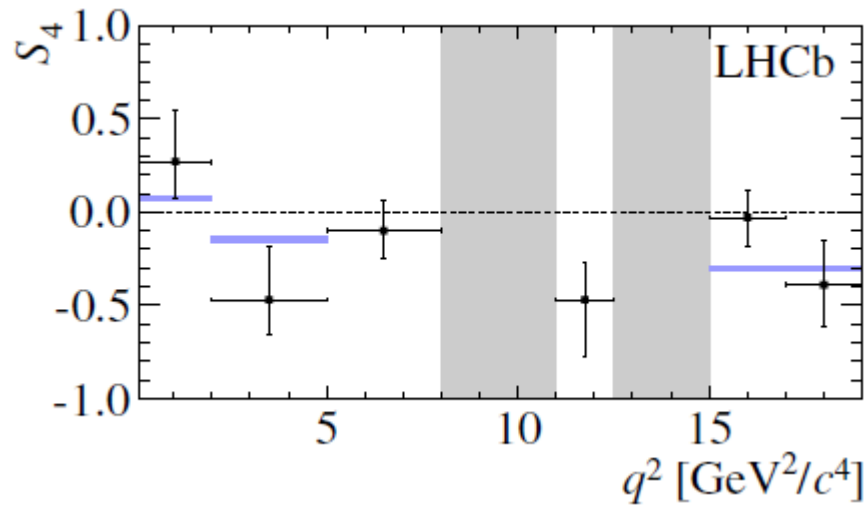
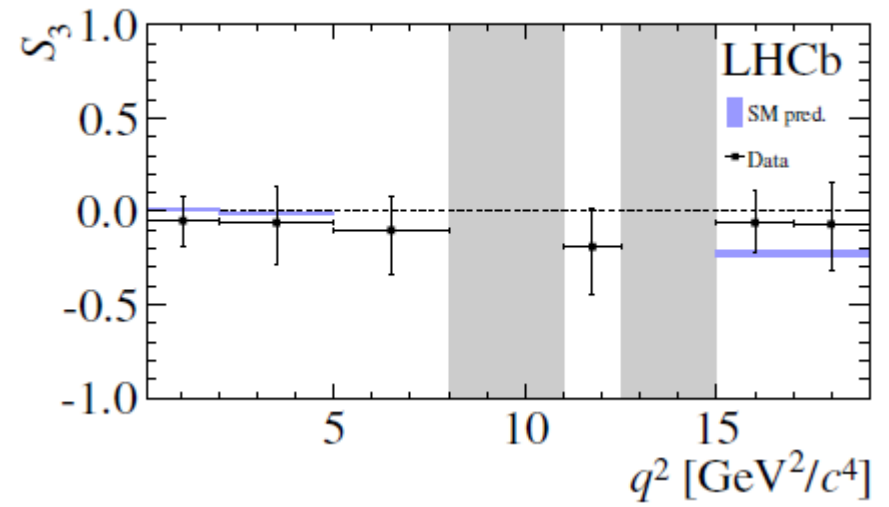
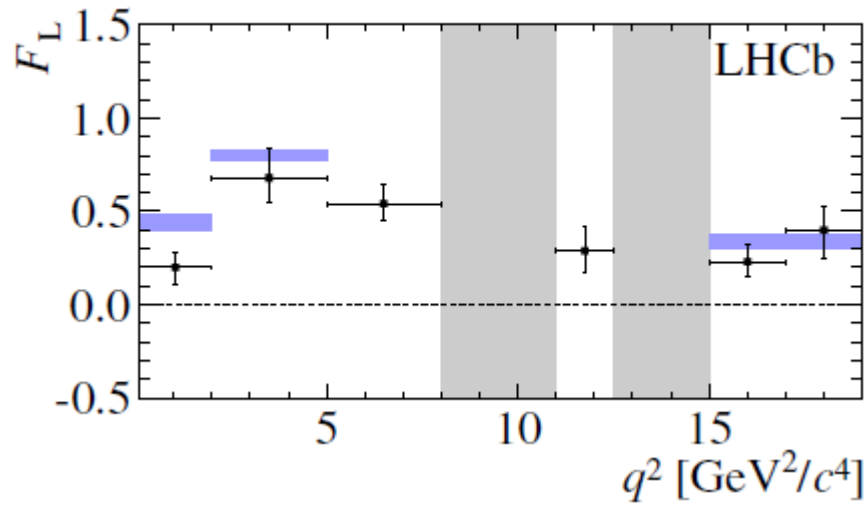
- Differential branching fraction and full angular analysis

$$\mathcal{B}(B_s^0 \rightarrow \phi \mu^+ \mu^-) = (7.97_{-0.43}^{+0.45} \pm 0.22 \pm 0.23 \pm 0.60) \times 10^{-7},$$

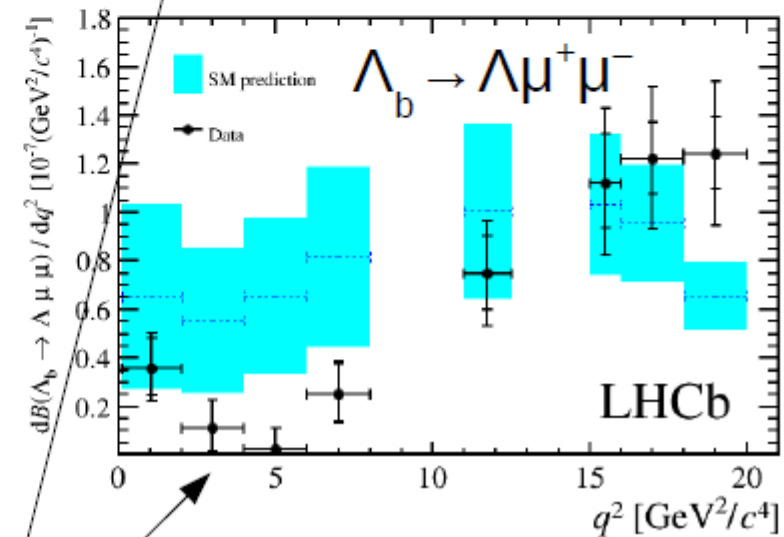
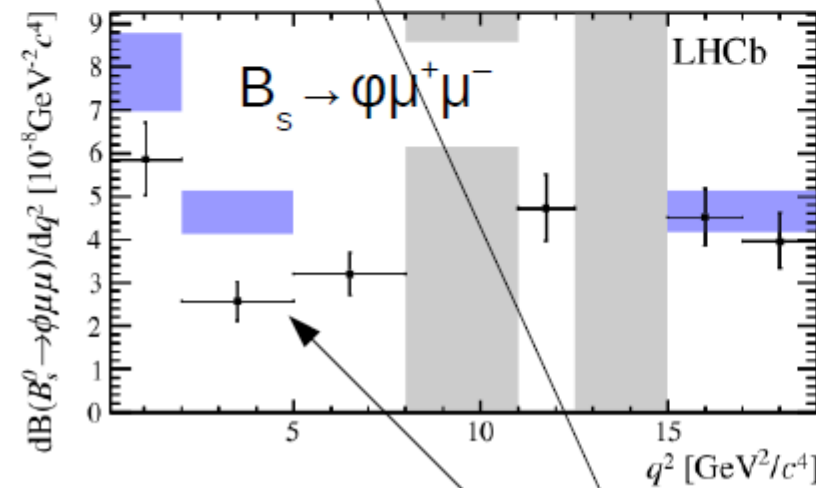
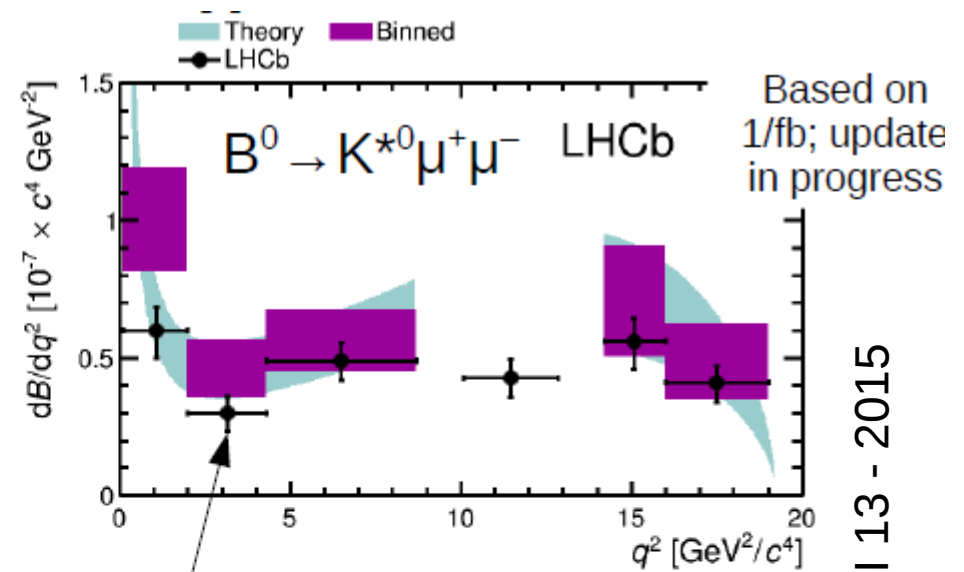
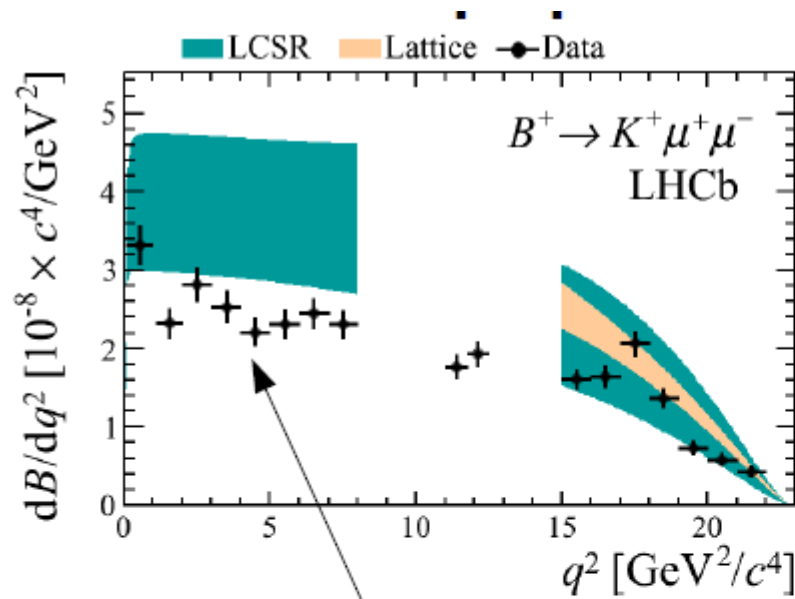


3.3  $\sigma$  inconsistency with SM





# $b \rightarrow s \mu^+ \mu^-$ branching fractions



Common trend to be below SM prediction at low  $q^2$

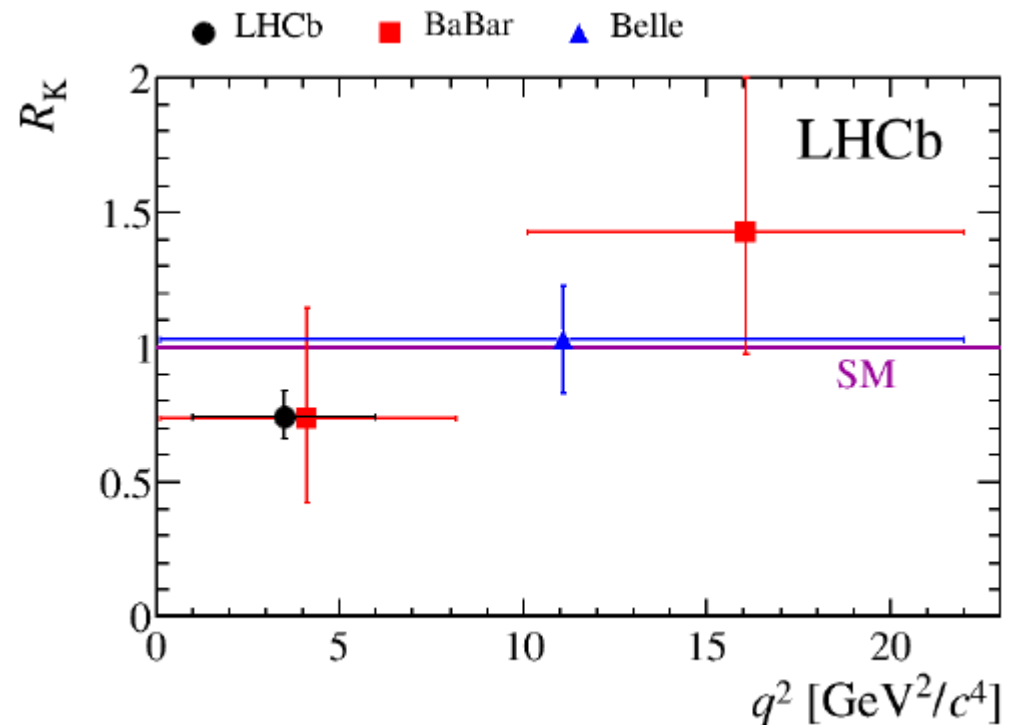
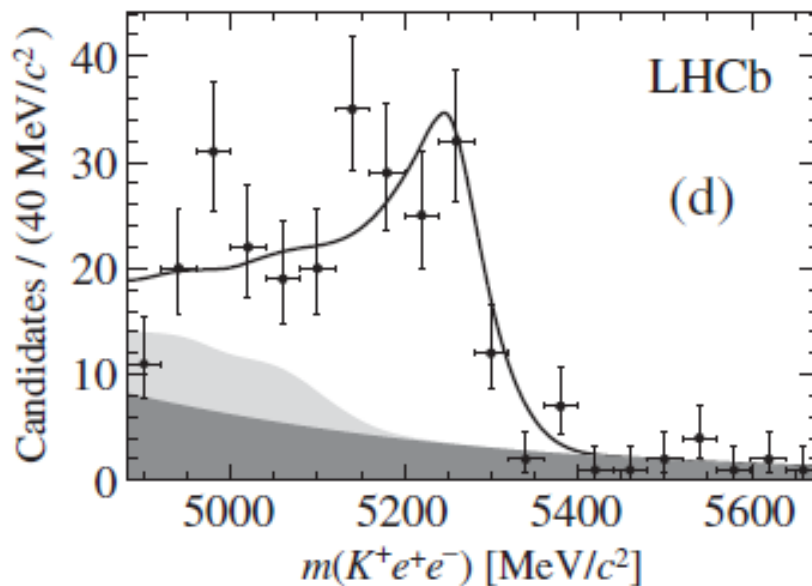
T. Gershon - GGI 13 - 2015

# Lepton universality in $B^+ \rightarrow K^+ \ell \ell$

- Normalize the BF to the  $B^+ \rightarrow K^+ J/\psi$  to cancel systematic effects mainly due to the different energy loss (see tail a low energy)
- $R_K$  measure in  $1 < q^2 < 6 \text{ GeV}^2/c^4$  is  $3\sigma$  from SM prediction

$$R_K = \left( \frac{\mathcal{N}_{K^+ \mu^+ \mu^-}}{\mathcal{N}_{K^+ e^+ e^-}} \right) \left( \frac{\mathcal{N}_{J/\psi(e^+ e^-) K^+}}{\mathcal{N}_{J/\psi(\mu^+ \mu^-) K^+}} \right) \times \left( \frac{\epsilon_{K^+ e^+ e^-}}{\epsilon_{K^+ \mu^+ \mu^-}} \right) \left( \frac{\epsilon_{J/\psi(\mu^+ \mu^-) K^+}}{\epsilon_{J/\psi(e^+ e^-) K^+}} \right),$$

$$R_K = 0.745^{+0.090}_{-0.074} (\text{stat}) \pm 0.036 (\text{syst}).$$



# Rare decays prospects

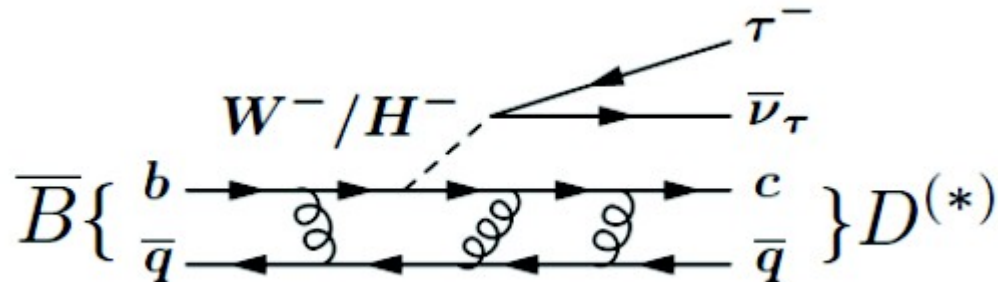
- Moving away from discovery and towards precision era
  - for  $B_s \rightarrow \mu\mu$
- Full angular analysis of  $b \rightarrow s\mu\mu$  shows several interesting  $3\sigma$  BSM effects
  - Ongoing effort to control hadronic uncertainties
- More decay modes to be studied to explore the full potential
- Combine all available informations with global fits to effective couplings (Wilson coefficients)
- $b \rightarrow d\mu\mu$  is becoming also interesting

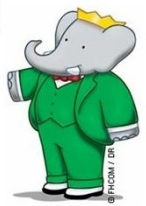
$$B \rightarrow D^{(*)} \tau \nu$$



# New physics in semileptonic decay channels

- Semileptonic decay channels are used to extract CKM matrix elements with good precision because of the factorization of the hadronic leptonic current
- In addition charged lepton universality assures that decays to  $e$ ,  $\mu$  and  $\tau$  should differ only in phase space and helicity suppression
- Many extensions of the SM instead predict a difference between lepton flavors, for instance a charged Higgs type II THDM would couple predominantly to decays into  $\tau$  leptons





# RD e RD\* in $B \rightarrow D^{(*)} \tau \nu$ : BaBar

- Can access both RD and RD\*

- Signal reconstructed in  $\tau \rightarrow \bar{\nu}_l \nu_\tau$

$$R(D) = \frac{BF(B \rightarrow D \tau \nu_\tau)}{BF(B \rightarrow D l \nu_l)_{l=e,\mu}}$$

- Normalization has same topology as signal

- Is also largest background

$$R(D^*) = \frac{BF(B \rightarrow D^* \tau \nu_\tau)}{BF(B \rightarrow D^* l \nu_l)_{l=e,\mu}}$$

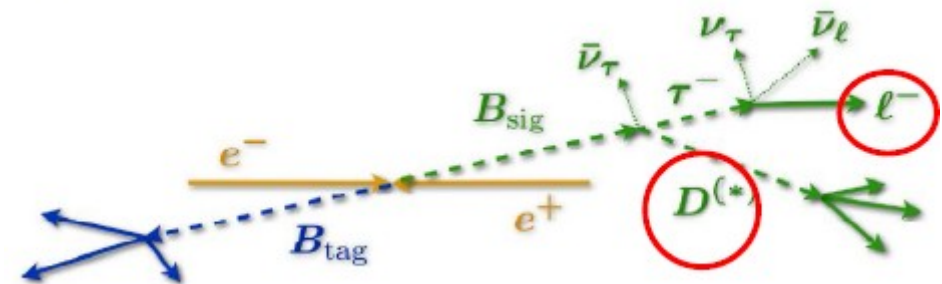
- Uncertainties cancel in the ratio

- $D^{(*)}$  reconstruction, PID, Tracking

- $|V_{cb}|$ , form factors

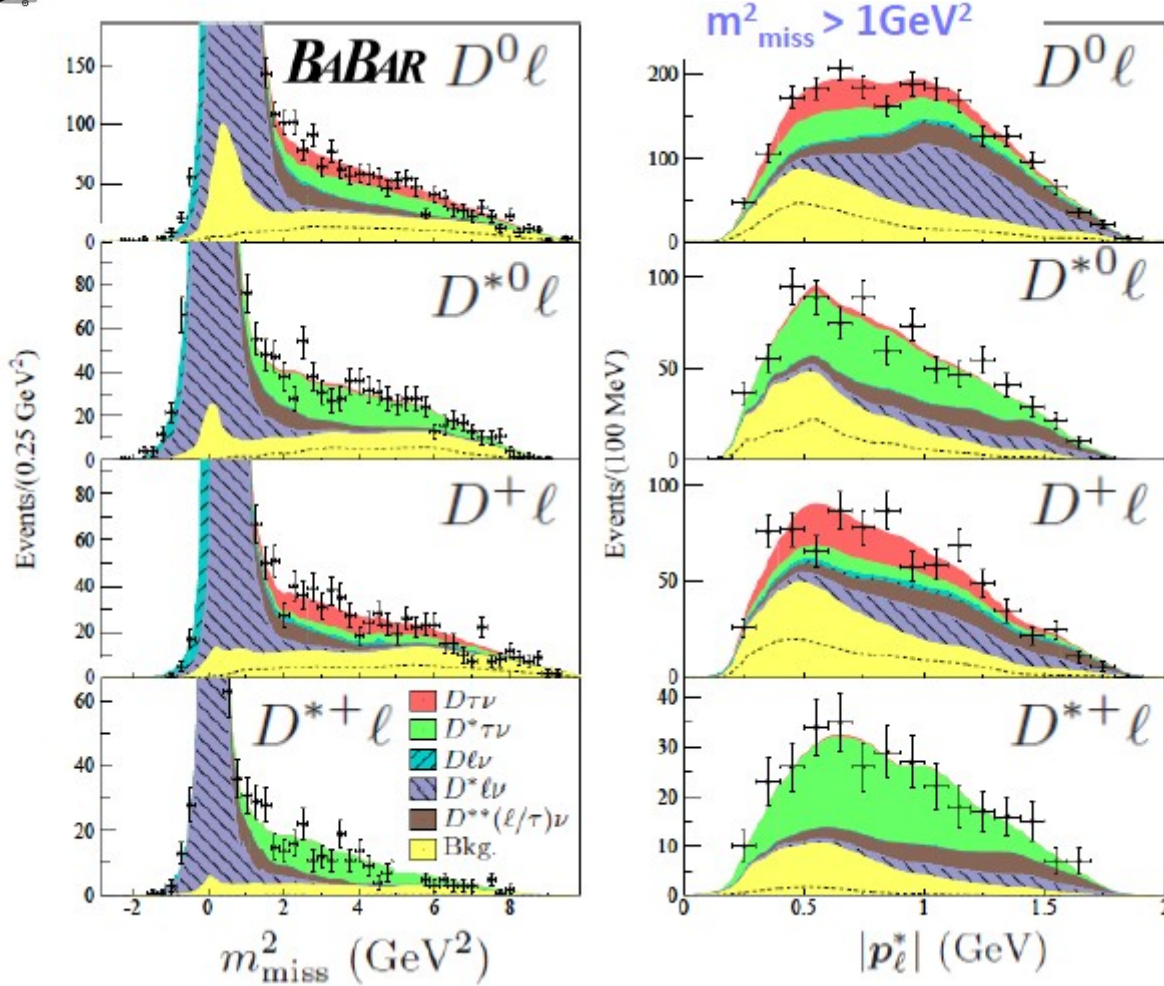
- $B_{\text{tag}}$  fully reconstructed into hadrons

- Allows to reconstruct charge and momentum of signal  $B$



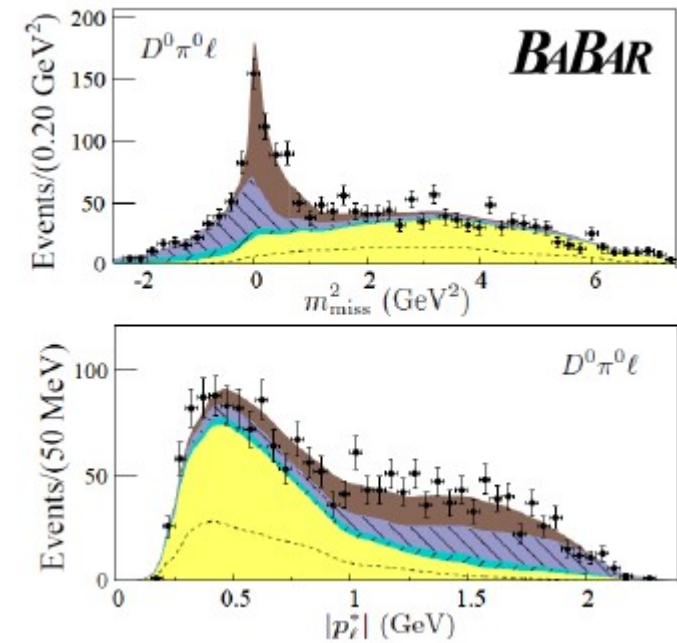
- Everything else should come from the signal  $B$

- Not possible at hadron colliders



$$m_{miss}^2 = (P_{e^+e^-} - P_{Btag} - P_l - P_{D^{(*)}})^2$$

- Main sist. related to  $D^{**}$  bkg fixed on control sample



- 4 signal samples ( $D^0, D^+, D^{*0}, D^{*+}$ ) $lv$
- 4 control samples ( $D^0, D^+, D^{*0}, D^{*+}$ ) $\pi^0 lv$  to derive  $D^{**} lv$  background.

Isospin constrained results:

$$R(D) = 0.440 \pm 0.058 \pm 0.042$$

$$R(D^*) = 0.332 \pm 0.024 \pm 0.018$$

Significance  $> 8 \sigma$



- Type II 2HDM can describe the higgs sector in the MSSM

- Differential decay rates

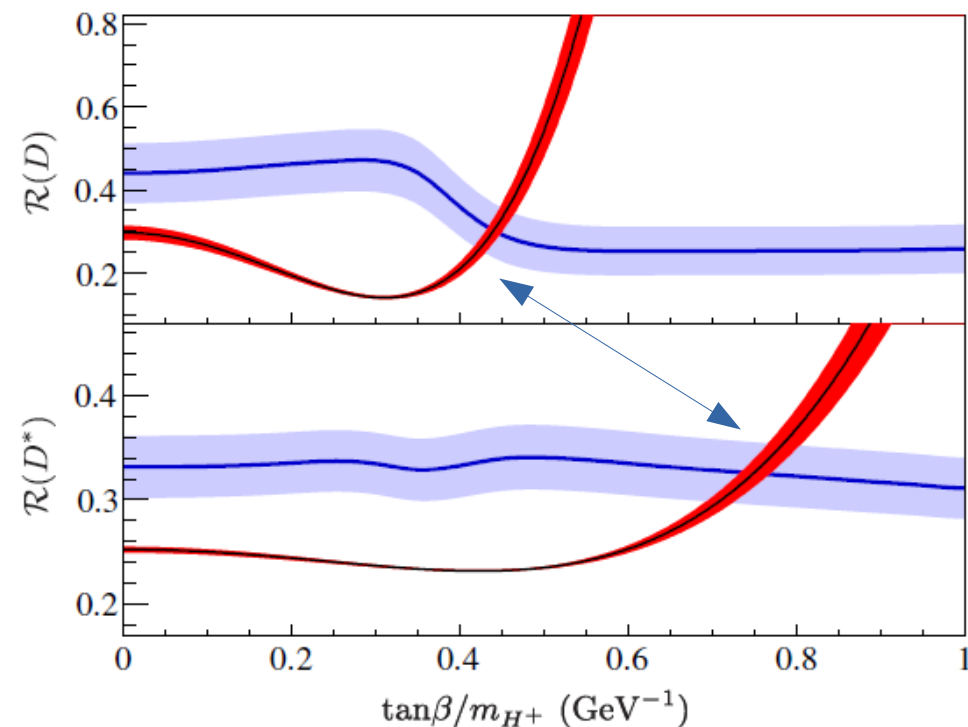
$$\frac{d\Gamma_\tau}{dq^2} = \frac{G_F^2 |V_{cb}|^2 |P_{D^{(*)}}^*|^2 q^2}{96\pi^3 m_B^2} \left(1 - \frac{m_\tau^2}{q^2}\right)^2 \left[ (|H_+|^2 + |H_-|^2 + |H_0|^2) \left(1 + \frac{m_\tau^2}{2q^2}\right) + \frac{3m_\tau^2}{2q^2} |H_s|^2 \right],$$

are modified by a charged higgs

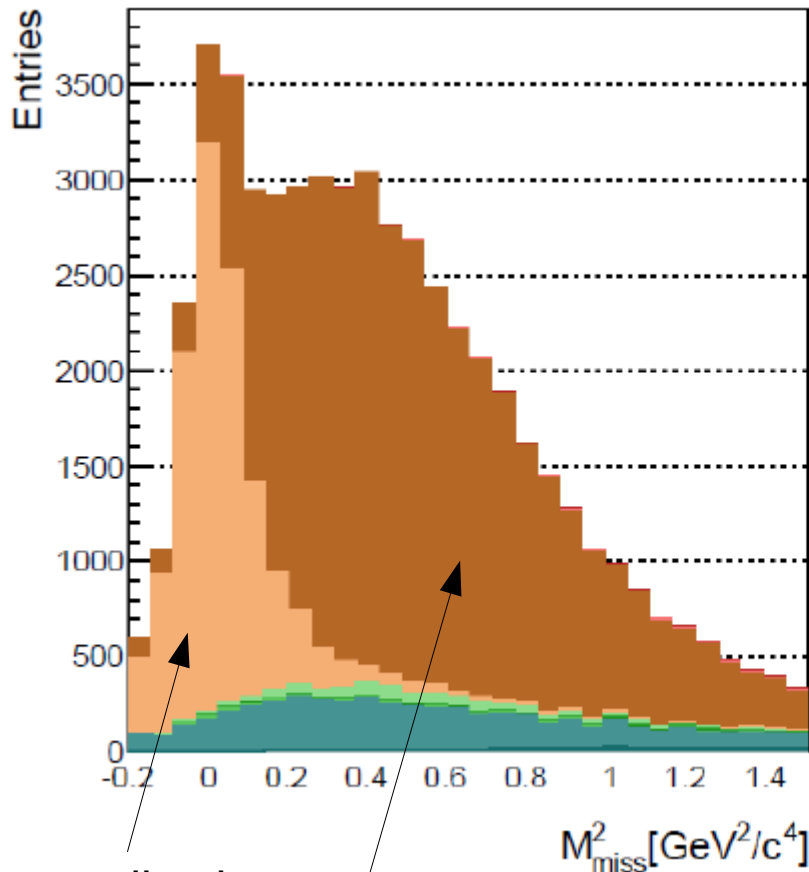
$$H_s^{2\text{HDM}} \approx H_s^{\text{SM}} \times \left(1 + (S_R \pm S_L) \frac{q^2}{m_\tau(m_b \mp m_c)}\right).$$

in a  $q^2$  dependent manner

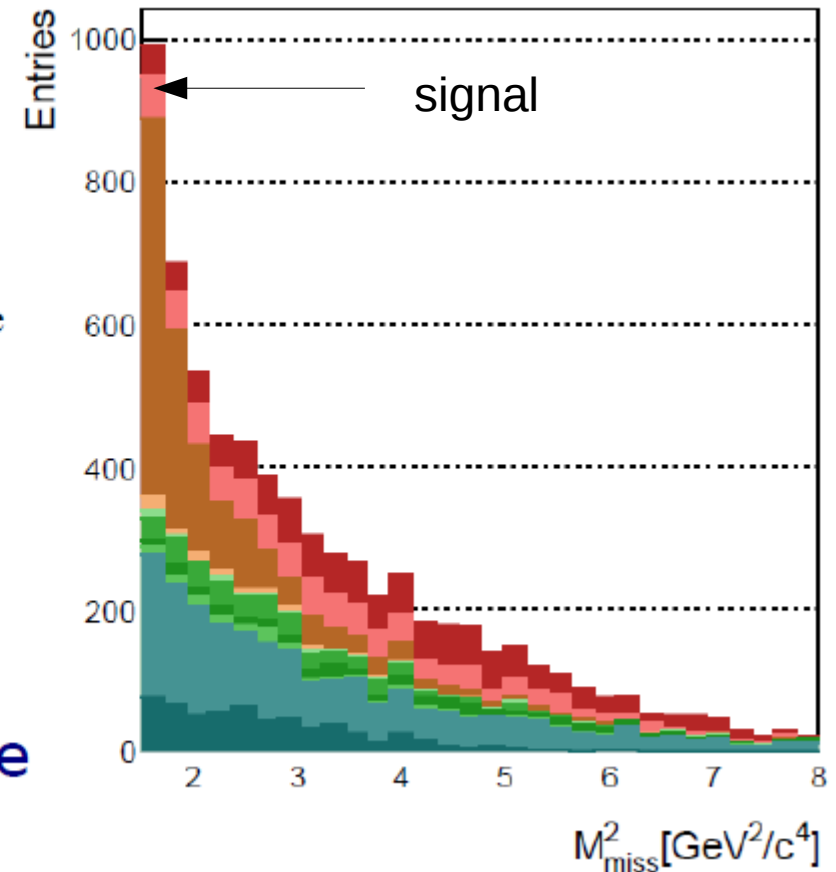
- Consider different values of  $\tan\beta/m_{H^+}$ : generate new signal shapes and measure  $R(D^{(*)})$ , compare with SM prediction
- Result is incompatible with 2HDM



- New result this summer with more sophisticated analysis
- Divide the sample in low  $m_{\text{miss}}$  to fit normalization sample and cross feed and high  $m_{\text{miss}}$  to extract signal from NN



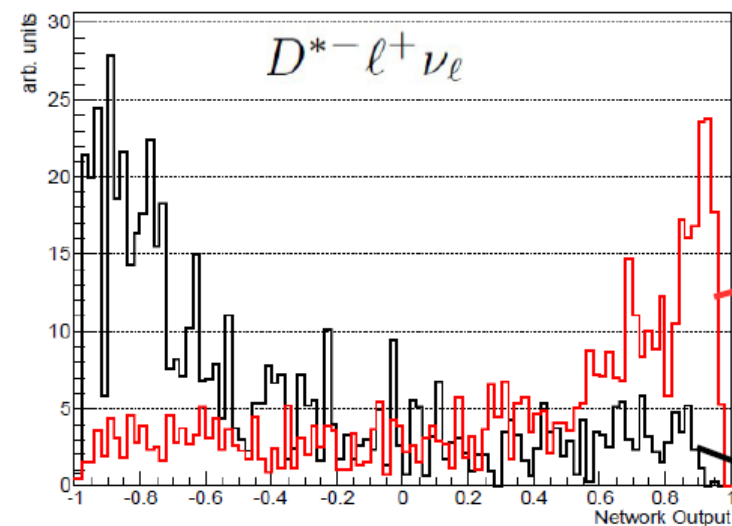
$D^0 l^-$  sample



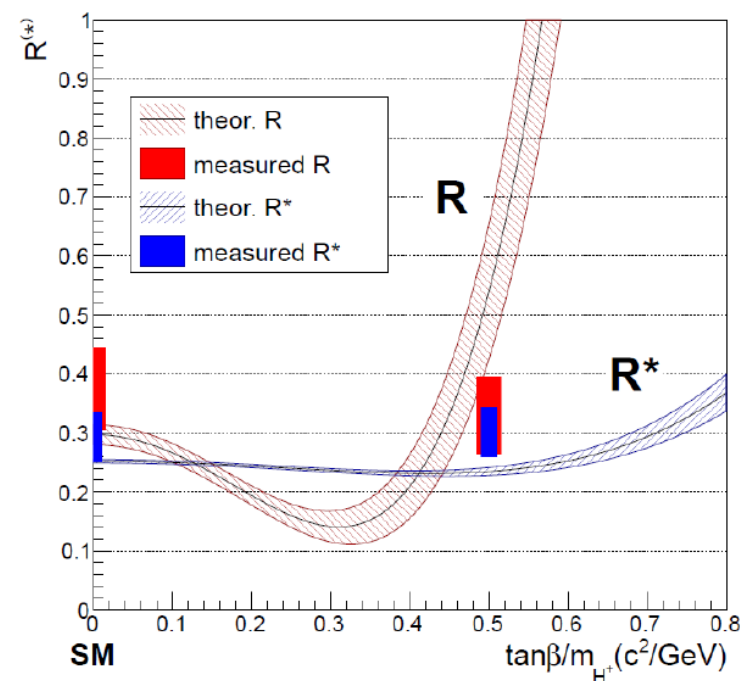
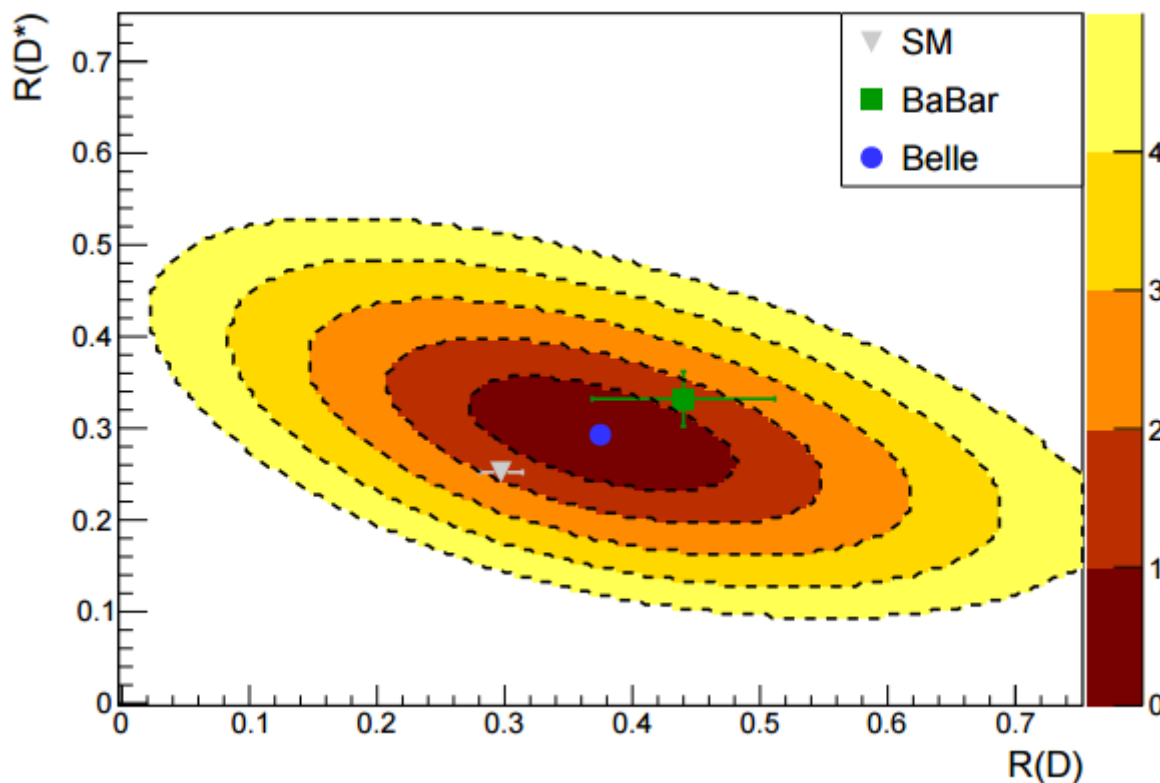
- Signal cross feed, wrong charge and combinatorial - constrained or measured from sidebands
- $D^{**}$  yield free parameter of the fit

- Result is compatible with both BaBar and SM and SM
- Repeat analysis for 2HDM type II with  $\tan\beta/m_{H^+} = 0.5 c^2/\text{GeV}$

## Neural network signal and background



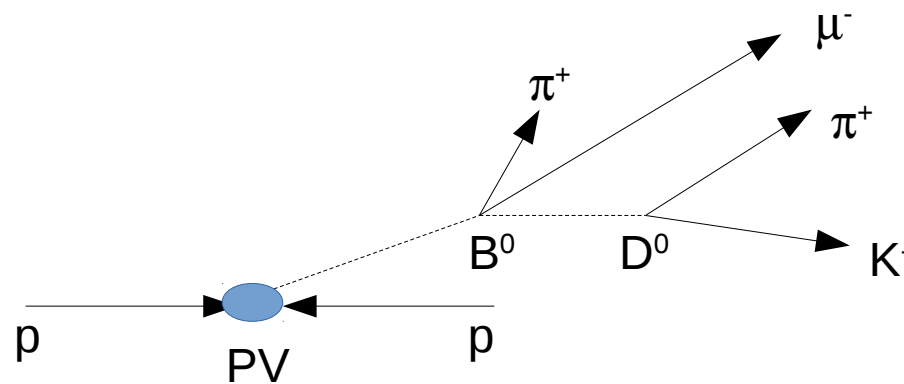
T.Kuhr FP CP 2015





- Trigger on the charm component
- Selection of signal exploiting excellent secondary vertex reconstruction and muon identification
- $\tau$  reconstruction efficiency 77.6% of the  $\mu$  normalization mode
- $B$  meson rest frame is not known: determine  $B$  direction from PV and  $B$  vertex
  - Approximate  $B$  boost along the beam direction with boost of the visible system

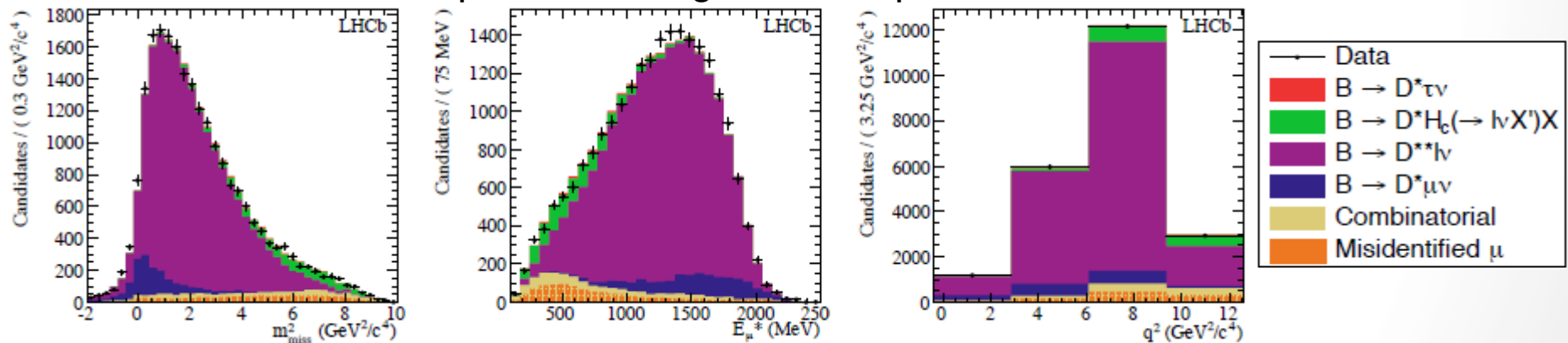
$$p_z(B^0) = \frac{m_{B^0}}{m(D^* \mu)} p_z(D^* \mu)$$



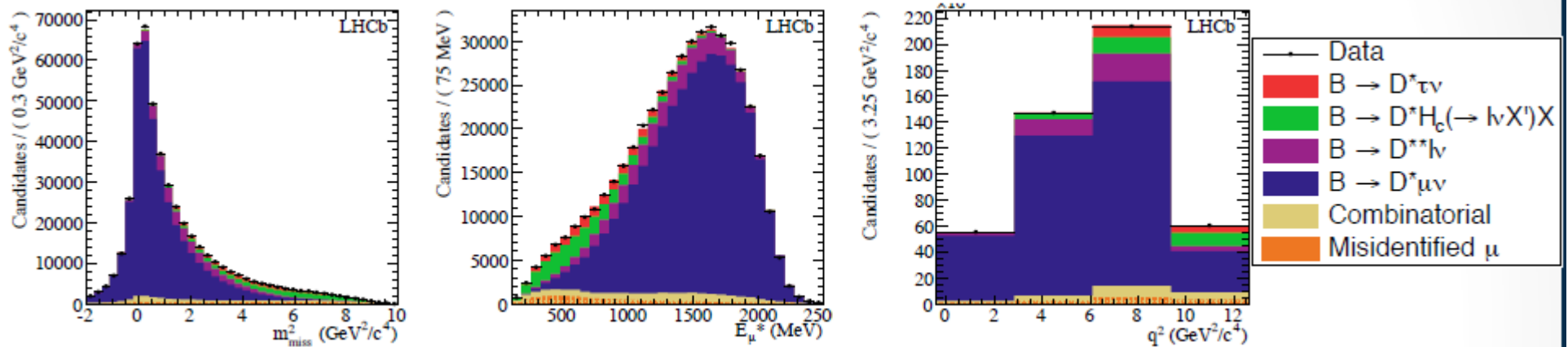


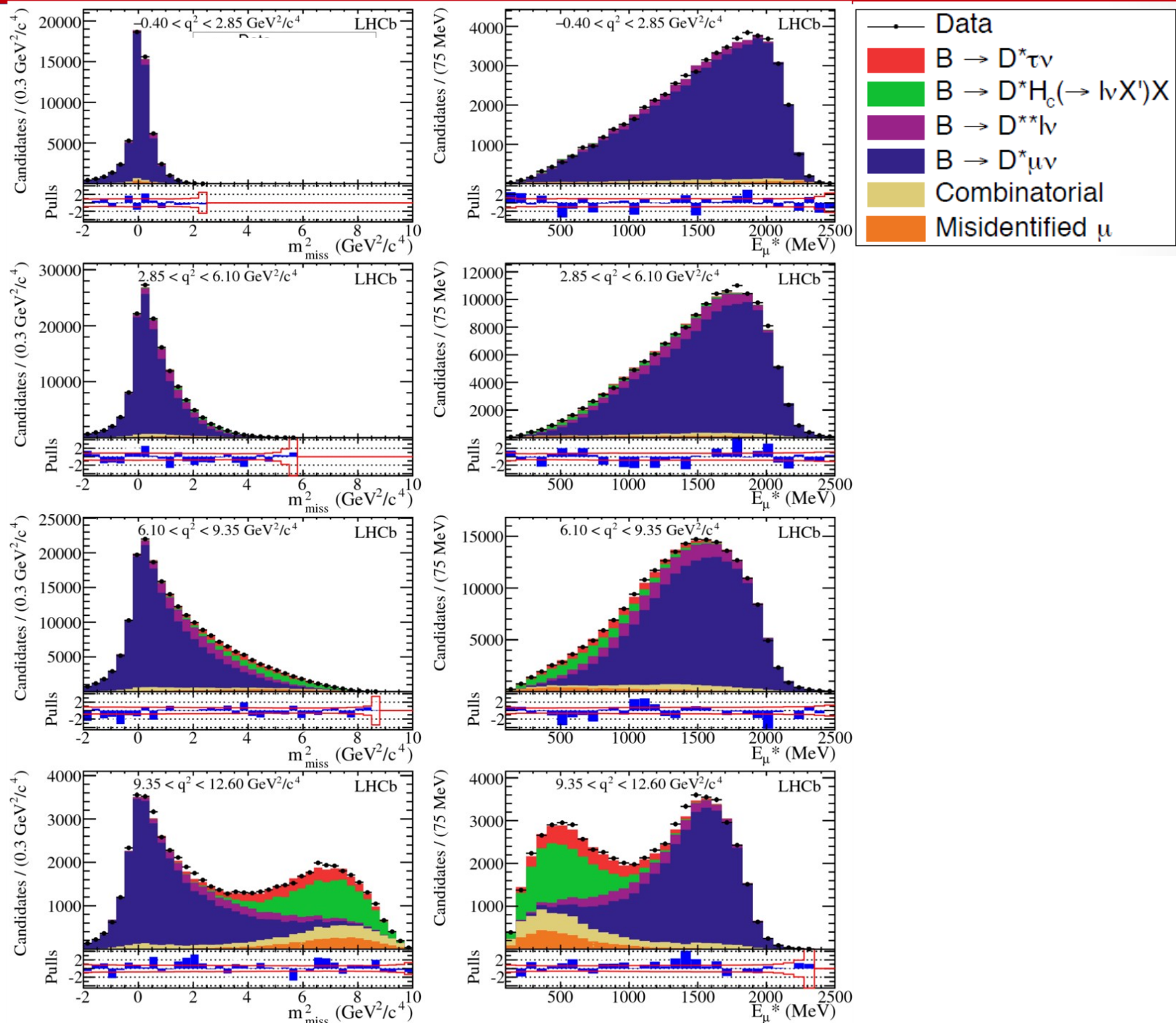
- ML fit to  $m^2_{\text{miss}}$ ,  $E_{\mu}$ ,  $q^2$  distributions with 3D templates representing signal, normalization and background sources
- Largest backgrounds from partially reconstructed semileptonic decays
  - $B \rightarrow D^{**} \mu \nu$ 
    - Known resonances modeled using template from MC, validated on  $D^{*+} \mu \pi^+$
    - Higher charm states models using ISGW2

Fit to the control sample with background templates

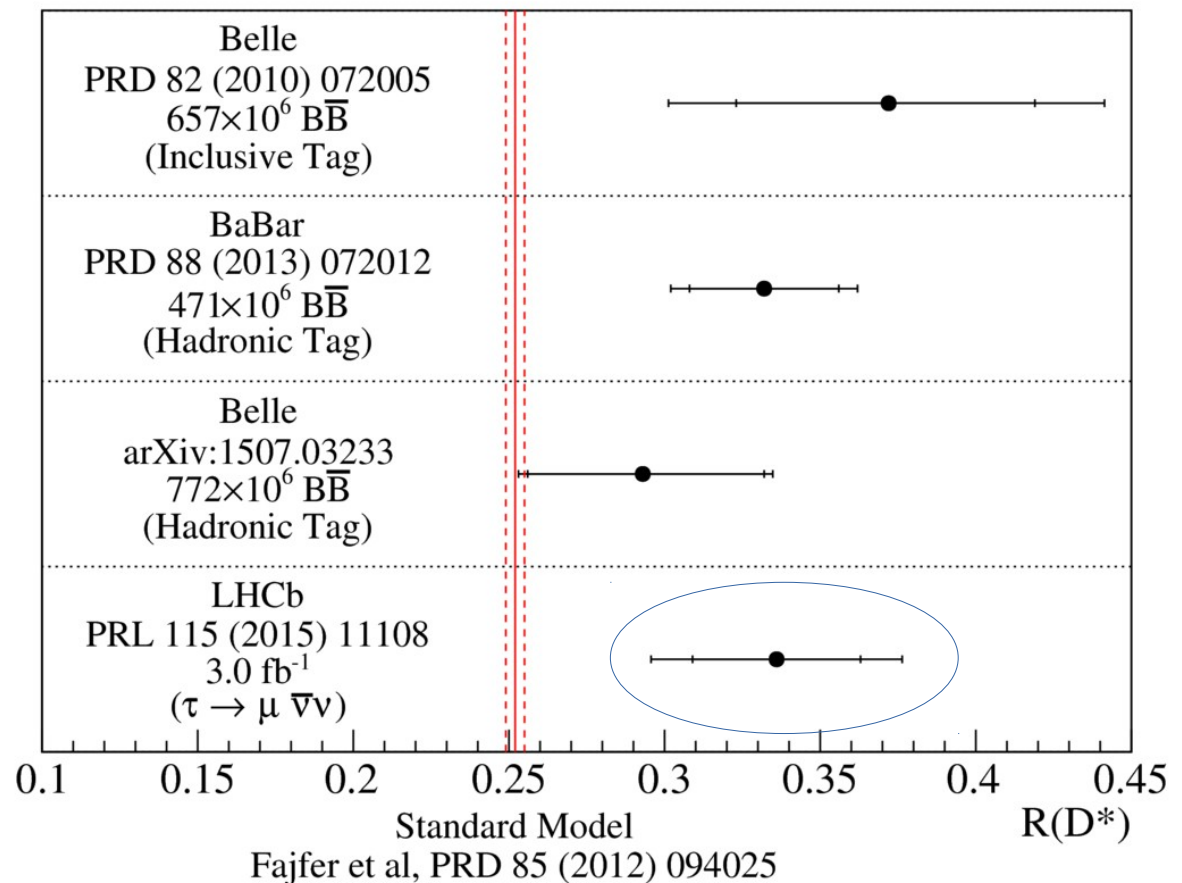


- Semileptonic decays of a second charm hadron in the event represent a significant source of background
- Empirical corrections to the templates extracted from fits to  $D^{*-}\mu\kappa^{\pm}$  control sample
- Muon misidentification and combinatorial also modeled with control samples

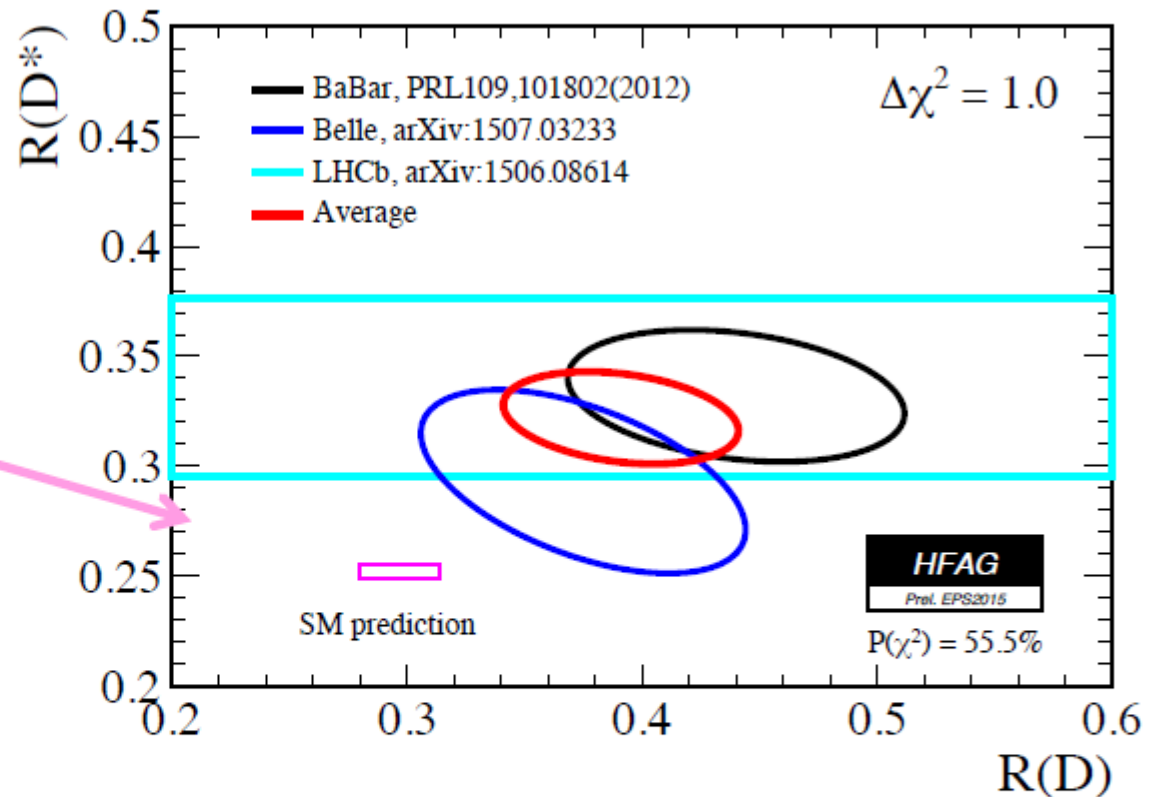




- $R(D^*) = 0.336 \pm 0.027(\text{stat}) \pm 0.030(\text{syst})$
- In agreement with previous measurement
- $2.1 \sigma$  higher than SM



M. Calvi  
EPS 2015



SM predictions

$$R(D^*) = 0.252 \pm 0.003$$

PRD 85 (2012) 094025

$$R(D) = 0.297 \pm 0.017$$

PRD 78 (2008) 014003

HFAG average

$$R(D^*) = 0.322 \pm 0.018 \pm 0.012$$

$$R(D) = 0.391 \pm 0.041 \pm 0.028$$

$$\text{Corr}(D, D^*) = -0.29$$

- Difference with the SM predictions at  $3.9 \sigma$  level.

# Lepton Universality prospects

- Expand the physics program searching for other modes
  - Measure a ratio similar to  $R_k(B \rightarrow K \mu^+ \mu^- / B \rightarrow K e^+ e^-)$  in decays to  $K^*, \phi, \Lambda$ , etc
  - Expand  $D\tau\nu$  to  $D^*, D_s, \Lambda_c$
  - Search for lepton number violation  $B \rightarrow \tau\nu, K\tau\nu, Ke\nu$  etc.
    - Mostly Belle-II

# Conclusions

- Many recent exciting results
- Sides and Angles
  - Reached high precision and consistency
  - There is margin to improve sensitivity on both  $\beta$  and  $\gamma$ ,  $|V_{ub}/V_{cb}|$ ,  $\Delta m_d$ ,  $\Delta m_s$  with more statistics (Belle II, LHCb upgrade)
- Rare decays
  - Hints of deviations from SM pointing in the same direction?
  - Is there a way to explain the observed pattern of deviations
- Lepton universality
  - Inconsistencies with 2HDM Type II confirmed
  - measurements pointing toward a bigger coupling to taus and electrons w.r.t. muons
    - puzzling
  - Need to improve experimental sensitivity: more decay channels and more data

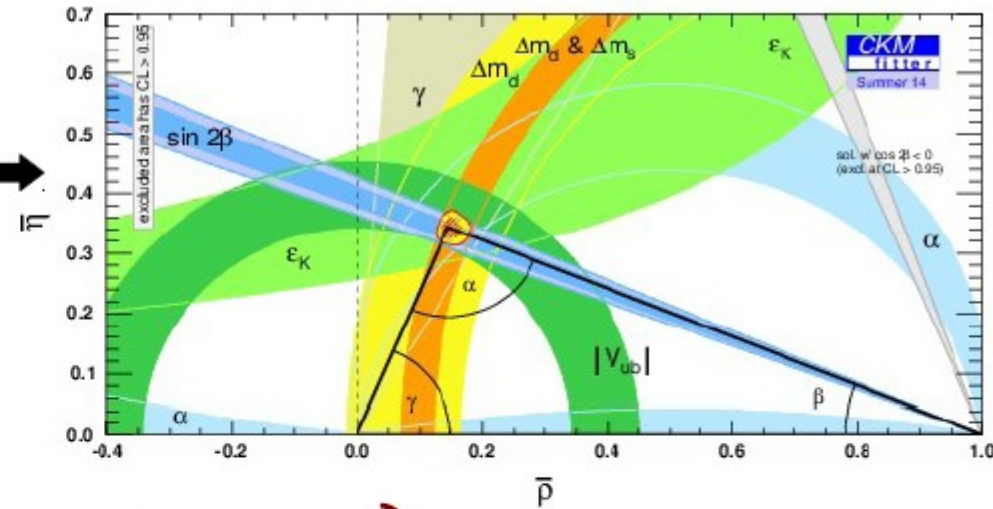
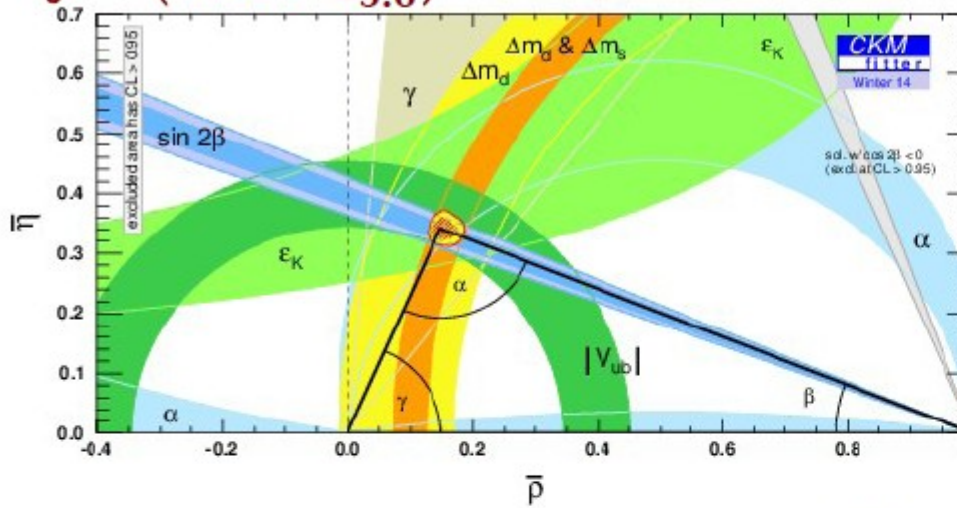


# Backup

# $\gamma$ determination

Trablesi, FPCP 2015

$$\gamma = (70.0^{+7.7}_{-9.0})^\circ$$

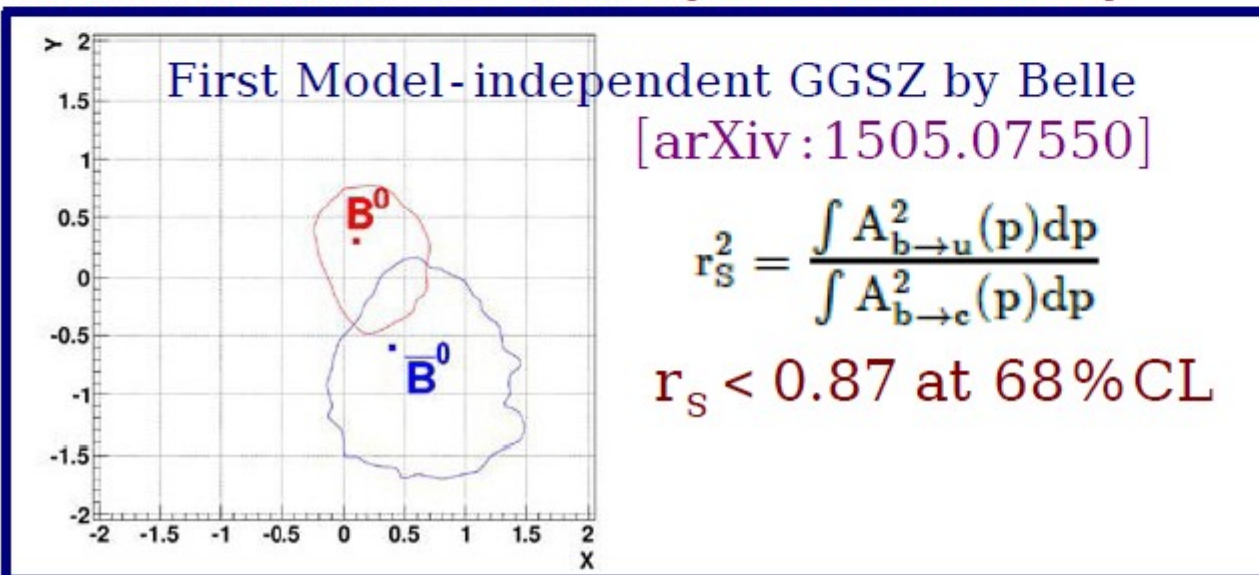


mostly from charged  $B \rightarrow D^{(*)} K$ :

$$\left. \begin{aligned} \gamma_{\text{BaBar}} &= (70 \pm 18)^\circ \\ \gamma_{\text{Belle}} &= (73^{+13}_{-15})^\circ \\ \gamma_{\text{LHCb}} &= (75 \pm 9)^\circ \end{aligned} \right\} \gamma = (73.2^{+6.3}_{-7.0})^\circ$$

$$r_B = 0.0970 \pm 0.0063$$

from neutral  $B \rightarrow DK^*$ : [Markus Rohrken]

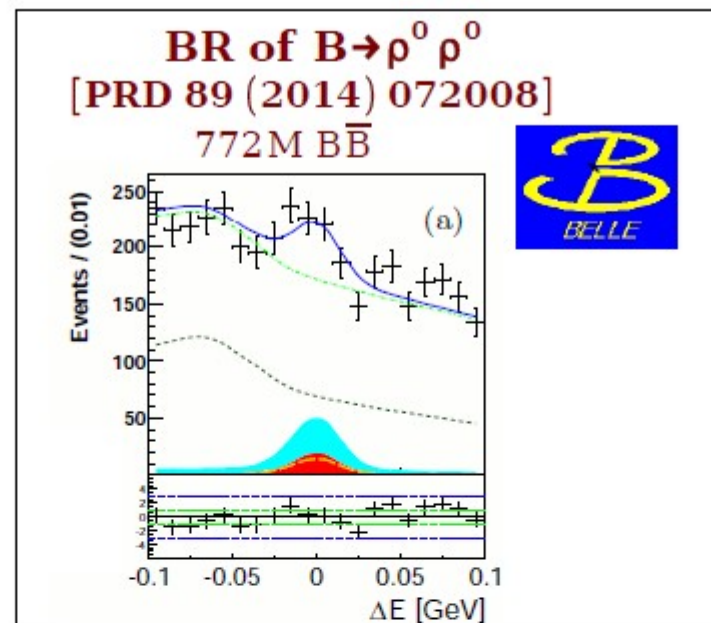
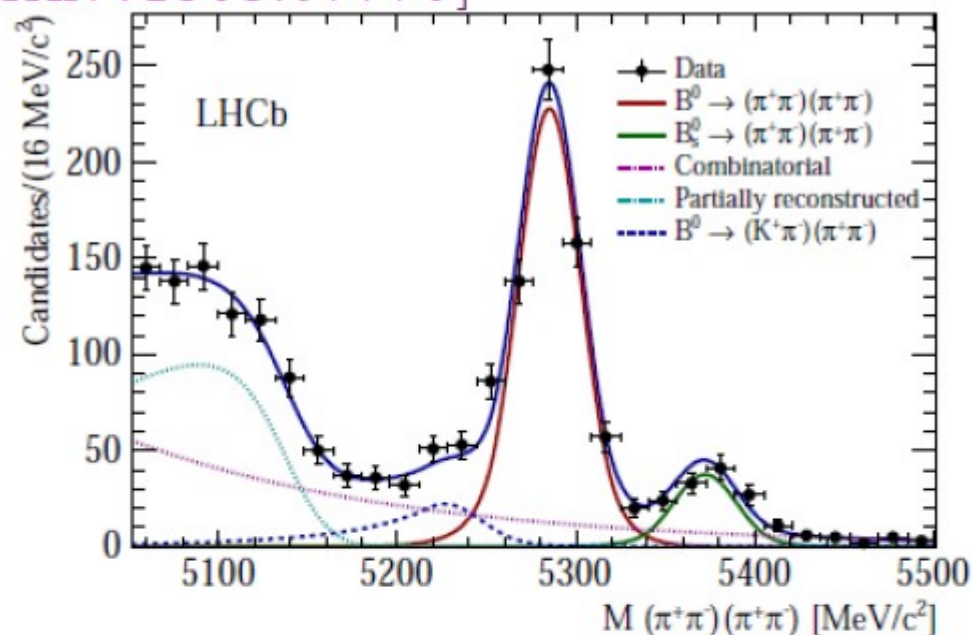


GLW+ADS by LHCb  
[arXiv:1407.8136]

$$r_S = 0.236^{+0.043}_{-0.052}$$

# $B \rightarrow \rho^0 \rho^0$ and $\alpha$ determination Trablesi, FPCP 2015

[arXiv:1503.07770]



Decay mode	Signal yields 2011	Signal yields 2012
$B^0 \rightarrow (\pi^+\pi^-)(\pi^+\pi^-)$	$185 \pm 15 \pm 4$	$449 \pm 24 \pm 7$
$B^0 \rightarrow (K^+\pi^-)(\pi^+\pi^-)$	$1610 \pm 42 \pm 5$	$3478 \pm 62 \pm 10$
$B^0 \rightarrow (K^+K^-)(K^+\pi^-)$	$1513 \pm 40 \pm 8$	$3602 \pm 62 \pm 10$
$B_s^0 \rightarrow (\pi^+\pi^-)(\pi^+\pi^-)$	$30 \pm 7 \pm 1$	$71 \pm 11 \pm 1$
$B_s^0 \rightarrow (K^-\pi^+)(\pi^+\pi^-)$	$40 \pm 10 \pm 3$	$96 \pm 14 \pm 6$
$B_s^0 \rightarrow (K^+K^-)(K^-\pi^+)$	$42 \pm 10 \pm 3$	$66 \pm 13 \pm 4$

600 signal  $4\pi^\pm$  events

$$B(B \rightarrow \rho^0 \rho^0) = (0.94 \pm 0.17 \pm 0.09 \pm 0.06) \times 10^{-6}$$

$$f_L = 0.745^{+0.048}_{-0.058} \pm 0.034$$

angular analysis

